

THE CLUSTER OF GALAXIES ABELL 1060

by

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## ABSTRACT

This thesis describes the implementation and results of a programme of photometry of the relatively nearby southern cluster of galaxies and X-ray source, Abell 1060 = Hydra I.

Photographic photometry in B and V wavebands of galaxies in a  $\sim 22$  degree<sup>2</sup> region centred on this cluster have been obtained on the basis of measurements of four U.K. Schmidt telescope plates by the COSMOS measuring machine at the Royal Observatory, Edinburgh. The feasibility of using the COSMOS machine to derive reliable isophotal magnitudes of bright and faint galaxies is demonstrated.

Multiaperture photoelectric (U, B, V and R-band) photometry has been secured of sixteen galaxies in the field of the cluster, including seven NGC and IC systems. Besides providing photometric calibration, these observations permit the total magnitudes and colours of several galaxies to be derived.

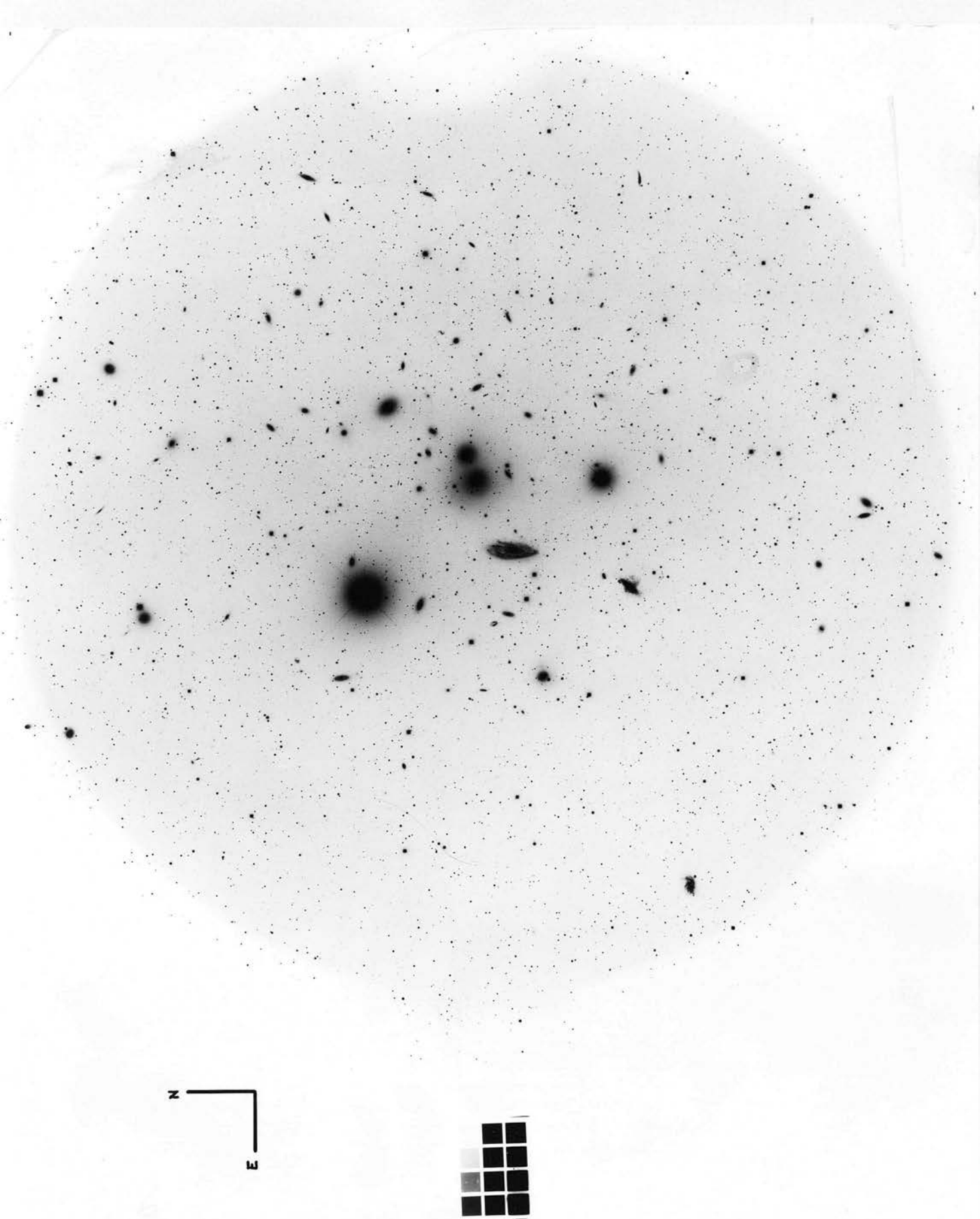
The positions and magnitudes of over 5000 galaxies, supplemented in many cases by morphological classifications, by micrometric measurements of dimensions and alignment, and by redshift data, have been analysed to determine the static and dynamical properties of A1060. The cluster is symmetric, centrally condensed, and rich in disk galaxies; members segregate radially by morphological type, with spiral galaxies generally farthest from the centre.

The V-band luminosity function of A1060 has been constructed. It is similar in form but different in detail from those of other clusters. The colour-magnitude relation of early-type galaxies in

A1060 is outlined, and is shown to be inconsistent with one prediction of a standard theory of the morphological evolution of galaxies and clusters.

Visual searches for low surface brightness members of A1060 are described. It is suggested that fewer such galaxies exist than an extrapolation of the luminosity function for all galaxy types to faint magnitudes would predict.

Peculiar properties of four galaxies near the centre of the cluster are interpreted in terms of the influence of their environments. Three are anaemic spirals, for which there is strong evidence that an era of gas depletion has recently begun. The first-ranked member, N3311, is an extended system of low mean surface brightness, and may represent a D galaxy in the early stages of evolution. These observations suggest that the cluster Abell 1060 is relatively young in dynamical terms.



The cluster of galaxies Abell 1060. Photographed by M.G. Smith at the prime focus of the 4-metre reflector at Cerro Tololo Interamerican Observatory. The central  $0.74 \text{ degree}^2$  of the cluster is represented, at a scale of  $18.56 \text{ arcsec mm}^{-1}$ . (For further particulars of plate 995, see page 169).



## DECLARATION

This thesis has been composed entirely by the undersigned candidate and with the exception of some of the astronomical observations described in Chapter Four, all work presented herein is his own. (Regulation 2.4.15)

Most of the material presented in Chapter Four was published in Monthly Notices of the Royal Astronomical Society 190, 631 (1980) under the title "Multiaperture Photoelectric Photometry of Galaxies in the Hydra I Cluster", by R.J. Smyth and R.S. Stobie. (Regulation 2.4.11)

Signed:

R.J. Smyth.

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## CHAPTER 1

### INTRODUCTION

#### 1.1 The Character of Confucius

The apparent contradiction of learning is not a contradiction.

The question is whether or not it is a contradiction.

What is the relationship between learning and thought?

Learning without thought is labour lost;

Thought without learning is perilous.

One of the main points of the Analects is that

learning is not a mechanical process.

- K'ung Fu-tzu (Confucius)

These are the main points of the Analects. Bk. ii

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## CHAPTER 1

### INTRODUCTION

#### 1.1 The Clustering of Galaxies

The apparent congregation of galaxies in clusters, observable over immense distances down to the limiting magnitude of present sky surveys, is one of the most striking characteristics of the large-scale distribution of visible matter in the universe. Studies of galaxies and their clustering are consequently vital to our comprehension of the physical world.

The cataloguing and characterisation of galaxies was underway long before their true nature came to be generally recognised. Even the mere thirty-three verifiable extragalactic "nebulae" in the third and final catalogue of Messier (1781) suffice to illustrate a concentration of bright galaxies in a region of the sky that contains the constellations of Virgo and Coma Berenices. The much larger compilation of W. Herschel (1784), with more than a thousand truly extragalactic objects, revealed the apparent large-scale gradients in numbers of galaxies across the northern sky (due to optical obscuration in the plane of our own Galaxy); this compilation was the forerunner of the General Catalogue (J. Herschel 1864) which reveals several clusters (Waters 1873). At least four distinct groups, in Virgo, Ursa Major and two in Fornax, are apparent in the distribution of the nearby galaxies listed in the catalogue of Shapley & Ames (1932).



Recognition of the vast numbers of clusters came with the advent of photographic sky surveys of the northern hemisphere, in particular the Lick Observatory 20-inch Astrographic Survey as inspected by Shane and Wirtanen (Shane & Wirtanen 1954, 1967; Shane 1956; Shane, Wirtanen & Steinlin 1959; Kron & Shane 1974) and the National Geographical Society - Palomar Observatory Sky Survey. Stringent criteria became necessary in the study of large volumes of space in order to restrict selection to manageable numbers. Abell (1958) crudely characterised 2712 of the very richest clusters visible on Palomar plates, and from the same material Zwicky et al. (1961-8) catalogued some 9700 recognisable associations, representing an estimated  $1.6 \times 10^7$  galaxies. Since the commissioning of the ESO and UK Schmidt telescopes, complementary surveys of the southern sky have been in progress, although specialised lists of southern clusters already exist (e.g. Duus & Newell 1977; Braid & MacGillivray 1978).

The surface distribution of galaxies as recorded by northern surveys has been subjected to statistical analysis since the 1950s (e.g. Neyman & Scott 1952). Correlation techniques involving angular separations have been developed by Peebles and his co-workers, and correlation functions determined which give the "excess" probabilities of finding 2-galaxy and 3-galaxy configurations (discussed by Fall et al. 1976). These functions vary as simple power laws of the angular scale, indicating no preferred scale of clustering at sizes up to 5 arc degrees. Clusters, therefore, may be only the high-density cores of extensive clouds of galaxies which may represent (cf. Soreira & Peebles 1977) substantially more than 80%

of all visible galaxies. Separate correlation analyses of the different morphological populations of galaxies reveals a greater clustering tendency among elliptical galaxies than other types (Davis & Geller 1976).

Important though statistical studies are, they are comparatively uninformative regarding the composition and dynamics of individual systems of galaxies. It has become increasingly evident in recent years that the location of a galaxy with respect to the centre of a cluster may play a significant part in determining its appearance and properties. A new classification scheme for galaxies was introduced by van den Bergh (1976) which explicitly recognised the existence of peculiar disk systems (anaemic spirals) that are often found in clusters. Strom & Strom (1977, 1978a,b,c,d, 1979a) have demonstrated that elliptical galaxies found in the cores of rich clusters where the density of galaxies is high tend to be of smaller characteristic size than those ellipticals of comparable central surface brightness that are found in the outskirts of clusters; they also find a corresponding variation of the disk colours of galaxies that are classified as S0 types (Strom & Strom 1978e, 1979b).

## 1.2 Properties of Rich Clusters

The subject of rich clusters and their properties has been extensively reviewed by Abell (1975), Bahcall (1977a) and by van den Bergh (1977a). Some aspects are considered in later chapters, and only brief outlines of definitions of cluster size and type are necessary here.

By general acquiescence the term "rich cluster" has come to

mean any sufficiently populous grouping of galaxies comparable to those catalogued by Abell (1958). The apparent population of a cluster may depend on the surroundings relative to which it is viewed, but the structures of rich clusters do vary widely and several different classification schemes have been proposed. Shapley (1933) drew an analogy between loose, irregular systems of galaxies and open star clusters of this Galaxy; Abell (1965) uses the term irregular clusters, to distinguish them from the more centrally concentrated regular clusters. Bahcall (1977a) introduced an intermediate class to make the transition smoother. The regular, intermediate and irregular clusters correspond quite closely to the compact, medium-compact and open categories of Zwicky et al. (1961-8). Morgan (1961) and de Vaucouleurs (1961b) each suggested discriminating according to galactic content, but this in itself is not a satisfactory criterion because it is difficult or impossible to classify galaxies at the distance of remote clusters; the same objection may be raised to the scheme of Melnick & Sargent (1977). However, the composition of clusters appears to correlate with their form.

Rood & Sastry (1971; RS) proposed a useful morphological classification system for clusters which bears strong resemblance to the original Hubble scheme for galaxies; six classes are distinguished according to the projected distribution of the ten brightest members in each cluster. Bautz & Morgan (1970; BM) exploit the optical contrast of the brightest galaxy relative to other member galaxies in their five-part system, in which the dominance of the first-ranked member is greatest in type I clusters and least in type III. White (1978) has described an elaboration of the Bautz-Morgan system that

is applicable to poor as well as to rich clusters. The BM classification, although vulnerable to chance superpositions of non-members, is of special interest in connection with possible secular changes in the distribution of luminosity among brighter cluster members. In common with RS types, it is often supplemented with estimates of cluster richness on Abell's (1958) scale, for neither scheme incorporates a "population" parameter.

Clusters, like galaxies, have no definite boundary, and three general approaches have been made to the problem of defining cluster size. Zwicky et al. (1961-8) adopt the mean diameter of an isopleth that represents a fixed fraction of the field density around a cluster, but this definition is distance-dependent (Peach & Beard 1969; Noonan 1972). An alternative, applicable only to regular and symmetric clusters, is to fit functions with characteristic length scales to the projected number density profiles of galaxies. Although the isothermal (Emden) sphere model has usually been adopted, and is considered later in this thesis, previous representations have involved the de Vaucouleurs (1948) " $r^{\frac{1}{4}}$ " law, Maxwellian (i.e. tri-variate normal) distributions (Scott 1962), and King (1966) dynamical models (e.g. Rood et al. 1972). Each of the above represent the observed profiles almost equally well over the definable range of number density, which in clusters is a mere two decades because of the contaminating effects of non-members. Thirdly, specific mean projected distances between cluster members may be estimated (e.g. Noonan 1972; Austin & Peach 1974; Hickson 1977), which have the advantage that they do not depend on the structure of the cluster. Attempts have been made to exploit the apparent sizes of clusters, on

the assumption that rich clusters are of fixed dimensions, to determine the cosmological deceleration parameter of cosmological world models. They have not, however, been conspicuously successful (e.g. Bruzual & Spinrad 1978).

### 1.3 Cluster Luminosity Functions

The relevance of the frequency distribution of luminosities of galaxies that represent particular samples of space, to a wide range of extragalactic problems has been stressed by Schechter (1976). These distributions are termed luminosity functions and may be determined for individual clusters.

Luminosity distributions have in principle been available since very early times. The Herschel symbols  $vB$ ,  $B$ ,  $F$  etc. were a rough attempt at the establishment of a nebular photometry. Ames (1923) published preliminary total magnitudes of 20 galaxies in the Virgo-Coma region, estimated by visual comparison with stellar images on photographic plates. The existence of a consistent local maximum in the numbers of bright members of the Virgo, Perseus and Ursa Major I clusters was commented on by Baade (1928).

A luminosity function for bright nearby galaxies that approximates a gaussian curve was derived by Hubble in the course of a famous study (Hubble 1936a,b). Doubts about its adequacy, however, followed Shapley's (1939) discovery of the Sculptor and Fornax dwarf members of the Local Group. Zwicky argued that no maximum in the luminosity function could exist, and presented evidence (Zwicky 1957) that the total number of galaxies per unit volume inside clusters



increased exponentially towards faint magnitudes. Subsequent studies (e.g. Abell 1976; Oemler 1974) have shown this to be a reasonable approximation, except at the bright end. Abell (1962, 1976) and Schechter (1976) have proposed alternative representations; these are contrasted with the Zwicky function in Figure 4 of Bahcall (1977a).

Three immediate questions concerning the luminosity functions of galaxies have yet to be answered. The first concerns which of the analytic representations is the most satisfactory; the second, whether it is constant from place to place in the universe (cf. Felten 1977); the third, whether it is constant in time. Universality of the luminosity functions of galaxies in rich clusters is an assumption implicit in, for example, the determination of the luminosity function of all galaxy systems from single galaxies to rich clusters (Bahcall 1979), otherwise the number of galaxies in a group in the range of the first few magnitudes (e.g. Abell 1958) would be a poor indicator of the total population of the group. Universality is obviously an attractive hypothesis, but recent photometry of rich clusters, and numerical simulations of their evolution, casts some doubt on the assumption. Further photometry of rich clusters to faint limits is required, and a substantial part of this thesis is devoted to the problem of photometry of one particular cluster that is sufficiently nearby for very faint members to be examined.

#### 1.4 The Organisation of this Thesis

It has been found necessary to dispense with a strictly chronological approach to presentation.



In Chapter 2, the distribution of galaxies in and around the Abell 1060 cluster is considered on three scales: within a region of dimensions 45 by 23 arc degrees, on the basis of published data; in a smaller area encompassing  $\sim 100$  degree<sup>2</sup>, examined on original photographic plates; and within a central  $\sim 22$  degree<sup>2</sup> that was scrutinised with the aid of the COSMOS machine. The structure of the cluster as inferred from these counts of galaxies is described, and the geometric parameters of the galaxies are investigated.

Chapter 3 presents a compilation of the radial velocities of galaxies in the field of this cluster. Additional redshifts of galaxies are determined from measurements of objective prism spectra.

Photoelectric photometry of galaxies in A1060 is detailed in Chapter 4. The (U-B) colour-magnitude relation for early-type galaxies in the cluster is derived and the total blue-band magnitudes of several members are estimated.

The derivation of photographic magnitudes of a large number of galaxies on the basis of measurements made by the COSMOS machine is discussed in some detail in Chapter 5. The photoelectric observations of Chapter 4 are shown to provide a satisfactory means of calibrating this photometry. The reliability of the V-band magnitudes is demonstrated by comparable reductions of a second plate.

In Chapter 6 an assessment of the combined photometric and optical information is presented. The luminosity function of A1060 is constructed and compared with those of other clusters. The properties of the brightest galaxy are discussed in terms of a theory

of evolution, and in terms of the previously constructed colour-magnitude diagram. Three peculiar spiral galaxies in the cluster are examined.

An account of a search for dwarf galaxies in the Hydra I cluster is given in Chapter 7.

In the eighth and final chapter, some previous results are employed to determine the average mass-to-luminosity ratio of the composition of the cluster; the main conclusions of previous chapters are summarised, and the implications discussed briefly.

The basic photometric data, morphological classifications and finding charts in respect of the brighter galaxies in the vicinity of A1060 are presented in a pair of Appendices that follow the main text.

Table 1.1

## Abbreviations and Catalogues

Abbreviations

AM	Abell & Mihalas (1966)
CM	Coarse Measurement <u>or</u> Coarse Mode (re COSMOS)
COSMOS	Acronym for the system at ROE that measures the Co-Ordinates, Sizes, Magnitudes, Orientations and Shapes of images on photographic plates (Pratt 1977)
CTIO	Cerro Tololo Interamerican Observatory
ESO	European Southern Observatory
HMS	Humason, Mayall & Sandage (1956)
MM	Mapping Measurement <u>or</u> Mapping Mode (re COSMOS)
ROE	Royal Observatory, Edinburgh
RI	Relative Intensity
SAAO	South African Astronomical Observatory
SRC	Science Research Council of the U.K.
UKST(U)	United Kingdom Schmidt Telescope (Unit)
VVC	See RC2

Catalogues

## a) Prefixes

3U	Giacconi et al. (1974)
4U	Forman et al. (1978)
2A	Cooke et al. (1978)
A	Abell (1958)
BS	Yale Catalogue of Bright Stars
IC $\equiv$ I	Dreyer (1895, 1908)
MCG	Vorontsov-Velyaminov et al. (1962-8)
NGC $\equiv$ N	Dreyer (1888)
SAO	Smithsonian Astrophysical Observatory Star Catalogue

## b) Non-prefixes

CGCG	Zwicky et al. (1961-8)
RC1	de Vaucouleurs & de Vaucouleurs (1964)
RC2	de Vaucouleurs et al. (1976)
RNGC	Sulentic & Tifft (1973)

## CHAPTER 2

### STATIC PROPERTIES

#### 2.1 A1060 Overviewed

The cluster of galaxies Hydra I, located at approximately  $\alpha_{1950} = 10^{\text{h}} 34^{\text{m}}$ ,  $\delta_{1950} = -27^{\circ} 16'$  (revised estimates of position and redshift are quoted herein) is a relatively nearby southern cluster visible through a moderately rich star-field at a galactic latitude of  $+26.5^{\circ}$ . It appears to have been recorded as a group of nine galaxies during compilation of the New General Catalogue of Galaxies and Clusters of Stars (Dreyer 1888), and is the fifth entry among twenty-five in a list devoted to groups of galaxies drawn from both galactic hemispheres (Shapley 1933). In the standard catalogue of rich clusters found north of declination  $-27^{\circ}$  (Abell 1958), Hydra I is the unrepresentative 1060th of 2712; it is too southerly and too local for inclusion in the homogeneous subsample of this catalogue. Only two clusters belong to Abell's distance class  $D = 0$ , and of these A1060 has the lowest redshift. By contrast, Hydra II = A732 was for many years the cluster with the highest measured velocity of recession and many past allusions to the "Hydra cluster" refer to Hydra II (e.g. Allen 1973).

The distribution of galaxies in Hydra I was first investigated by Zwicky (1941, 1942), who emphasised its spherical symmetry and its structural similarity to the Coma = A1656 cluster, and quoted the first observational determination of its redshift ( $3950 \text{ km s}^{-1}$  for the member NGC 3309). Zwicky determined the radial dependence of the

number density distribution of galaxies from counts on 0.5m Schmidt telescope plates, and estimated the structural index of the cluster on the isothermal sphere model to be four minutes of arc. His counts were later analysed by Sharov (1959) to deduce the three-dimensional structure of Hydra I. Kwast (1966) undertook independent counts of galaxies on 1.2m Schmidt plates, in two wavebands, and for the brighter systems confirmed Zwicky's estimate of the structural index. Noonan (1972) reanalysed all this data to derive a mean radius of the cluster of  $\bar{r}_c \approx 48$  minutes of arc.

Abell (1958) assigned to A1060 a richness class of  $R = 1$ , implying between 50 and 79 members within the two magnitude range faintwards of the third brightest member and within  $4.6 \times 10^5 / cz$  mm.  $\approx 103/z$  arcseconds of the cluster centre on the scale of the Palomar Sky Survey plates, where  $z$  denotes mean redshift. 51 galaxies were counted by Sandage & Hardy (1973) within an angular diameter of  $137(1+z^2)/z$  arcseconds and in the magnitude range from the third brightest galaxy in A1060 to  $2\frac{1}{2}$  magnitudes fainter. Zwicky (1941) inferred "roughly 200" members in the range of the brightest three magnitudes from counts that covered some 50 square degrees in the vicinity, and suggested that this represented less than half the population of the cluster, which he considered to have a diameter of at least ten degrees.

Hydra I has been classified by Rood & Sastry (1971) as a type C and by Bautz & Morgan (1970) as a type III cluster, indicative of a core-halo structure with no dominant galaxies. Although the evident symmetry and central concentration of Hydra I (e.g. Kwast 1966;



Rudnicki 1967) is characteristic of rich regular clusters (Abell 1965), it also has properties more typical of irregular clusterings such as that in Virgo. Modern X-ray observations have confirmed (Ives & Sanford 1976) that A1060 (= 3U 1044-30 = 4U 1033-26 = 2A 1033-270) is a relatively weak source of X-ray emission, with a 2-10 keV luminosity  $L_x \approx 2 \times 10^{43} \text{ erg s}^{-1}$  on the basis of both the third UHURU (Giacconi et al. 1974) and second Ariel V (Cooke et al. 1978) satellite observations (Gursky & Schwartz 1977; Bahcall 1977a). A discrete emitter at radio wavelengths has also been reported in the direction of the centre of A1060, which may be associated with either or both of the giant member galaxies N3309 and N3311 (Disney & Wall 1977).

Photographic photometry of bright members of the system of globular clusters centred on N3311 (Smith & Weedman 1976), supplemented by photometric parameters for the three brightest "elliptical" galaxies, were the basis of a recent estimate,  $\Delta(\text{Hya I}) = 40.5(1 \pm .10)$  Mpc, of the distance of the Hydra I cluster (de Vaucouleurs 1977b). An arithmetic error and an underestimate of the total magnitude  $B_T$  of N3311 impair this determination, but it is unlikely to be in error by more than 40%. The working approximation  $\Delta(\text{Hya I}) = 40$  Mpc is adopted hereafter, except where otherwise implied by a stated value of the Hubble 'constant'  $H_0$ . For a cosmological recessional velocity of  $3390 \text{ km s}^{-1}$  (Chapter 3),  $\Delta = 40$  Mpc corresponds to a local value of  $H_0$  of  $85 \text{ km s}^{-1} \text{ Mpc}^{-1}$ .

## 2.2 Visual Impressions of the Large-Scale Structure

Since an angular extent of  $10^\circ$  (Zwicky 1941) corresponds to a linear



diameter of 7 Mpc at 40 Mpc, the question arises as to whether A1060 is part of a galaxy complex or supersystem such as those described by de Vaucouleurs (1971), of characteristic scale length  $\Lambda \approx 10$  Mpc.

A study of the surface distribution of galaxies brighter than the 13th magnitude in the catalogue of Shapley & Ames (1932) reveals a large ellipsoidal cloud complex containing galaxies of various types, at low galactic latitudes in Hydra (de Vaucouleurs 1956a,b). This cloud extends westwards from A1060 ( $l = 270^\circ$ ,  $b = +26^\circ$ ) and the Antlia group ( $l = 273^\circ$ ,  $b = +19^\circ$ ) across some  $40^\circ$  in galactic longitude and  $20^\circ$  in latitude, nearly parallel to the galactic plane. A plot of the distribution of the galaxies in this area that are listed in the Revised New General Catalogue (Sulentic & Tifft 1973) reveals all three associations (Figure 2.1).

Visvanathan & Griersmith (1979) have decomposed the main Hydra complex into clouds A and B with a velocity differential of  $1400 \text{ km s}^{-1}$ , the north-westerly cloud B with a redshift similar to that of A1060. However, an enlarged sample of redshifts drawn from Sandage (1978) suggests that this demarcation is artificial (inset, Figure 2.1). The mean velocity of recession of the Hydra cloud as a whole is about  $2250 \text{ km s}^{-1}$  relative to the Local Group, and most members are probably in the foreground of A1060 and the Antlia group. The loose N2997 group = G8 (de Vaucouleurs 1975) at  $l = 250^\circ$ ,  $b = +19^\circ$  is closer still, and is represented by the peak at  $\sim 500 \text{ km s}^{-1}$  in the frequency distribution of velocities.

To gain an independent insight into the large-scale structure

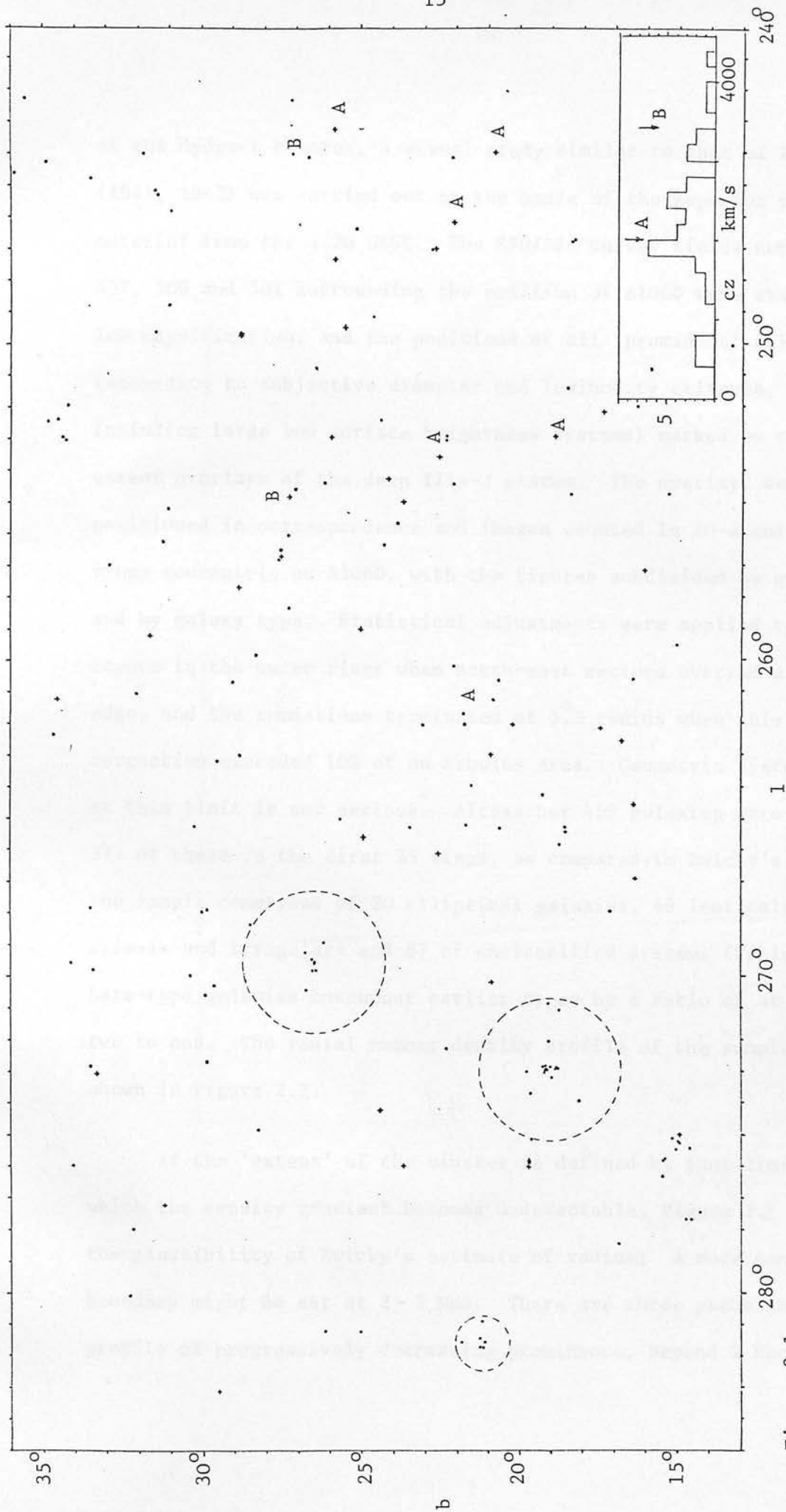


Figure 2.1

The surface distribution in galactic coordinates (1,b) of galaxies listed in the RNGC (Sulentic & Tifft 1973) which are visible in a  $45^\circ \times 23^\circ$  region encompassing the Hydra I cluster (269 $^\circ$ ,26 $^\circ$ ), Antlia group (273 $^\circ$ ,19 $^\circ$ ) and N3557 group (282 $^\circ$ ,21 $^\circ$ ). The Hydra cloud extends westwards from  $l \approx 275^\circ$ . The inset shows the frequency distribution of "indicative velocities" relative to the Local Group of those galaxies (outside the dotted circles, and shown as crosses in the plot) of known redshift. "A" refers to Hydra cloud A, "B" to cloud B after Visvanathan & Griersmith (1979).

of the Hydra I cluster, a visual study similar to that of Zwicky (1941, 1942) was carried out on the basis of the superior plate material from the 1.2m UKST. The ESO/SRC Survey fields number 436, 437, 500 and 501 surrounding the position of A1060 were examined at low magnification, and the positions of all 'prominent' galaxies (according to subjective diameter and luminosity criteria, and including large low surface brightness systems) marked on transparent overlays of the deep IIIa-J plates. The overlays were then positioned in correspondence and images counted in 10-arcminute rings concentric on A1060, with the figures subdivided by quadrant and by galaxy type. Statistical adjustments were applied to the counts in the outer rings when north-east sectors overran a field edge, and the summations terminated at  $5.3^{\circ}$  radius when this correction exceeded 10% of an annulus area. Geometric distortion at this limit is not serious. Altogether 459 galaxies were recorded, 371 of these in the first 25 rings, as compared to Zwicky's 374; the sample comprised of 20 elliptical galaxies, 49 lenticulars, 303 spirals and irregulars and 87 of unclassified systems (Table 2.1). Late-type galaxies outnumber earlier types by a ratio of at least two to one. The radial number density profile of the sample is shown in Figure 2.2.

If the 'extent' of the cluster is defined by that limit at which the density gradient becomes undetectable, Figure 2.2 confirms the plausibility of Zwicky's estimate of radius; a more conservative boundary might be set at 2 - 3 Mpc. There are three peaks in the profile of progressively decreasing prominence, beyond 2 Mpc; all

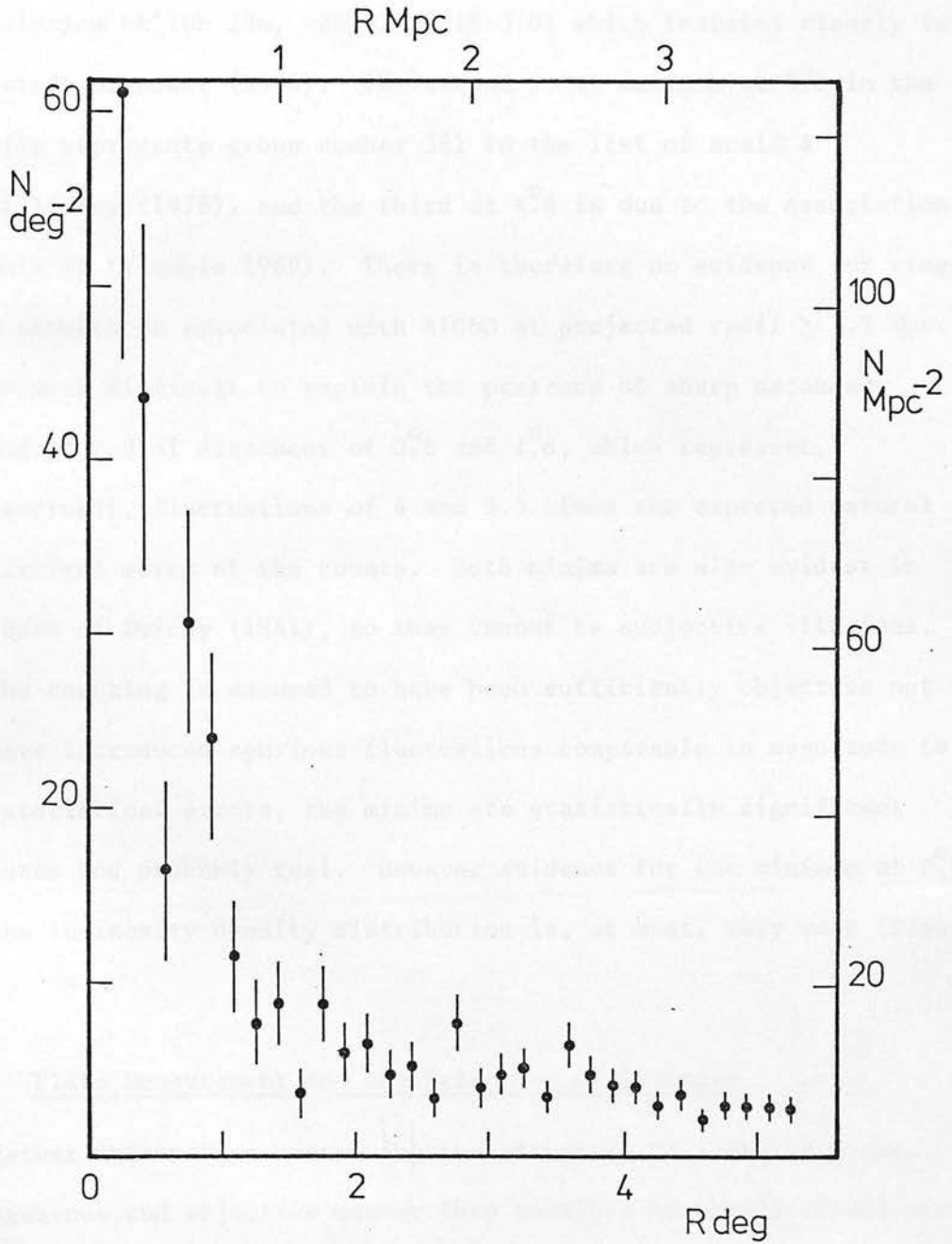


Figure 2.2

The radial surface density distribution of "prominent" galaxies within  $5.3^\circ$  of the centre of A1060 as seen in projection. Discrete peaks in the profile beyond  $\sim 1.5$  Mpc are due to identifiable groups of galaxies.

are attributable to localised clusterings. The most obvious, at a projected separation of  $2.75^\circ$  from the centre of A1060, is the group of galaxies at 10h 23m,  $-26^\circ 24'$  (1950.0) which features clearly in the study of Kwast (1966). The second local maximum at  $3.6^\circ$  in the profile represents group number 381 in the list of Braid & MacGillivray (1978), and the third at  $4.4^\circ$  is due to the association Klemola 11 (Klemola 1969). There is therefore no evidence for ring-like structures associated with A1060 at projected radii  $\geq 1.5$  Mpc. It is more difficult to explain the presence of sharp secondary minima at radial distances of  $0.6^\circ$  and  $1.6^\circ$ , which represent, respectively, fluctuations of 4 and 3.5 times the expected natural statistical error of the counts. Both minima are also evident in the data of Zwicky (1941), so they cannot be subjective illusions. If the counting is assumed to have been sufficiently objective not to have introduced spurious fluctuations comparable in magnitude to the statistical errors, the minima are statistically significant features and probably real. However evidence for the minimum at  $0.6^\circ$  in the luminosity density distribution is, at most, very weak (Figure 6.7).

### 2.3 Plate Measurement and the Selection of Galaxies

To gather information concerning the structure of A1060 in a more homogeneous and objective manner than possible by purely visual means, the COSMOS machine was assigned to measure the central  $\sim 22$  degree<sup>2</sup> of three short-exposure (B, V and R waveband) UKST plates on which were recorded the cluster and its surroundings. The particulars of the V and R plates are included in Table 5.1. Operating in its



Coarse Mode (CM), COSMOS measured the positions, extents, minimum transmissions (measures of maximum surface brightness), areas, and quadrants of the orientation of tens of thousands of images that registered on each plate, and transferred the information to magnetic tape for off-line access by computer.

Rapid separation of the parameter sets which refer to galaxies from those that refer to foreground "field" stars is usually attempted at this stage by the application of quantitative classification criteria which are optimised for each plate in question. Until recently such easily measurable gross parameters as image shape (Tyson & Jarvis 1979) or surface brightness (e.g. Oemler 1974) have been the only practicable discriminators of image type; see Godwin & Peach (1977) for some comments on the efficacy of these criteria. MacGillivray et al. (1976a) have shown that the most successful discriminator for the sorting of COSMOS CM data of individual plates is a plot of the logarithm of the image area against the minimum transmission in each image; galaxies and stars segregate by surface brightness, the demarcation varying with image size and frequently with position on the plate. Such variations can sometimes be adjusted for, but this was not the case for the plates of A1060. As a result the application of classification criteria determined for sample areas on each plate failed to sort galaxies from stars over the full areas of the scans with any reliability. Additional factors contributing to the low (~40%) success rate were the comparatively high proportion of blended star images recorded at this galactic latitude, some of which mimic the properties of galaxies, and the



relative granularity of the emulsions involved (Table 5.1). Following the precedents of Godwin & Peach (1977) and Dressler (1978b), the elimination of the stars was instead accomplished by visual inspection.

The dataset pertaining to plate R1347 was selected for sorting, as a compromise between the desirability of a faint limiting magnitude to detect low surface brightness galaxies on one hand, and conversely of a bright limit to minimize the generation of spurious shredded images within the haloes of bright objects, and which presented a manageable exercise in manual data handling, on the other. COSMOS detected 158000 images during measurement of this plate, 55000 of which were less than ten units ( $1 \text{ unit} = 64 \mu\text{m}^2$ ) in area and attributable mainly to grain clumping in the emulsion. The remainder were treated as valid images and displayed in the form of computer listings and finding charts.

Scott (1962) notes that the probability that a galaxy is identified and counted depends on, among other factors, the state of the observer at the time of counting. Sorting proceeded at the rate of roughly one square degree per thirteen hours, so to mitigate the effect of possible day-to-day variation of the identification criteria, the 49 sub-regions of the scan area were inspected in semi-random order, though tending from centre outwards. Coverage of different quarters of the scan was provided by four different IIIa-J plates; that representing the south-west was affected by atmospheric haze, whereas the plate of the north-east quarter had a denser sky background exposure and the images on it were slightly trailed. Offsetting

this patchwork inhomogeneity, however, was the gain of some  $3^m$  in limiting magnitude and the high definition afforded by the IIIa-J plates, as compared to the 098 emulsion scanned by machine.

To check the repeatability of the selection of galaxies, 3 of the 49 sub-regions were re-examined after a lapse of some months, with the comparison plates rotated through  $90^\circ$  relative to their original orientation. The original and secondary classifications of the images in each area are compared in Table 2.2. Whilst fewer galaxies and more objects of uncertain nature were recorded during the repeat sessions, the actual numbers of galaxies recognised are stable to within  $\sim 10\%$ . This implies a consistency of 99% in the visual sorting process, a success rate infrequently achieved by automatic image sorters except under favourable conditions. The automatic detection process, on the other hand, has a repeatability of much less than 99% for faint images, and moreover different observers are liable to apply different visual criteria to any selection task.

At the conclusion of this process, CM parameters for 5185 individual galaxies and 227 instances of image blends had been placed on a master card file. 54 additional images of galaxies were found to have been split into double or multiple components; area-weighted centroids were assigned to the low surface brightness systems in this category, and the locations of the remainder were estimated either by direct inspection of the plates, or by transformation of the coordinates listed during the course of independent photographic photometry (Chapter 5). A total of 220 images,

distributed widely throughout the scan area, were of uncertain classification and were excluded from further analysis.

Following the elimination of the foreground star component, visual inspection of plates was resumed in order to classify as many galaxies as possible within the framework of a simplified de Vaucouleurs (= revised 3-dimensional Hubble) system. An attractive feature of this scheme (see e.g., de Vaucouleurs 1959) is its additive property, whereby ever higher degrees of refinement of classification can be accommodated according to the suitability of the plate material. A simple numerical coding scheme was devised which took advantage of this fact, based on the system employed in the RC1 (de Vaucouleurs & de Vaucouleurs 1964), to facilitate later computer analysis of the morphological statistics. Those parameters which describe stage and DDO luminosity class were ignored, since only the brightest systems were photographed with sufficient resolution.

The microcmetric axis dimensions of all galaxies with measurable diameters ( $\geq 1\frac{1}{2}$  graticule divisions or  $\geq 10$  arcsec) were also recorded during classification. More images were measured than could be classified; it was found to be particularly difficult to assign morphological types to overexposed disk systems at high inclinations to the line of sight, whereas the more lightly exposed non-IIIa-J emulsions were too coarse-grained to be of assistance. The preponderance of spiral galaxies in the classified sample is partly a reflection of the fact that they are the easiest type to identify out to moderate distances.

The projected distribution of the different types of galaxies is considered in Section 2.6.

#### 2.4 The Surface Distribution of Galaxies

It is convenient to discuss the distribution of the  $\sim 5500$  galaxies recognised on plate R1347, in terms of apparent photographic magnitude. The method whereby magnitudes are estimated from CM data is described later in Section 5.6.

Figure 2.3 shows the surface distribution of galaxies which have V-band magnitudes at the 25th mag arcsec<sup>-2</sup> isophote ( $V_{25}$ ), or approximations thereof (as explained in Chapter 5), brighter than 17<sup>m</sup>. This limit is chosen because at fainter magnitudes, the sample becomes significantly incomplete. Apart from the omission of small strips along the north and west boundaries, Figures 2.3 - 2.5 show approximately the full area of the CM scan. Note the high degree of symmetry in the central region of the cluster and the general decrease of surface density towards the edges of the field (Figure 2.3). This is a quantitative statement of what the charts in Appendix I make clear more qualitatively.

The indistinctness of the group of galaxies Klemola 13 (Klemola 1969) in Figure 2.3 deserves comment. This group is visible south preceding the core of Hydra I at an angular separation of only 0.8 Mpc, and is dominated by a large barred spiral and includes a pair of interacting galaxies. In the second finding chart in Appendix I, the group may be found at coordinates  $X = 14.9$  cm,  $Y = 8.5$  cm. Klemola 13 is also listed as the 378th entry in the catalogue of

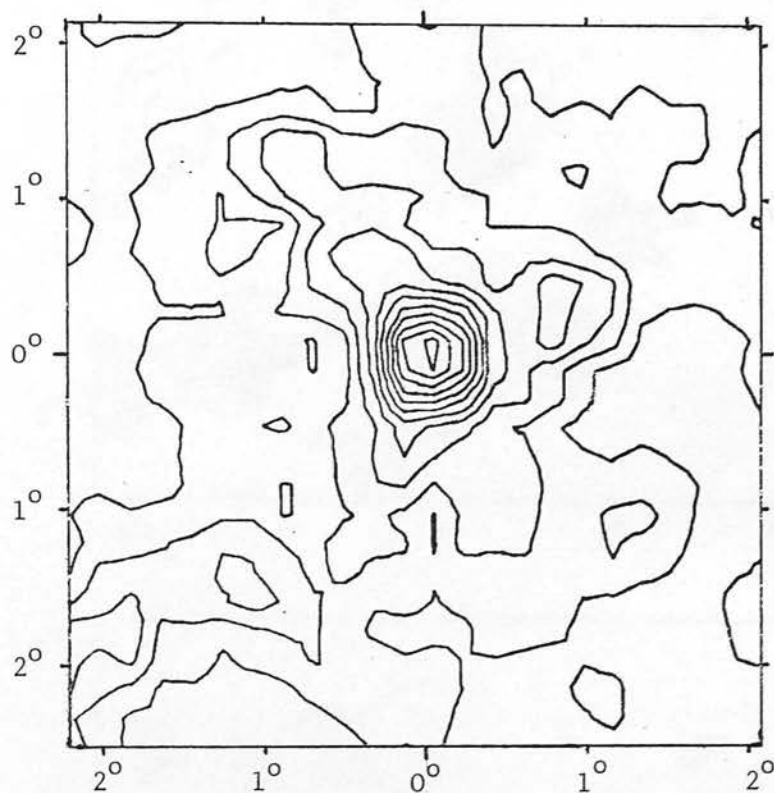


Figure 2.3

Isopleths of galaxies brighter than  $V_{25} = 17$  in a region of dimensions  $4.51 \times 4.85$  approximately. Contours are in steps of 20 galaxies degree<sup>-2</sup> from 20 to 220 galaxies degree<sup>-2</sup>.

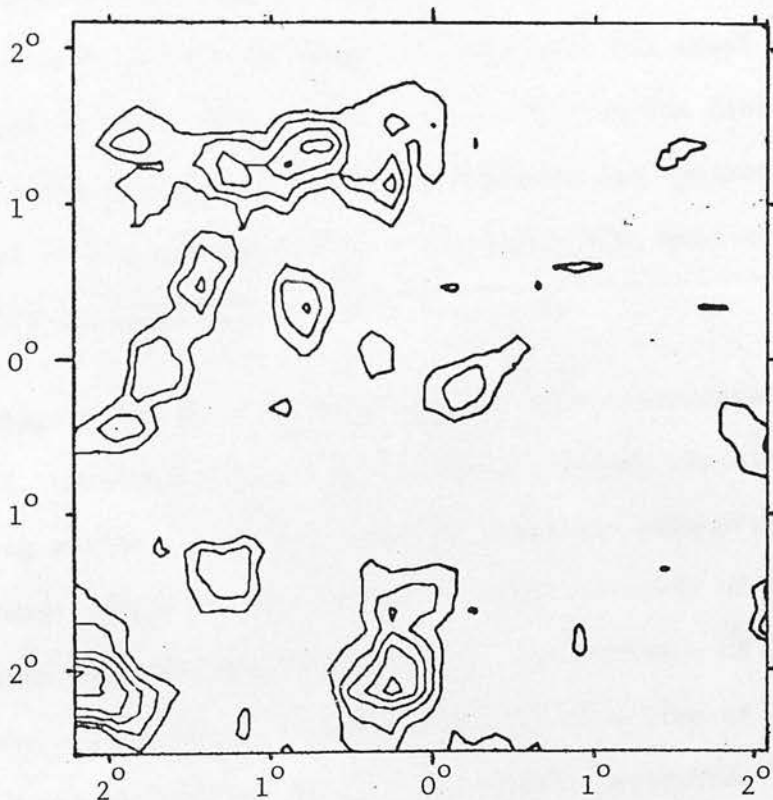
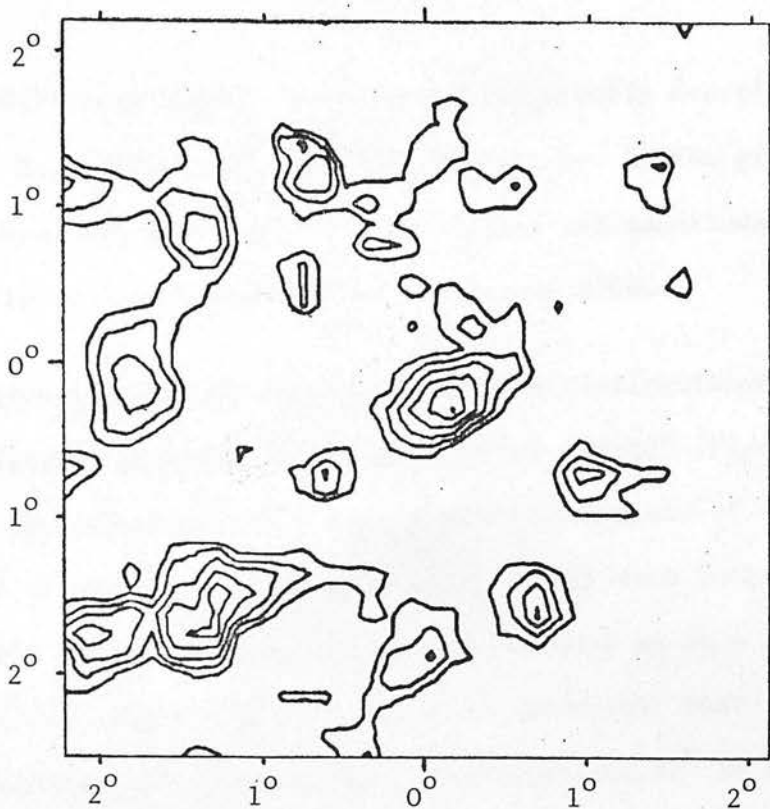


Figure 2.4 (top) Isopleths of 1387 galaxies of  $17.0 < V_{25} \leq 18.2$  contoured in steps of 20 galaxies degree<sup>-2</sup> from 100 to 200 galaxies degree<sup>-2</sup>.

Figure 2.5 (bottom) Isopleths of 3257 galaxies fainter than  $V_{25} = 18.2$ , contoured in steps of 40 galaxies degree<sup>-2</sup> from 200 to 360 galaxies degree<sup>-2</sup>.



Braid & MacGillivray (1978), where it is incorrectly described as very rich. No redshifts of the brighter galaxies in the group have been published, but their apparent dimensions and magnitudes suggest that they lie at a distance similar to that of A1060.

At fainter magnitudes Hydra I cannot be distinguished relative to the density of galaxies in the surrounding "field" (Figure 2.4). To a first approximation the number density of members of the cluster and of the field (both considered later) each increase exponentially with magnitude, but the latter does so at a much faster rate (cf. Figures 6.3 & 6.4), with the result that field galaxies dominate the total population even before  $17^m$  is reached. The most interesting feature in Figure 2.4 is a flattened cloud of galaxies, central in the diagram, which is only  $0.2^\circ$  south preceding the position of A1060 and almost certainly behind the cluster. In Figures 2.4 and 2.5 other density enhancements are apparent mainly to the east of the central meridian through A1060, most of which probably lie at large line-of-sight distances.

At magnitudes still fainter ( $\geq 20^m$ ) than illustrated in Figure 2.5, extremely distant rich clusters become visible. Although very strong subjective effects tend to frustrate attempts to select a homogeneous sample of these faint clusters, a study of their distribution suggests two observations: the presence of greater optical obscuration in the (south-westerly) direction of the galactic plane, and the absence of strong intracluster absorption in A1060 (other than what is attributable to the individual bright cluster members). There is no indication that the number of galaxy

associations visible in the background are fewer along the line of sight through this cluster than outside it. Zwicky (e.g. 1957) has reported that galaxy clusters appear to be fewer in number behind intervening clusters than in other directions, but Abadi & Edmunds (1978) consider that this effect is due merely to problems of recognition.

To locate the position of the centre of the cluster the projected distributions of galaxies brighter than  $V_{25} = 15, 16$  and  $17^m$  were each smoothed and contoured in the manner of Figure 2.3, and the centres of symmetry of the isopleths estimated. This is essentially a refined version of the conventional procedure of defining the cluster centre by means of strip counts. The position of maximum number density in A1060 is at  $X = 126.8 \pm 1.4$  mm,  $Y = 142.8 \pm 1.0$  mm in the system of the COSMOS Coarse Mode measurements, and corresponds to the 1950.0 coordinates  $\alpha = 10^h 34^m 17^s \pm 6s$ ,  $\delta = -27^\circ 13' \pm 1'$ . The two brightest galaxies in the cluster, N3309 and N3311, are located just 30 kpc immediately south of this position (cf. Figure 6.5). The centre of (luminous) mass in the cluster probably coincides with this giant binary.

There is evidence (e.g. MacGillivray et al. 1976b) that most or all distant, rich clusters of galaxies are elongated systems. The degree of apparent flattening of A1060 may be estimated by the iterative fitting of a series of ellipses to the surface density distribution illustrated in Figure 2.3. Stable (or mildly oscillating) solutions were obtained in which each ellipse assumed a shape and orientation most consistent with the distribution of galaxies that

it enclosed. The mean axial ratio of the best-fitting models was  $0.87 \pm 0.02$ , with no marked dependence on the proportion of the sample of galaxies bounded (that is, on the size of the ellipse). If it is maintained that A1060 has a disk-like structure, the disk must be aligned so as to present an almost face-on aspect.

The high degree of circular symmetry makes A1060 amenable to representation by the isothermal sphere model (cf. Zwicky 1942, 1957; Bahcall 1972, 1973a,b,c, 1974, 1975; Austin & Peach 1974; MacGillivray et al. 1976b). The density distribution in an infinite isothermal "gas" sphere of self-gravitating point masses (Emden 1907) as projected onto a plane has been tabulated by Zwicky (1957). He uses rather large increments in the radial variable to describe the profile to very great distances from the nucleus, but the central peak of the profile is not well defined, and some interpolation of his figures was necessary. A  $\chi^2$ -minimising computer procedure was applied to fit the observed surface distribution of galaxies brighter than  $V_{25} = 17$  to this projected Emden sphere (Figure 2.6).

The structural index of A1060 is  $\beta = 2.1$  minutes of arc, corresponding to a core radius (Bahcall 1973a) of  $R_c = 3\beta = 73$  kpc. If the innermost core of the observed distribution is excluded from the fit, the derived  $\beta$  increases by 20% to 2.5. Butcher & Oemler (1978) have carried out visual counts of galaxies in A1060 to an estimated limiting magnitude of  $V \approx 16.6$ ; a fit of their data to the standard projected sphere yields  $\beta = 2.6$ , which is in good agreement. The matching procedure is not particularly sensitive to the choice of cluster centre, but estimates of  $\beta$  may vary with the

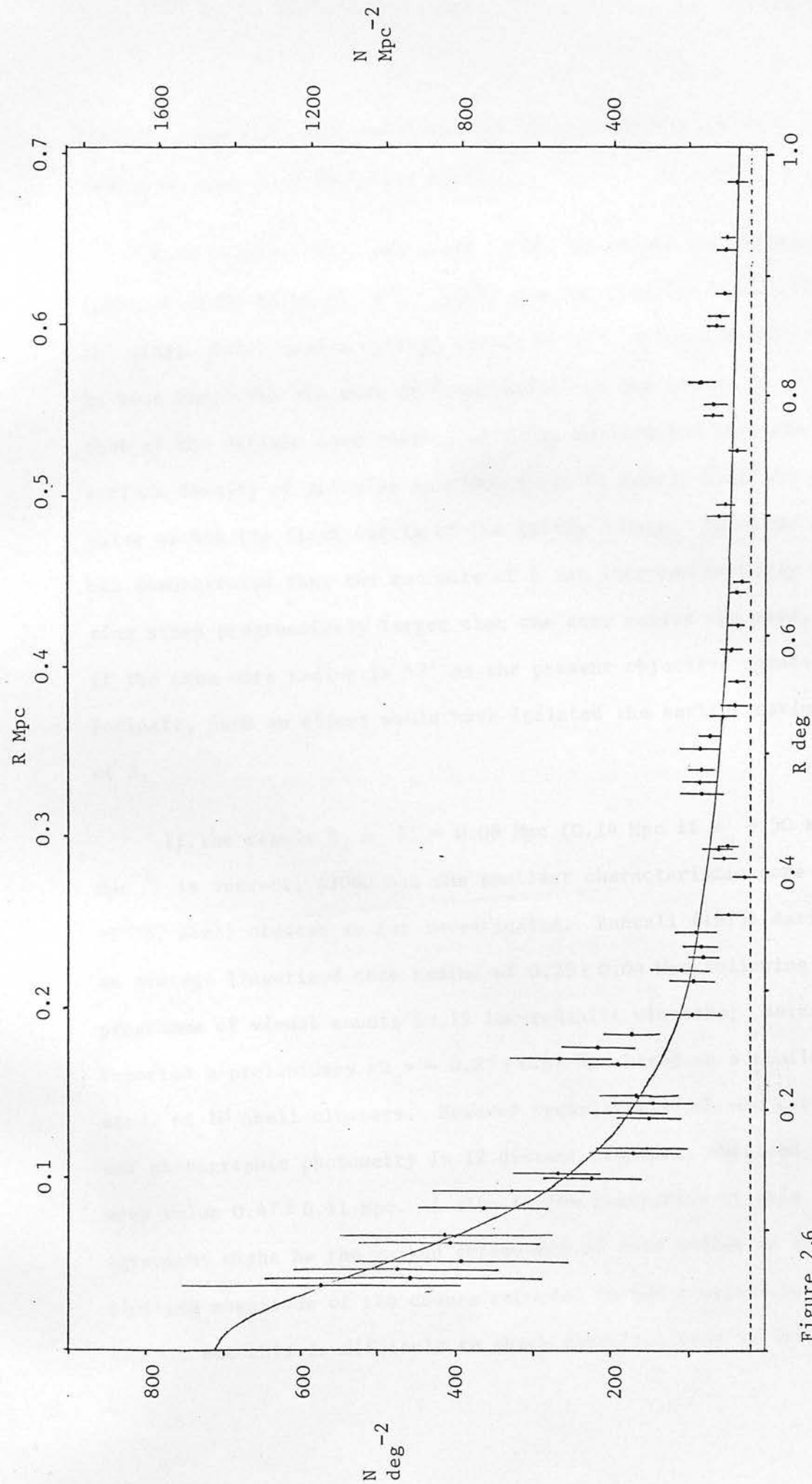


Figure 2.6

The radial surface density distribution of galaxies of  $V_{25} < 17$  from counts in annuli of various widths, and the best isothermal sphere representation. The adopted background level is shown by the dotted line.

limiting magnitude (and completeness) of the counts, and with the choice of ring size (Quintana 1979).

Both Zwicky (1942) and Kwast (1966) estimated the structural index of A1060 to be  $\beta = 4'$ . Zwicky counted galaxies within rings 10' wide; Kwast used a central circle of 15' followed by 10' rings. In both cases the bin size or 'resolution' of the counts is about that of the derived core radius, with the implication that the surface density of galaxies in A1060 falls to nearly half its central value within the first circle of the Zwicky counts. Quintana (1979) has demonstrated that the estimate of  $\beta$  can increase markedly if ring sizes progressively larger than the core radius are used, and if the true core radius is  $\sim 7'$  as the present objective counts indicate, such an effect would have inflated the earlier estimates of  $\beta$ .

If the result  $R_c \approx 7' = 0.08 \text{ Mpc}$  ( $0.14 \text{ Mpc}$  if  $H_0 = 50 \text{ km s}^{-1} \text{ Mpc}^{-1}$ ) is correct, A1060 has the smallest characteristic core size of any Abell cluster so far investigated. Bahcall (1975) derived an average linearised core radius of  $0.25 \pm 0.04 \text{ Mpc}$  following a programme of visual counts in 15 low-redshift clusters; Baier (1976) reported a preliminary  $\langle R_c \rangle = 0.23 \pm 0.08 \text{ Mpc}$  based on a similar study of 10 Abell clusters. However Dressler (1978c), who carried out photographic photometry in 12 distant clusters, obtained the mean value  $0.47 \pm 0.11 \text{ Mpc}$ . A clue to the resolution of this disagreement might be the marked dependence of core radius on the limiting magnitude of the counts recorded in two clusters by Quintana (1979), but this is difficult to check quantitatively in the present



case because very large statistical uncertainties in the density of galaxies in the core arise if the modest numbers of galaxies in A1060 are subdivided by magnitude range.

## 2.5 The Correction for Non-members

The redshifts of all but a very small number of galaxies in the sample are unavailable, so any correction intended to eliminate galaxies which are not members of A1060 must inevitably be statistical in nature. If the distribution of "field" (i.e. non-member) galaxies of all types within the volume of space surveyed is homogeneous and isotropic, then the number,  $N(\leq m)$ , of field galaxies brighter than apparent magnitude  $m$  varies as  $N(\leq m) \propto 10^{0.6m}$ . The constant of proportionality may be fixed by the requirement that the expression should agree with the number of field galaxies to a given limiting magnitude in the regions immediately beyond the boundary of the cluster.

The surface density of galaxies brighter than  $V_{25} = 17$  in the scan area decreases from approximately 31 galaxies degree<sup>-2</sup> at 0.91 Mpc from the cluster centre to 20 galaxies degree<sup>-2</sup> at 1.56 Mpc. A similar decline of number density with increasing radial distance has been found for galaxies in the magnitude range  $17 < V_{25} < 18$ . Fainter than  $V_{25} = 18$ , the number distribution fluctuates strongly due to the presence of background clusters but shows no corresponding decrease from centre to edge. It is clear from these observations and from Figure 2.2 that there are still appreciable numbers of galaxies associated with A1060 at 1 to 1.5 Mpc from the centre, and



that to determine the field density, either photometrically calibrated number counts that refer to a wider local region must be considered (cf. Dressler 1978b,c) or some form of extrapolation to larger angular distances must be applied using the existing counts (cf. Butcher & Oemler 1978; Quintana 1979).

Neither the CGCG (Zwicky et al. 1961-68) nor the Lick Observatory galaxy counts (e.g. Shane & Wirtanen 1954) enumerate the distribution of galaxies in this region, but earlier estimates of the local field density exist. Shapley (1933) determined an average background density of 34.6 galaxies degree<sup>-2</sup> to a limiting photographic magnitude,  $m_{pg}$ , of 17.3 for galaxies, over a nine square degree region excluding the cluster core. As this would correspond roughly to the annulus about Hydra I with inner and outer radii of 0.5 Mpc and 1 Mpc respectively, his figure is an overestimate. Zwicky (1941) found 2.8 galaxies degree<sup>-2</sup> brighter than  $m_{pg} \approx 16$  at a projected distance of  $5^{\circ}$  (3.5 Mpc). If it is assumed that both  $m_{pg}$  systems approximate the International photographic scale, then (Brown 1974)

$$B \approx m_{pg} + 0.1 + 0.14(m_{pg} - 15)$$

where  $B$  is a blue-band magnitude presumably equivalent to the standard face-on magnitude  $B(0)$  defined in the RCl (de Vaucouleurs & de Vaucouleurs 1964), which correspond to the  $\sim 23.5$ th  $B$  mag arcsec<sup>-2</sup> isophote in ellipticals and lenticulars and  $\sim 24.5$ th  $B$  isophote for later types (Freeman 1976). With  $B_{25} \approx B$  and a mean colour index  $\langle B_{25}^c - V_{25}^c \rangle = 1.1$  for lenticular galaxies at  $V_{25} = 17$ , it follows that

Shapley's counts imply  $N(V_{25} \leq 17) \approx 58$  galaxies degree<sup>-2</sup> in this region, whereas Zwicky's imply  $\sim 37$  galaxies degree<sup>-2</sup>. These figures are very uncertain, not least because the original estimates of limiting magnitude were crude. By contrast, Butcher & Oemler (1978) estimate that there are only 7 non-members per square degree brighter than magnitude 16.6 on V-band plates.

A fit of all galaxies of  $V_{25} < 17$  to the Emden sphere model gives a value of 17.6 galaxies degree<sup>-2</sup> at a projected radius of  $5^{\circ}$ . With proper regard of the above considerations, an estimate of the background which exceeds this value slightly but which is less than the actual counts at a projected radius of 1.5 Mpc ( $2.1^{\circ}$ ) is adopted, viz.  $19 \pm 6$ : galaxies degree<sup>-2</sup> to  $V_{25} = 17$ , which corresponds to the logarithmic cumulative luminosity function

$$\log N(\leq V_{25}) = 0.6V_{25} - 8.921$$

for field galaxies in each square degree, as dimmed by extinction. For comparison, the function adopted by Sandage & Tammann (1976) [ $\log N(\leq m) = 0.6m - 9.30$ ] predicts 7.2 galaxies degree<sup>-2</sup> to this limit, after allowance for extinction ( $A_V = 0.13$ , Sandage 1973a) and for their use of total magnitudes ( $\approx V_{25} - 0.2^m$ ). The field correction of Bucknell et al. (1978) and previous papers in that series imply  $N(\leq 17) = 19.4$  galaxies degree<sup>-2</sup>, very close to the figure assumed here.

The adopted field correction predicts that at most, 276 of the 526 galaxies brighter than  $V_{25} = 17$  and within 1.5 Mpc of the centre of A1060 are likely to be non-members. Since this subtracted

'background' tends to exceed the density of galaxies observed at the corners of the  $\sim 22$  degree<sup>2</sup> scan area, 250 is probably a minimum estimate of the number of systems associated with A1060. The relative contribution of this background to the observed number density of galaxies at  $V_{25} = 17$  is illustrated in Figure 2.6. Knowledge of the field correction at each magnitude is essential to the later discussion, in Section 6.2, of the cluster luminosity function.

## 2.6 The Radial Variation of the Cluster Composition

The surface distributions of the galaxies that are brighter than  $V_{25} = 17$  (a few are fainter systems blended with a star image) and which belong to different morphological types is displayed in Figures 2.7 to 2.9. Some galaxies which were not categorised on visual inspection and which happened to lie within the smaller scan area of the primary photographic photometry were identified as late-type systems on the basis of integrated B-V colour (Section 6.1). These were assigned to the spiral and irregular galaxy category, and marked by crosses in Figure 2.9 to distinguish them from visually classified systems. The distribution of the remaining unclassified galaxies brighter than  $17^m$  is illustrated in Figure 2.10.

The differences of the distributions of the three recognised types of systems is striking. Elliptical galaxies form a rather compact central core superposed on an apparently distinct, homogeneous field distribution. The lenticulars do not exhibit this dichotomy, but instead are smoothly concentrated toward the centre of the cluster, and trail off towards the outskirts with no

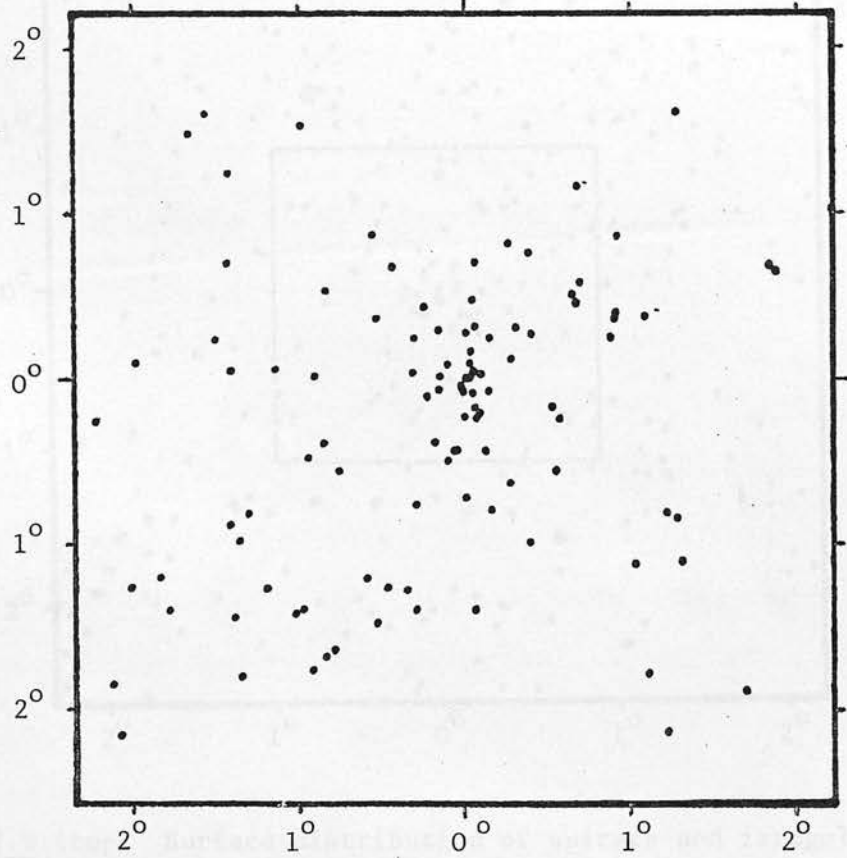
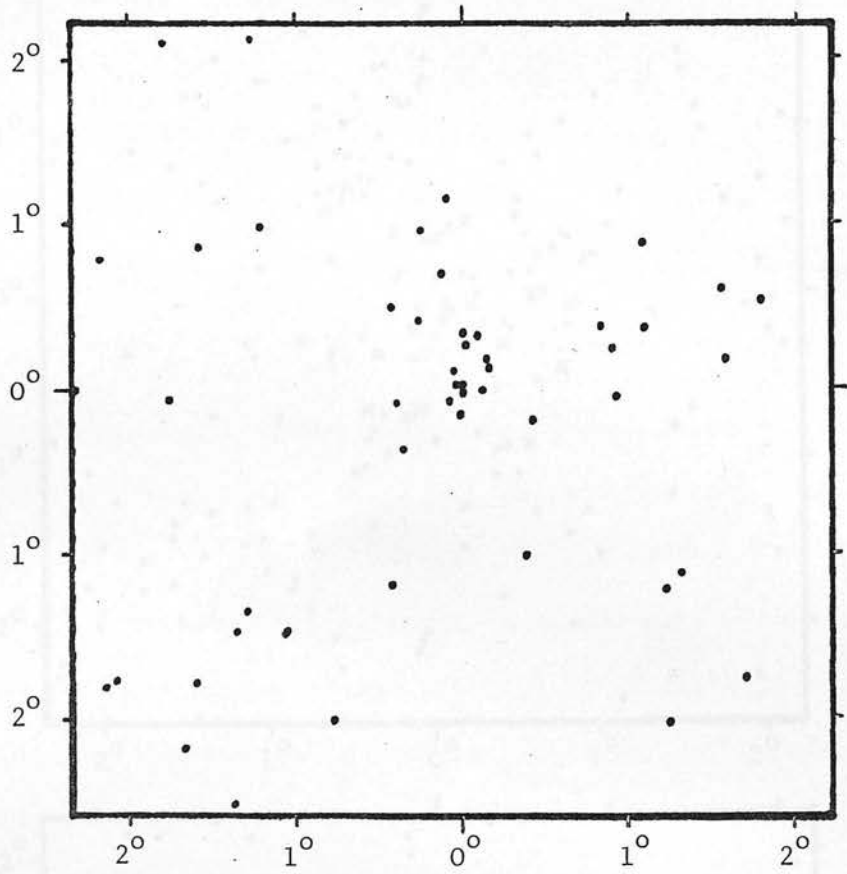


Figure 2.7 (top) Surface distribution of ellipticals,  $V_{25} < 17$

Figure 2.8 (bottom) Surface distribution of lenticulars,  $V_{25} < 17$

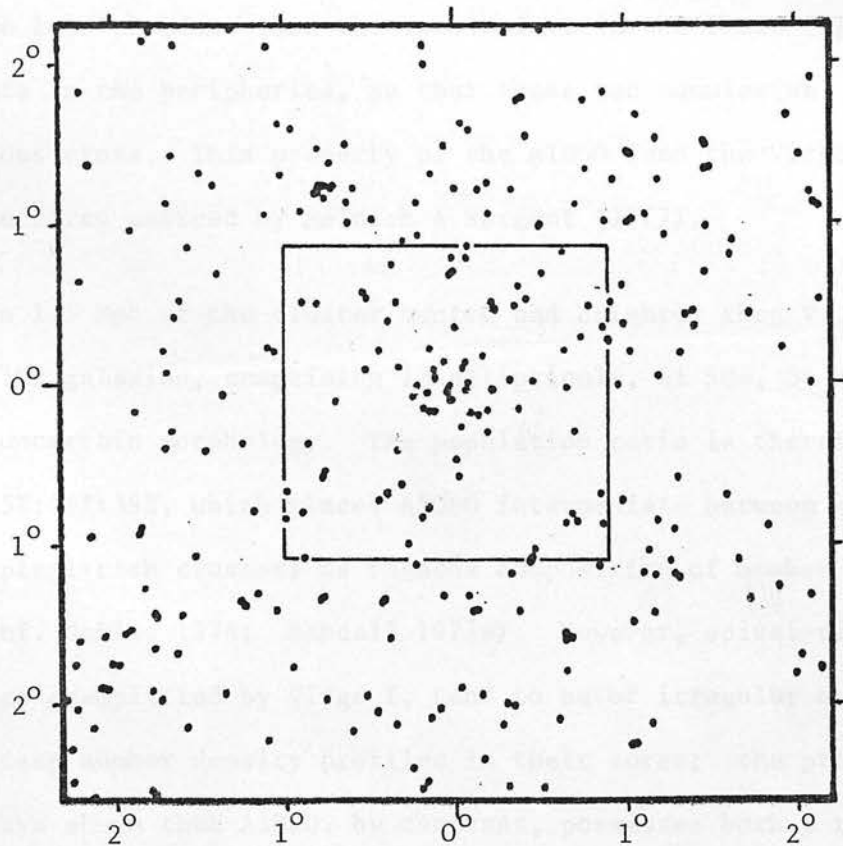
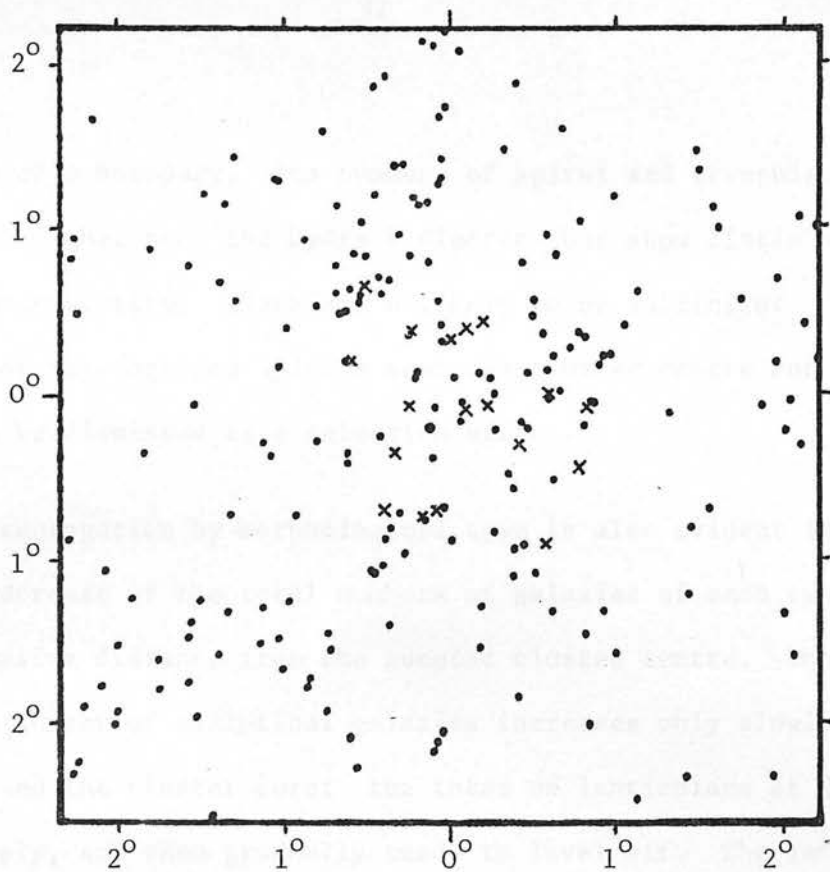


Figure 2.9 (top) Surface distribution of spirals and irregulars,  $V_{25} < 17$ .

Figure 2.10 (bottom) Surface distribution of unclassified galaxies,  $V'_{25} < 17$ .

suggestion of a boundary. The numbers of spiral and irregular systems are higher near the Hydra I cluster, but show little or no central concentration; there are unlikely to be sufficient instances of unrecognised spirals near the cluster centre for this paucity to be dismissed as a selection effect.

This segregation by morphological type is also evident in the rates of increase of the total numbers of galaxies of each type with increasing distance from the adopted cluster centre. The cumulative number of elliptical galaxies increases only slowly with radius beyond the cluster core; the total of lenticulars at first rises steeply, and then gradually tends to level off. The late-type systems are less abundant than the lenticulars in the inner regions, but dominate in the peripheries, so that these two cumulative distributions cross. This property of the A1060 (and the Virgo I) cluster was first noticed by Melnick & Sargent (1977).

Within 1.5 Mpc of the cluster centre and brighter than  $V_{25} = 15$  there are 102 galaxies, comprising 13 ellipticals, 41 S0s, 34 spirals and 14 of uncertain morphology. The population ratio is therefore E:S0:S = 15%:46%:39%, which places A1060 intermediate between spiral-poor and spiral-rich clusters as regards composition of member galaxies (cf. Oemler 1974; Bahcall 1977a). However, spiral-rich clusters, as exemplified by Virgo I, tend to be of irregular outline and lack steep number density profiles in their cores; the preceding sections have shown that A1060, by contrast, possesses both a regular structure and strong central condensation. The Hydra I cluster may therefore be classified as of intermediate type in the sense of Bahcall (1977a).



## 2.7 Geometric Parameters of the Galaxies

Supplementary information concerning the visual geometry of the brightest galaxies was collected on a manually controlled plate comparator. In this machine image acquisition and display are facilitated by a television system, and a teletype provides positional readout in steps of 12.5 microns over some 20 degree<sup>2</sup> of plate. Scanning in a boustrophedonic fashion, an initial finding list of over 500 bright galaxies was prepared. The axial diameters,  $D$  and  $d$ , and the orientations of the projected major axes were then determined for many of these images by the following procedure. The coordinates of each end of the major axis were noted, and next those of the minor axis: this supplied four points which defined an approximating ellipse, and the procedure was repeated. The ratio of the average lengths of the two semi-major and semi-minor axes defined the apparent ellipticity of the image. The projected orientation was taken to be the average of the mean major axis and mean minor axis orientations, the latter shifted into correspondence by the addition of  $90^\circ$ , weighted by axis length; the meridian passing through the cluster centre defined the axis of reference, angles increasing north through east. Arigo et al. (1978) have described a quadruple averaging procedure for the reduction of similar data, and their prescription was found to yield comparable figures for axis ratio and orientation.

For direct comparison with the above, a second set of values was derived with the plate rotated through  $112^\circ$ . This angle was chosen following Thompson (1976), to avoid unconscious biases which

can occur at multiples of  $45^\circ$ . Not all galaxies in peripheral regions could be measured twice, because the accessible region of the plate varied with alignment. The internal errors in the data may be quantified by evaluating (cf. Hawley & Peebles 1975)

$$\Delta P = \frac{1}{2} \langle P_1 - P_2 \rangle \quad \text{and} \quad \sigma(P) = \frac{1}{2} \langle (P_1 - P_2)^2 \rangle^{\frac{1}{2}}$$

where the angular brackets  $\langle \rangle$  denote average value and the subscripts 1 and 2 refer to the first and second determinations of parameter  $P$  respectively. For the axis ratios  $R = d/D (\leq 1)$  and orientations  $\phi$  determined for 167 galaxies at both plate alignments,

$$\Delta R = -0.01 \qquad \sigma(R) = 0.05$$

$$\Delta \phi = -0.81^\circ \qquad \sigma(\phi) = 6.80^\circ$$

These indicate no important systematic disagreement of the two sets of measured shapes and orientations. The dispersion about the mean orientation,  $\sigma(\phi)$ , falls to  $\sim 4^\circ$  if 33 images of apparent ellipticity  $e = 10(1-R)$  less than 2.5 are excluded.

When the axis ratios of non-spherical images with gradual luminosity gradients are measured by eye rather than by machine, their ellipticities tend to be overestimated (Holmberg 1946). To determine whether, and how seriously, the axis ratios derived from the above measurements with a manual plate comparator ( $R_3$ ) and from micrometric axis diameters collated during galaxy classification ( $R_2$ ) were biased, each set of apparent axis ratios was compared with the values ( $R_1$ ) obtained from direct measurements of

the extents of the 25th V mag arcsec<sup>-2</sup> isophotes of bright galaxies on contour plots. The ratios  $R_1$  should be free from errors of the type described by Holmberg since the isophotal contours were constructed by machine from objective data. In analogy with the previous analysis, comparisons of 73 galaxies yielded

$$\begin{aligned} \langle R_2 - R_1 \rangle &= -0.11 & \langle (R_2 - R_1)^2 \rangle^{\frac{1}{2}} &= 0.16 \\ \langle R_3 - R_1 \rangle &= 0.00 & \langle (R_3 - R_1)^2 \rangle^{\frac{1}{2}} &= 0.15. \end{aligned}$$

These figures confirm that the micrometric measurements do suffer from a "Holmberg effect" which causes ellipticities to be overestimated by 1 unit on average. The data from the plate comparator appear to be unbiased, but because these measurements refer typically to contours of higher luminosity ( $\lesssim 23.5$  V mag arcsec<sup>-2</sup>), the comparison is not strictly valid; the brighter isophotes of the galaxies may be more circular. The accuracy of the estimates of  $R_1$  and  $R_2$  is about  $\pm 0.1$  for the brighter systems, deteriorating for smaller images.

The most extensive and hence useful body of axis ratio estimates is the set  $R_2$ . These were first corrected for the Holmberg effect by means of the expressions given in RC1:

$$\begin{aligned} R_c &= .86R_2 - .14(1 - R_2)^6 + .14 && \text{for ellipticals} \\ &= .81R_2 - .19(1 - R_2)^6 + .19 && \text{for other types,} \end{aligned}$$

with the results

$$\begin{aligned} \langle R_c - R_1 \rangle &= -0.02 & \langle (R_c - R_1)^2 \rangle^{\frac{1}{2}} &= 0.11 \\ &&& (74 \text{ systems}), \end{aligned}$$

which are clear improvements over the unadjusted quantities and fully justify the correction procedure.

The frequency distribution of ellipticities corresponding to the corrected axis ratios  $R_c$  for galaxies within 1.5 Mpc of the cluster centre is shown for each of three ranges of  $V_{25}$  magnitude in Figure 2.11. Apart from N3314, which consists of two spirals superposed in the line of sight, the sample representing the galaxies brighter than  $V_{25} = 15$  is complete (upper histogram, Figure 2.11); fainter than this, omissions become progressively more numerous and contamination by non-members of A1060 more severe. The data for galaxies which satisfy  $V_{25} < 16$  may usefully be compared with the frequency distribution, shown in Figure 4(a) of MacGillivray & Dodd (1979), of the apparent ellipticities of galaxies in the field of the rich cluster at  $\alpha_{1950} = 00^{\text{h}} 47^{\text{m}}$ ,  $\delta_{1950} = -29^{\circ} 46'$ . That field is a square of side  $\sim 3.2$  Mpc at the distance of the cluster, of similar size to the 3 Mpc-diameter circular boundary applied here to A1060, if distances are scaled simply as the redshifts. A greater proportion of elongated ( $e > 3$ ) systems are found in A1060, about 50% as compared to 40% in the other cluster. This indicates that the population of Hydra I is relatively richer in disk systems.

Strom & Strom (1979a) have reported a relatively greater incidence of flattened elliptical (E) galaxies in the visibly elongated cD cluster A401 than in the less flattened cD cluster A2670, and suggest that as a general rule, the flattening of a rich cluster may correlate with the shape of the member galaxies. At a very different scale, the average ellipticity of poorly populated galaxy

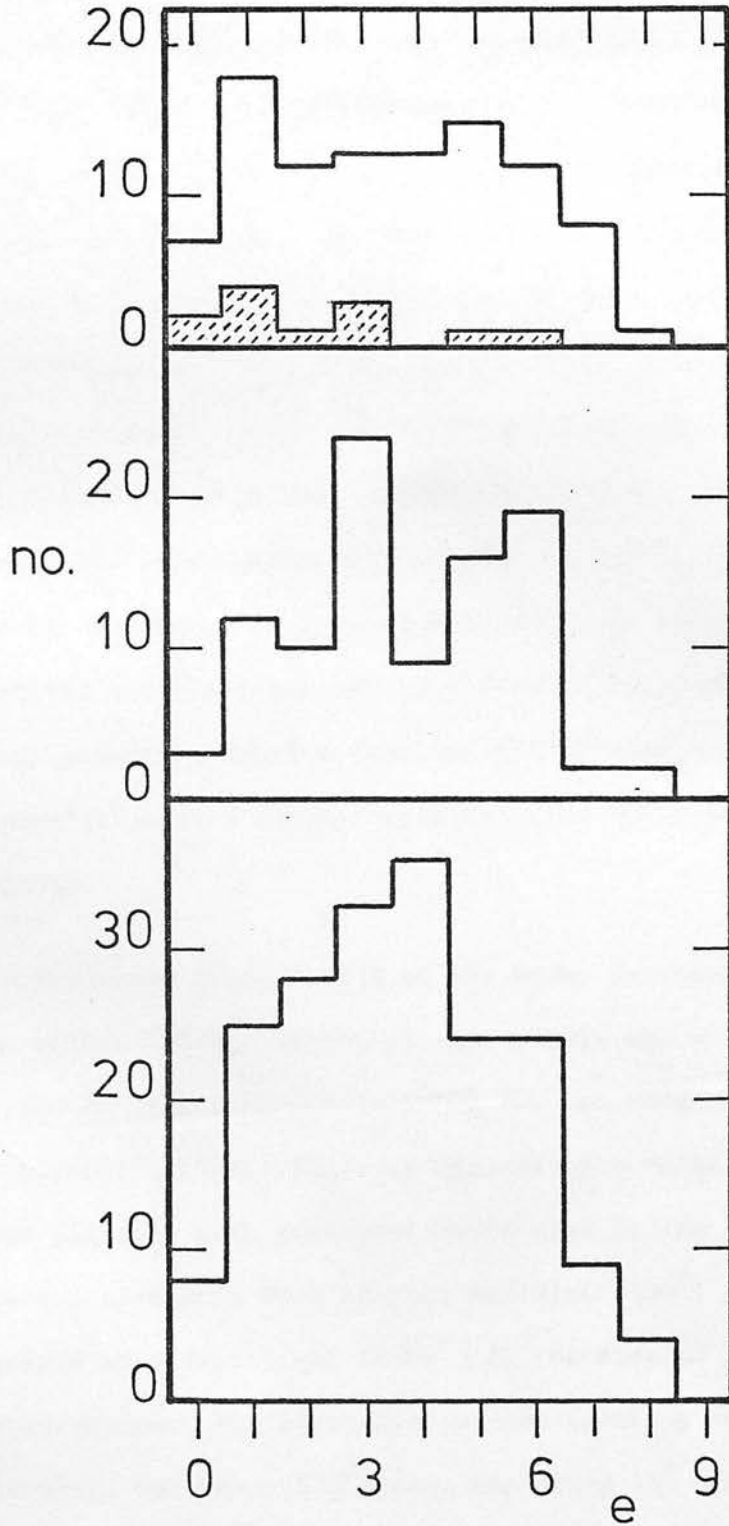


Figure 2.11

Frequency distributions of the ellipticities of those galaxies within 1.5 Mpc of the cluster centre and of magnitude (i)  $V_{25} < 15$  (top; semi-hatched histogram represents E galaxies); (ii)  $15 \leq V_{25} < 16$  (centre; completeness 92%); and (iii)  $16 \leq V_{25} < 17$  (bottom; completeness 56%).



groupings is comparable to the mean for individual elliptical galaxies (Rood 1979). Both observations may constitute evidence that there exists some physical process which governs both the distribution of stars in a galaxy and the distribution of galaxies in an association. The population of E galaxies in A1060 which are brighter than  $V_{25} = 15$ , however, is barely sufficient for a meaningful comparison to be made with the frequency distributions of ellipticities of E galaxies in other clusters: the few Es in this sample are represented by the semi-hatched histogram in Figure 2.11. Inclusion of fainter E galaxies is the obvious remedy, but redshifts would be required to eliminate non-members. The relationship between cluster type and the flattening of the E galaxy population is a further consideration which has yet to be investigated.

The frequency distribution of the known orientations,  $\phi$ , of galaxies within 1.5 Mpc of the cluster centre and with apparent axis ratios of less than 0.9 is shown for two ranges of  $V_{25}$  magnitude in Figure 2.12(a). There is apparently a relative deficiency of bright galaxies with projected major axes in the east-west plane (orientation eastwards from central meridian  $\approx 100^\circ$ ) and most systems are observed at orientations of  $60^\circ \pm 20^\circ$  or else at right-angles to this direction. The impression of non-isotropy does not hold up statistically, however. A  $\chi^2$  test, employing  $15^\circ$  bins of  $\phi$ , reveals the observed frequency histogram to be significantly non-random only at the 40% level, and on the modelled amplitude test (see Hawley & Peebles 1975, or Thompson 1976) the probability of a chance



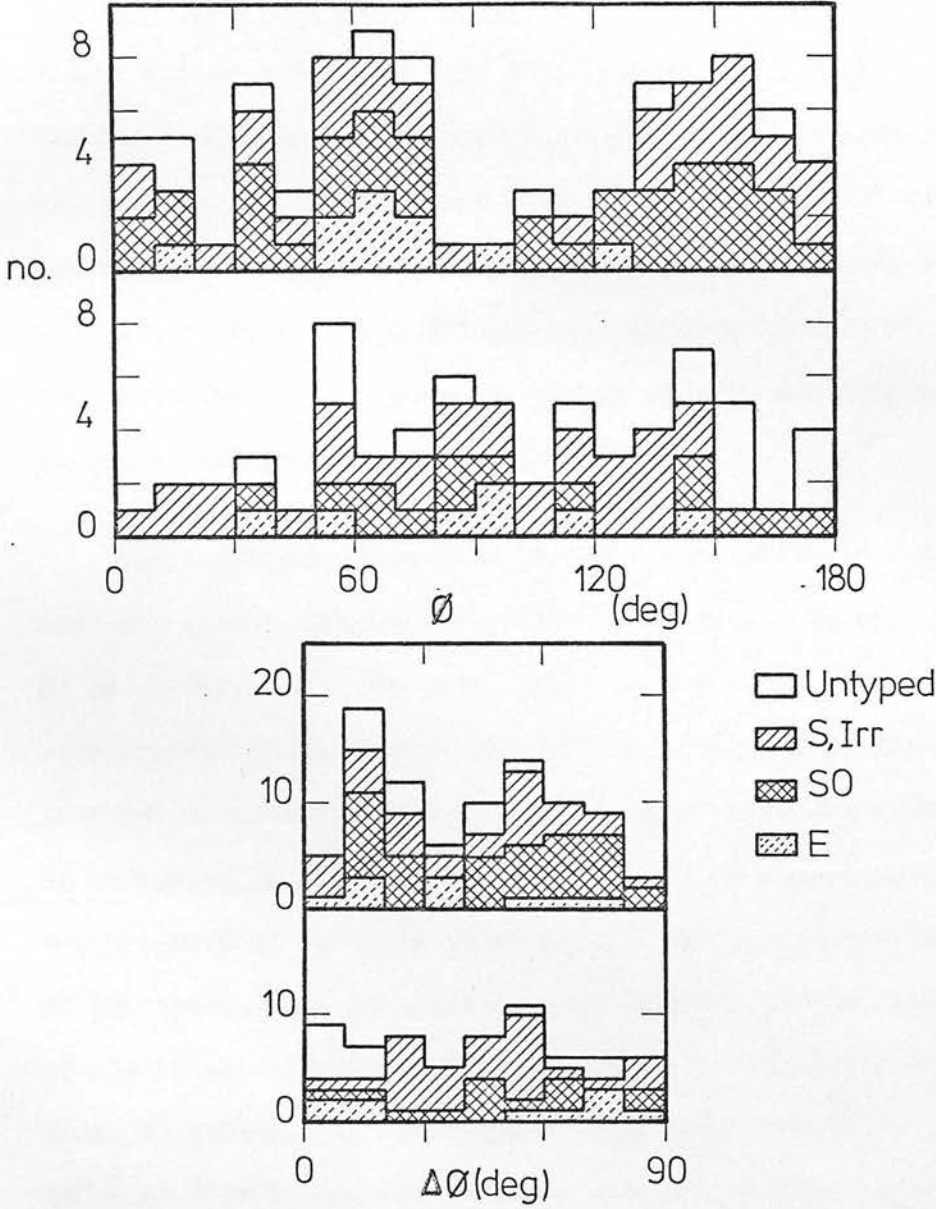


Figure 2.12 (a) (top)

Frequency distributions of orientations  $\phi$  of those galaxies within 1.5 Mpc of the cluster centre and of magnitude (i)  $V_{25} < 15$  (upper histogram; completeness 95%); and (ii)  $15 \leq V_{25} < 16$  (lower histogram; completeness 41%).

Figure 2.12 (b) (bottom)

Frequency distributions of the radial position angles  $\Delta\phi$  of the same samples (i) and (ii).

occurrence of this distribution is 66%, which is high. Nonetheless, it is worth mentioning that many galaxies in the cluster have approximately the same orientation as the brightest member, N3311 ( $\sim 50^\circ$ ). N3308 and N3309 (Figure 6.5) are conspicuous examples of near-parallel alignment, and most of the bright E galaxies satisfying  $V_{25} < 15$  have similar orientations. A second noteworthy point is the correspondence of the orientation of the cluster ( $\sim 121^\circ$ ) as determined by iterative ellipse fitting, with the axis defined by the central four bright NGC galaxies numbers 3308, 3309, 3311 and 3312 ( $\sim 130^\circ$ ) and also with the second peak in the frequency distribution of galaxy orientations ( $\sim 150^\circ$ ).

The frequency histograms of the radial position angles,  $\Delta\phi$ , of the same galaxy samples are plotted in Figure 2.12(b). The angle  $\Delta\phi$  is defined to be the orientation of the major axis of each galaxy with respect to the imaginary radius vector linking the centre of that galaxy to the centre of the cluster. Galaxies which appear to be elongated in the direction of the cluster centre have  $\Delta\phi = 0^\circ$  and are said to be radially aligned. The frequency distribution of  $\Delta\phi$  for the fainter galaxies sampled is emphatically random; that of the brighter systems is less readily interpreted, but the breakdown, 47 galaxies in the range  $0^\circ \leq \Delta\phi \leq 44^\circ$  and 42 in the range  $45^\circ \leq \Delta\phi \leq 89^\circ$ , does not indicate that these galaxies are preferentially aligned either radially or tangentially.

Before discussing these results, it should be pointed out that they are based on two-dimensional, projected representations of the true spatial orientations and shapes of the galaxies and of their

distribution. A good example of the peculiar distortions which may result from projection effects is provided by the pair of lenticular galaxies numbers 22 and 38 in the list of Appendix II. Number 38 is an unusual system intermediate in appearance with regard to N128 and N7332 (depicted in the Hubble Atlas of Sandage, 1961); it consists of a box-shaped nuclear region within an envelope that resembles a peanut with pointed ends. The immediate neighbour, number 22, is a normal S0 with a slightly elongated nucleus. It is evident from the appearance of the pair that they lie almost parallel to each other and nearly edge-on to the line of sight; yet, because their intrinsic major axes are foreshortened, they are assigned quite different values of  $\phi$ . Thus although statistically the gross features of the spatial alignments in a large assembly tend to be preserved in projection, inferences regarding isolated cases are unreliable.

MacGillivray & Dodd (1979) have ventured that systematic radial alignment trends are more likely to be observed in clusters classified as early in the scheme of Bautz & Morgan (1970), viz. classes I, I-II and II, whereas a tendency of member galaxies to align with the apparent major axis of the cluster as a whole may be more common amongst clusters of later BM type. A special case of the latter effect may arise in clusters which are dominated by a highly luminous giant galaxy; this first-ranked member is often aligned with the apparent major axis defined by the positions of the fainter galaxies (Sastry 1968; Dressler 1978b; Green 1978). Thompson (1976) suggests that linear alignment of a cluster and its membership is more probable amongst highly elongated clusters in general, or in

types L (linear) and F (flattened) in the classification system of Rood & Sastry (1971), in particular.

The trends observed in A1060, even though they are not statistically impressive, provide a modicum of support for the above proposals. The exception is the special case: whilst a number of galaxies in excess of average are oriented in parallel with the apparent major axis of the first-ranked system, this axis does not coincide with the direction of elongation of A1060. The cluster is classified as of BM type III, and the bright galaxies in this area do show a mild preference to be aligned with the distributive elongation axis; at the same time, some of them are aligned roughly parallel to the direction defined by the central four NGC galaxies, as appropriate for a cluster of RS type L. Actually the classification of Hydra I is C on the RS system, but the direction of slight flattening appears to be well defined, and is stable even if the brightest galaxies are excluded.

Table 2.1

Visual counts of "prominent" galaxies within circular annuli of width 10 arcmin centred on A1060, broken down by galaxy type.

'Ring 1' is a circle of radius 10 arcmin.

E = elliptical galaxies, SO = lenticulars, S = spirals and irregulars, and U = untyped.

<u>Ring No</u>	<u>E</u>	<u>SO</u>	<u>S</u>	<u>U</u>	<u>Total</u>
1	3	4	3	0	10
2	0	8	3	5	16
3	2	6	7	4	19
4	0	0	5	5	10
5	0	3	15	6	24
6	1	3	12	7	23
7	2	1	6	4	13
8	0	1	7	2	10
9	1	2	4	6	13
10	0	1	5	0	6
11	2	3	7	4	16
12	1	0	11	0	12
13	0	0	11	3	14
14	0	2	7	2	11
15	0	0	8	5	13
16	0	1	7	1	9
17	1	3	14	4	22
18	0	0	12	0	12
19	0	1	12	2	15
20	2	1	12	2	17
21	0	1	9	2	12
22	1	1	20	2	24
23*	1	0	17	0	18
24*	1	0	13	2	16
25*	0	0	13	3	16
26*	2	1	6	3	12
27*	0	3	12	0	15
28*	0	1	8	0	9
29*	0	0	9	4	13
30*	0	1	11	1	13
31*	0	1	7	5	13
32*	0	0	10	3	13
	<hr/>	<hr/>	<hr/>	<hr/>	<hr/>
	20	49	303	87	459

\* ring overran a field edge.

Table 2.2

Comparison of the numbers of images, by category, as classified during the initial (A) and repeated (B) inspections of IIIa-J plates.

	Subregion 8		Subregion 19		Subregion 37	
	Centre 10h44m -25°45'		Centre 10h30m -26°31'		Centre 10h40m -28°45'	
	(A)	(B)	(A)	(B)	(A)	(B)
(1)	1800	1806	2472	2477	2658	2659
(2)	100	95	137	120	201	194
(3)	9	6	1	14	2	6
(4)	4	6	6	5	4	6
(5)	0	0	1	1	1	1

Description of categories:

(1): stars and spurious images; (2): galaxies; (3): uncertain objects; (4): image blends; (5): dissociated images (single images recorded by the measuring machine in two or more discrete components).



## CHAPTER 3

### REDSHIFTS

Knowledge of the redshifts of a sample of galaxies in the field of any cluster is very valuable, because it provides the only reasonably reliable means of identifying those systems not gravitationally associated with major clustering, and because only redshifts can define the internal dynamics of a cluster. Both applications assume that the redshift is purely a Doppler effect. Controversial evidence for intrinsic quantised components in the redshifts of galaxies in the cores of clusters has been reviewed by Tifft (1974).

Redshifts are inferred from spectra, obtained either individually at high wavelength dispersion (spectrograms) or en masse at lower resolution (e.g. objective prism plates). All new redshifts reported in this chapter for galaxies in the field of A1060 have been obtained from prism plates, and are useful only in assessments of cluster membership. However, the opportunity has also been taken to reassess the velocity dispersion of the cluster, as defined by accurate radial velocities that are already in print. This is done in conjunction with estimates of the luminosities of the galaxies for which redshifts are available (Chapter 5).

#### 3.1 Redshifts from Objective Prism Spectra

Very large numbers of galaxy spectra can be recorded simultaneously on fine-grain photographic plates exposed behind an objective prism, and in particular that of the U.K. Schmidt telescope (Nandy et al. 1977). The redshifts of early-type galaxies in the sample may be

estimated by comparison of the redshifted and rest-frame positions of characteristic features in the galaxy spectra, as recorded, for example, in graphical form (Cooke et al. 1977) or in digitized form, following measurements by the COSMOS machine (e.g. Nandy 1979). The integrated spectra of such galaxies - E, SO and Sa/Sb - is dominated by emission from late-type stars which exhibit a sharp drop in intensity shortwards of about  $4000\text{\AA}$  in their rest frame. This "step" is sufficiently clear in low-dispersion prism spectra to serve as a redshift indicator.

UKST prism plate UJ4067P, an exposure centred on the Hydra I cluster, was used for this purpose. The particulars of UJ4067P are summarized in Table 3.1. Galaxies were initially selected for tracing on the basis of their apparent brightness (approximately  $14.5 < V_{25} < 17.5$  in the notation of Chapter 5), position within the COSMOS CM scan area of the previous chapter (for identification purposes), and freedom from blending on the CM charts. In practice, spectral overlaps resulted in many rejections, whereas a few image blends were in fact successfully decomposed. Unfortunately an important area of plate 47 arcminutes wide and symmetric about the length of the central meridian through the cluster, was physically inaccessible to measurement on the microdensitometer. Consequently, no systems within 270 kpc of the cluster centre could be examined.

The spectra of 108 galaxies satisfying the luminosity, isolation and accessibility criteria were recorded; twenty of these were traced once or more during independent recording sessions. Calibration was provided by tracings of field stars. Nineteen



distinct lines and blends in the spectra of four stars, fitted to the dispersion curve of the UKST objective prism (Nandy et al. 1977), gave  $\lambda_c = 5433 \pm 11$  ( $1 \pm .17$ )(s.e.) Å as the value of the wavelength cutoff of the emulsion of plate UJ4067P. In turn, this fiducial mark fixed the wavelength of the 4000Å feature, defined in nine late-type stars, as  $\lambda_{4000} = 3990 \pm 9$  ( $1 \pm .24$ )(s.e.) Å, assuming zero error in  $\lambda_c$ . This figure is consistent with the published value of  $\lambda_{4000} = 4010 \pm 30$  Å (Nandy et al. 1977). The wavelength cutoff for UJ4067P differs significantly from the value,  $5360 \pm 30$  Å, of Nandy et al., but one should perhaps be wary of possible systematic external errors in the present work. The values of  $\lambda_{4000}$  and  $\lambda_c$  determined for UJ4067P were adopted on the grounds that they lead to internally consistent redshifts, a decision that avoids the spurious blueshifts that would otherwise result.

Of the recorded sample, sixty percent, spiral galaxies figuring prominently, lacked an unambiguous 4000Å feature and had to be discarded. The redshifts obtained for the remaining 43 galaxies are presented in Table 3.2. Nine of the listed values represent the means of repeated determinations, and the mean absolute internal error for these estimates is only  $\langle |z| \rangle = 0.005$ . A better indication of the accuracy of the redshifts would follow if they could be compared with independent measurements of a second prism plate or, preferably, a high-dispersion spectrogram obtained on some other telescope. As matters stand, the errors in the tabulated redshifts are probably less than  $\pm 0.02$  (cf. Cooke et al. 1977).

Eight galaxies in Table 3.2 are identified as probable outlying members of A1060 on the evidence of their redshifts. Not one is fainter than  $V_{25} = 16$ , but the background would be expected to dominate at faint magnitudes in the outskirts of the cluster. The low density of cluster members in regions distant from the core of Hydra I would be accentuated in any sample of redshifts obtained by the procedure outlined above, because the proportion of late-type members in the cluster outskirts is higher and this class of galaxy does not exhibit the spectral features that permit redshifts to be determined on objective prism plates. Table 3.2, nevertheless, is sufficient to establish that two galaxy groupings which exist in this field have mean redshifts of  $\sim 0.07$  and  $0.08$  respectively, clearly identifying both as distant groups at seven times the distance of Hydra I.

### 3.2 The Velocity Dispersion

The mean velocity of recession and the internal velocity dispersion of A1060 has previously been estimated by Vidal & Peterson (1975), Faber & Dressler (1976, 1977) and by Rood & Dickel (1978). Harrison & Noonan (1979), however, have illustrated that considerable ambiguity exists in the interpretation of redshift data, so it seems worthwhile to repeat the velocity analysis below on the basis of previously published heliocentric redshifts and newly determined photographic magnitudes.

The information sources are listed in the References section to Table 3.3, which summarizes the existing redshift data for all

galaxies within a projected radius of 5 Mpc from the cluster centre. Weights inversely proportional to the square of the internal errors have been adopted in the calculations of velocity averages, after first converting the probable errors quoted by Ref. 1 and the mean errors of Refs. 4 and 5 to corresponding standard deviations as used by Ref. 3, and assigning, arbitrarily, an internal standard deviation of  $75 \text{ km s}^{-1}$  to the velocity estimates of Ref. 2. Where only one velocity is available the standard deviation of that determination has been substituted for the standard error of the mean, the latter an obviously uncertain quantity since multiple observations of any system are few.

The distribution of the averaged heliocentric redshifts (expressed as "indicative velocities",  $cz_0$ ) is shown in Figure 3.1, which is evidently ambiguous. As a spatial distribution along the line of sight it may be distorted greatly by peculiar motions. The small sample does not present a classical gaussian velocity distribution (cf. Yahil & Vidal 1977), yet none of the redshift extremes differs by an amount, 2.5 or 3 times the standard deviation of the original velocity distribution, that would justify its elimination as a non-member. This casts doubt on the analysis of Rood & Dickel (1978) of Sandage's (1978) limited sample of galaxies in A1060, since Rood & Dickel reject three of his seven systems.\* On the other hand, the line-of-sight depth of the sample would be large -  $20(100/H_0) \text{ Mpc}$  - if no peculiar motions were involved, and those bright spirals in the low-velocity tail of the distribution, including N3312, have redshifts similar to that of the foreground

\* A1060 = Hydra I is misidentified as A732 = Hydra II due to a misprint in Sandage (1978).



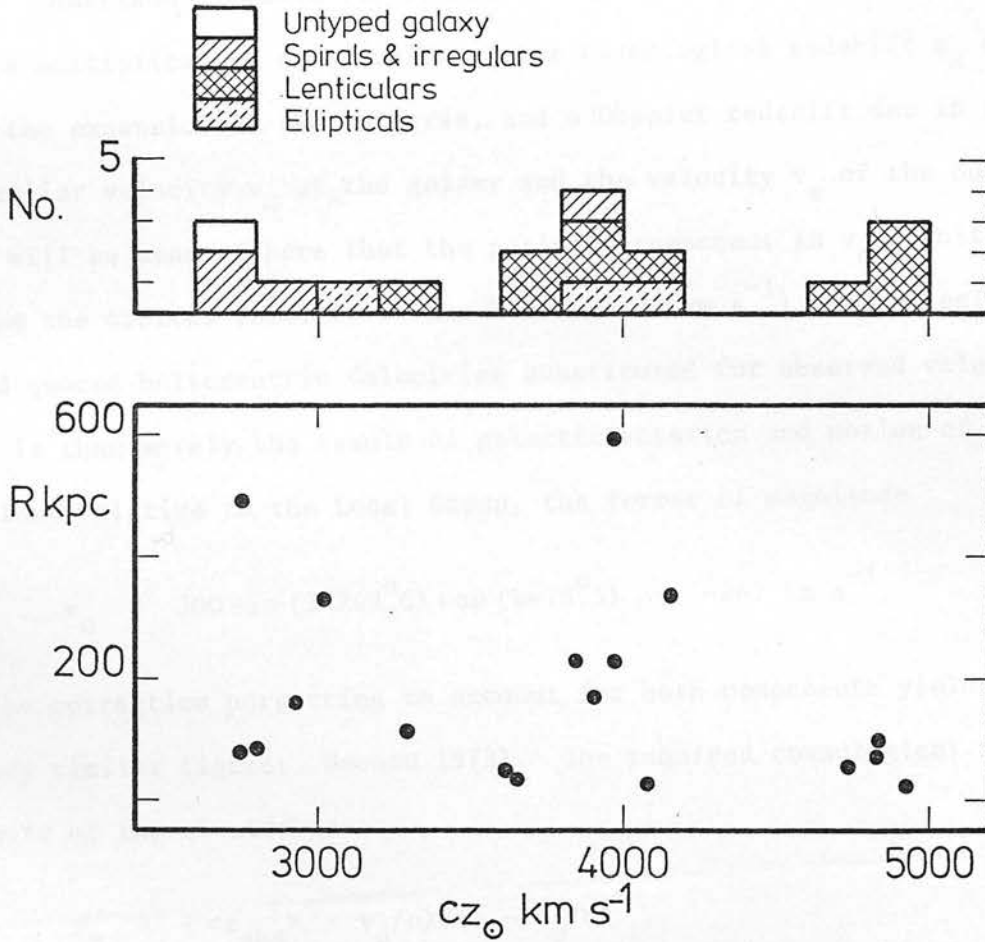


Figure 3.1

The distribution of the indicative heliocentric radial velocities listed in Table 3.3.  $R$  is the projected separation in kpc from the adopted cluster centre.



Hydra cloud (Chapter 2). For the present discussion, all galaxies in Table 3.3 will be treated as members of A1060.

Harrison & Noonan (1979) remark that any observable redshift  $z$  is a multiplicative combination of the cosmological redshift  $z_R$  caused by the expansion of the universe, and a Doppler redshift due to the peculiar velocity  $v_G$  of the galaxy and the velocity  $v_o$  of the observer. It will be assumed here that the periodic component in  $v_o$  resulting from the orbital velocity of the Earth ( $\leq 30 \text{ km s}^{-1}$ ) can be neglected and quoted heliocentric velocities substituted for observed values;  $v_o$  is then merely the result of galactic rotation and motion of the galaxy relative to the Local Group, the former of magnitude

$$v_o = 300 \sin(l=269.6) \cos(b=26.5) = -267 \text{ km s}^{-1}.$$

(The correction purporting to account for both components yields a very similar figure: Noonan 1973). The required cosmological redshift of the cluster is

$$z_R = (\langle z_{\text{obs}} \rangle + v_o/c) / (1 - v_o/c)$$

where  $z_{\text{obs}} \approx z_o$ , and the mean square dispersion of the radial velocities in the cluster is

$$s^2 \equiv \langle v_G^2 \rangle = (\langle z_{\text{obs}}^2 \rangle - \langle z_{\text{obs}} \rangle^2) / (1 + z_R)^2$$

in which the angular brackets  $\langle \rangle$  denote mass-weighted values. The galaxy masses are calculated on the basis of absolute  $V_{25}$  luminosities and representative mass-to-light ratios of 5, 10 and  $15 M_\odot/L_\odot$  for spiral galaxies, lenticulars and ellipticals respectively. If

galaxies one and twelve in Table 3.1 are discounted because of conflicting redshift estimates and the remaining galaxies treated as cluster members, the recessional and dynamical parameters of Hydra I are

$$\langle z_{\text{obs}} \rangle = 0.01221 \pm 0.00042 \text{ (s.e.)},$$

$$z_{\text{R}} = 0.01131 \pm 0.00043 \text{ (s.e.)},$$

$$cz_{\text{R}} = 3390 \pm 129 \text{ km s}^{-1},$$

$$S = 496 \text{ km s}^{-1}.$$

The cosmological radial velocity of the centre of mass,  $cz_{\text{R}}$ , determined above differs insignificantly from the value ( $3324 \text{ km s}^{-1}$ ) given by Faber & Dressler (1976), but the velocity dispersion of  $496 \text{ km s}^{-1}$  is some 40% lower. This is mainly a consequence of the lack of mass weighting in the calculations of Faber & Dressler (and of Vidal & Peterson 1975). A fully rigorous computation would in fact yield a still lower estimate of  $S$ , since part of the measured dispersion arises from observational errors in the redshifts (cf. Materne 1974). Estimation of the confidence limits of the velocity dispersion is not straightforward; naive application of the  $\chi^2$  distribution implies limits for 68.3% confidence of  $459 \leq S \leq 651 \text{ km s}^{-1}$ , but this neglects the uncertainties in the galaxy masses. The 68.3% confidence limits for the cosmological radial velocity  $cz_{\text{R}}$  have been calculated on the basis of the  $t$ -distribution, assuming error-free masses and a 10% uncertainty in the velocity  $v_0$  of the observer.

The downward revision of  $S$  from  $771 \text{ km s}^{-1}$  (the value adopted

for A1060 by Faber & Dressler 1976) to  $496 \text{ km s}^{-1}$  is of some consequence regarding correlations between  $S$  and other properties of clusters generally. The correlation between velocity dispersion and X-ray luminosity is a particularly relevant example, in connection with the evaluation of models of X-ray emission from clusters. Because A1060 has the lowest X-ray luminosity ( $L_X$ ) of any Abell cluster yet observed, a large change in the estimate of the velocity dispersion (or in that of the  $L_X$ ) of A1060 can alter appreciably the form of the correlation between  $L$  and  $S$  at low  $L_X$  (see, for example, Figure 5 of Dawe et al. 1977 or Figure 3, Vidal & Peterson 1975). Specifically, a more linear relationship between  $S$  and  $L_X$  results from the downward revision noted above.

It should be emphasised, however, that most (or all) of the previous determinations of the velocity dispersions in X-ray clusters (as opposed to nearby galaxy groups) have neglected the matter of mass-weighting of individual galaxies, usually in the absence of suitable photometry. If the masses of the galaxies observed in a cluster are appreciably unequal, this neglect may result in the velocity dispersion being overestimated, at least in a cluster dominated by a few bright and centrally located galaxies. It seems likely that the depth into the cluster luminosity function that is represented by each sample of redshifts is greater in nearby, poorer clusters such as A1060 than in richer and more distant clusters in which spectroscopic observations of faint members are not easily secured. Mass-weighting of the redshifts would then be essential for the derivation of comparative values of  $S$ , and this requires

optical photometry of the cluster members involved, in conjunction with realistic mass-to-light ratios.

Table 3.1

Particulars of UKST prism plate UJ4067P

Date exposed	4 - 5 April 1978
R.A. (1950)	$10^{\text{h}} 34^{\text{m}}$
Dec. (1950)	$-27^{\circ} 24'$
Emulsion	IIIa-J, hypersensitized
Grain type	Fine
Exposure (minutes)	35
Prism orientation (degrees)	180
Trailing	none

Table 3.2

The redshifts of 43 galaxies, determined by measurements of UKSTU prism plate UJ4067P. Notation is that used in Appendix II.

<u>No.</u>	<u>Z</u>	<u>Suggested Membership</u>	<u>No.</u>	<u>Z</u>	<u>Suggested Membership</u>
90	.00	A 1060	374	.10	
95	.02	A 1060	378	.10	
99	.00	A 1060	404	.05	
133	.00	A 1060	433	.08	C1 1037-260
136	.01	A 1060	449	.10	
157	.04		451	.05	C1 1030-269
169	.07		472	.04	
178	.09		496	.08	
179	.02	A 1060	509	.10:	
183	.00	A 1060	538	.08	C1 1037-260
220	.02	A 1060	563	.14	C1 1037-260?
236	.08	C1 1030-269	574	.06	
249	.05		579	.09	
268	.08	C1 1030-269	587	.07	C1 1037-260?
275	.07:		627	.07	
278	.09		742	.06	C1 1037-260
291	.07		775	.12	
318	.08		803	.06	C1 1030-269
319	.08		814	.07	
323	.09		848	.08	
352	.12		958	.09	C1 1037-260
364	.04				



Table 3.3

Compilation of radial velocities of galaxies in the field of A1060.

VP	FD	No	$\alpha$ (1950.0)			$\delta$			$cz_{\odot}$	s.e. n		Identi- fication	Refs	
			h	m	s	0	'	"						
1	J	70	10	33	42.5	-27	14	59	(1)			A1033-2715	3,5	
2	S	96	10	33	49.6	-27	11	39	4829	25	1		3	
3	10	6	16	10	33	49.9	-26	54	09	3968	4	2	N3305	3,4
4	C	7	5	10	34	00.5	-27	10	43	3611	61	3	N3308	3,4,5
5		11	19	10	34	03.1	-26	44	23	4158	88	1		4
6	O	66	10	34	03.1	-27	19	20	2742	56	1		3	
7		10	53	10	34	05.9	-27	03	34	3292	88	1		4
8	M	46	10	34	07.3	-27	13	29	4927	133	1		3	
9	B	1	2	10	34	14.3	-27	15	31	4084	11	3	N3309	3,4,5
10	P	37	10	34	20.0	-27	18	03	4736	50	1	A1034-27A	5	
11	A	2	1	10	34	21.5	-27	16	06	3650	41	4	N3311	2,3,4,5
12	L	61	10	34	23.4	-27	12	35	(2)				3,6	
13	8	14	10	34	37.0	-27	55	01	2753	16	1		3	
14	D	3	3	10	34	41.1	-27	18	18	2796	43	4	N3312	1,3,4,5
15	F	B	10	34	51.5	-27	25	24	2916(3)	120	1	N3314	3	
16		8	24	10	34	55.5	-27	12	30	4836	4	2	A1034-27B	4,5
17		5	10	10	34	57.5	-26	55	55	3840	88	1		4
18	E	9	9	10	35	15.9	-27	20	01	3906	61	3	N3316	1,3,4
19	5	4	4	10	35	25.5	-26	49	12	3020	9	2		3,4
20	2	7	10	37	55.3	-27	30	57	3970	-	1		3	

Column 1: running number; col. 2: VP = designation from Vidal & Peterson (1975); col. 3: FD = designation from Faber & Dressler (1977); col. 4: designation in Appendix II; cols. 5-10: precise epoch 1950 coordinates,  $\pm 2''$  (supersede approximations in Appendix II); col. 11: adopted weighted mean heliocentric velocity in  $\text{km s}^{-1}$ ; col. 12: notes; col. 13: standard error in  $cz_{\odot}$ ,  $\text{km s}^{-1}$ ; col. 14: number of observations; col. 15: designation as in RC2; col. 16: references.

#### Notes

(1):  $cz_{\odot} = 2389 \pm 96$  (s.d.) (Vidal & Peterson 1975);  $4147 \pm 50$  (m.e.) (Sandage 1978). Refer to different galaxies? A1033-27 in RC2 is probably object 32, Appendix II.

(2):  $cz_{\odot} = 2772 \pm 59$  (s.d.) (Vidal & Peterson 1975);  $2280 \pm ?$  (Fairall 1979).

(3): N3314 consists of two spiral galaxies superposed along line of sight.

#### References

1: Mayall & de Vaucouleurs (1962); 2: Peterson (1970); 3: Vidal & Peterson (1975); 4: Faber & Dressler (1977); 5: Sandage (1978); 6: Fairall (1979).

## CHAPTER 4

### PHOTOELECTRIC PHOTOMETRY

The most reliable method of calibration of photographic photometry of galaxies utilising the COSMOS measuring machine is shown, in chapter 5, to involve the comparison of the projected distributions of photographic and photoelectric intensities within suitable standard galaxies. The present chapter describes a series of multicolour and multiaperture photoelectric observations of sixteen galaxies in the field of A1060, carried out on the 1.0m and 0.5m reflectors at S.A.A.O., Sutherland in early 1978 for the purpose of plate calibration. The characteristics of the photoelectric systems are outlined, and an assessment is made of the photometric errors. The derivation of "total" magnitudes of galaxies from the photoelectric data is also discussed.

#### 4.1 Instrumentation

The observations were secured between 31 January to 7 February 1978 with the St Andrews Photometer (StAP) on the S.A.A.O. 1.0m Elizabeth reflector, and during 28 March to 11 April 1978 initially with the 1.0m and StAP and later with the People's Photometer (PP) on the 0.5m reflector. The StAP and PP are single-channel pulse-counting photometers with EMI 6256A S11 photocathodes for UVB photometry operating at a nominal  $-10^{\circ}\text{C}$ ; VBR response is catered for in the PP by an EMI 9659A extended S20 which functions optimally at  $-25^{\circ}\text{C}$ . The filter combinations applying to each of these systems is given in Table 4.1. Red leak across the Schott UG-2 filters defining the U wavebands is

assumed to be negligible (cf. Shao & Young 1965).

Trials were carried out to ascertain the behaviour of the photometric systems in their operational configurations. Short integrations performed as a star drifted centrally across the largest useful apertures in each of two mutually perpendicular directions confirmed that the response across the fields of the photocathodes was acceptably uniform. The dead time constants of the blue channels of the StAP and PP, which characterise the losses of counts caused by simultaneous arrival of photons, were determined to be 86 and 48 nanoseconds respectively; that of the red channel of the PP was approximately 30 ns.

Focal plane scales were estimated by timing the intervals between bisections of a star image by opposite edges of selected diaphragms with the telescopes stationary. A value of  $13.8 \text{ mm}^{-1}$  was obtained for the 1.0m, in agreement with Corwin (1980), but a trend exists whereby smaller diaphragms give numerically larger plate scales. This effect may also be indicated in data from the 0.5m, for which a scale of  $25.1 \text{ mm}^{-1}$  is adopted. The linear dimensions (S.A.A.O. and Corwin, private communications) and corresponding angular sizes of the 1.0m and 0.5m diaphragms are presented in Tables 4.2a and 4.2b. In general the tabulated values differ substantially from those quoted in the S.A.A.O. Facilities Handbook for 1977, and supersede them.

Multiaperture measurements of a standard star on the 1.0m telescope reveal a dependence of measured magnitude on diaphragm size caused by light being scattered out of the smaller apertures. For larger diaphragms the correction required was independent of seeing

and amounted to less than  $0.01^m$  for apertures greater than 20 arc-seconds. The use of a 20" aperture to isolate the comparison stars therefore introduced negligible systematic error.

#### 4.2 Observations and their Reduction

A single observation on the 1.0m or 0.5m represented the programmed filter sequence V, B, U, U, B, V, with 30-second integrations in each waveband, without recentring. Observations were alternated between each galaxy and one of a set of "sky holes" determined beforehand by reference to deep IIIa-J UK Schmidt plates to be free from contamination by field stars within apertures of interest down to a magnitude of  $B \approx 20$ . All sky holes were at least 10 arcminutes distant from the nominal cluster centre (NGC 3311), beyond which any background glow of intracluster origin should be negligible. Typically 3 galaxy and 2 sky observations were made per selected object on any night; where a galaxy was measured on more than one night, averages were adopted.

The measurements were tied to the UBV system through the use of Cousins' (1973) photometric standards (expanded compilation, available at S.A.A.O.), which were observed at least twice each night at dusk and dawn on the 0.5m and at least three times a night on the 1.0m. Linear transformation coefficients, with a non-linear correction in B-V, determined from measurements on the 1.0m have been used to reduce instrumental magnitudes to the UBV system. Data from the 0.5m have been reduced through linear and non-linear colour equations derived from observations with the blue and the red channels

of the PP separately. Throughout, mean extinction coefficients have been used.

Strictly, to obtain accurate galaxy photometry a correction for light scattered out of the smaller measuring diaphragms, which depends on the luminosity profile of the galaxy, should be taken into account. However since the main purpose of this photometry is to calibrate photographic plates where presumably a similar scattering function is already present, no attempt has been made to correct the photometry for this effect. It appears that only Weedman (1975) has applied corrections of this sort to his data.

In the course of the observations three aberrant stars were found, the implied colour indices of which differed by modest but consistent amounts from those listed by Cousins (1973). The correct indices and V magnitudes are listed in Table 4.4. A fourth star, E393, was also discrepant, but because the measurements were not repeated, the possibility of misidentification cannot be discounted. In addition, the star E336 may have a faint companion.

#### 4.3 Results

The results of the photoelectric photometry are presented in Table 4.3. The galaxy designations in the first column of the table are those employed in Appendix II, which may be consulted to identify the galaxies in the Appendix I finding charts. In Table 4.3, for convenience, standard NGC designations are also given. Column (2) contains the morphological index, T, corresponding to the classification of the galaxy in the de Vaucouleurs (1959) system. Column (3)



lists the telescope aperture in metres, column (4) the diaphragm diameter in arcseconds, and the measured UBV magnitudes and colours are tabulated in columns (5) to (7). Column (8) notes the number of nights on which observations were repeated, where an asterisk indicates that a U-B value on one night was discarded (1.0m) or was not available during VBR photometry (0.5m). V-R colour indices from the latter and their repetition factors are given in columns (9) and (10).

The R band is in the Cape-Kron system (Bingham & Cousins 1974). At an effective wavelength of about 6500 Å the Cape-Kron R is some 450 Å shortward of Johnson's R band (e.g. Johnson 1966); in particular, direct comparison is precluded of the present V-R indices with those measured for galaxies in several southern clusters by Sandage (1975).

The internal accuracy of the photometry may be ascertained by the consistency of measurements from one night to the next. The mean absolute values, and their formal standard deviations, of 53 residuals in each colour from among the 1.0m photometry are

$$\begin{aligned} V : \langle |\Delta| \rangle &= 0.012 \pm .013 \\ B - V : \langle |\Delta| \rangle &= 0.008 \pm .008 \\ U - B : \langle |\Delta| \rangle &= 0.020 \pm .020 \end{aligned}$$

In general the residuals are magnitude dependent and increase from about half the quoted value at  $V = 13$  to twice the quoted value at  $V = 15$  (Figure 4.1). Similar analysis of the 0.5m data on the brighter galaxies suggests that the mean internal errors of the 0.5m photometry



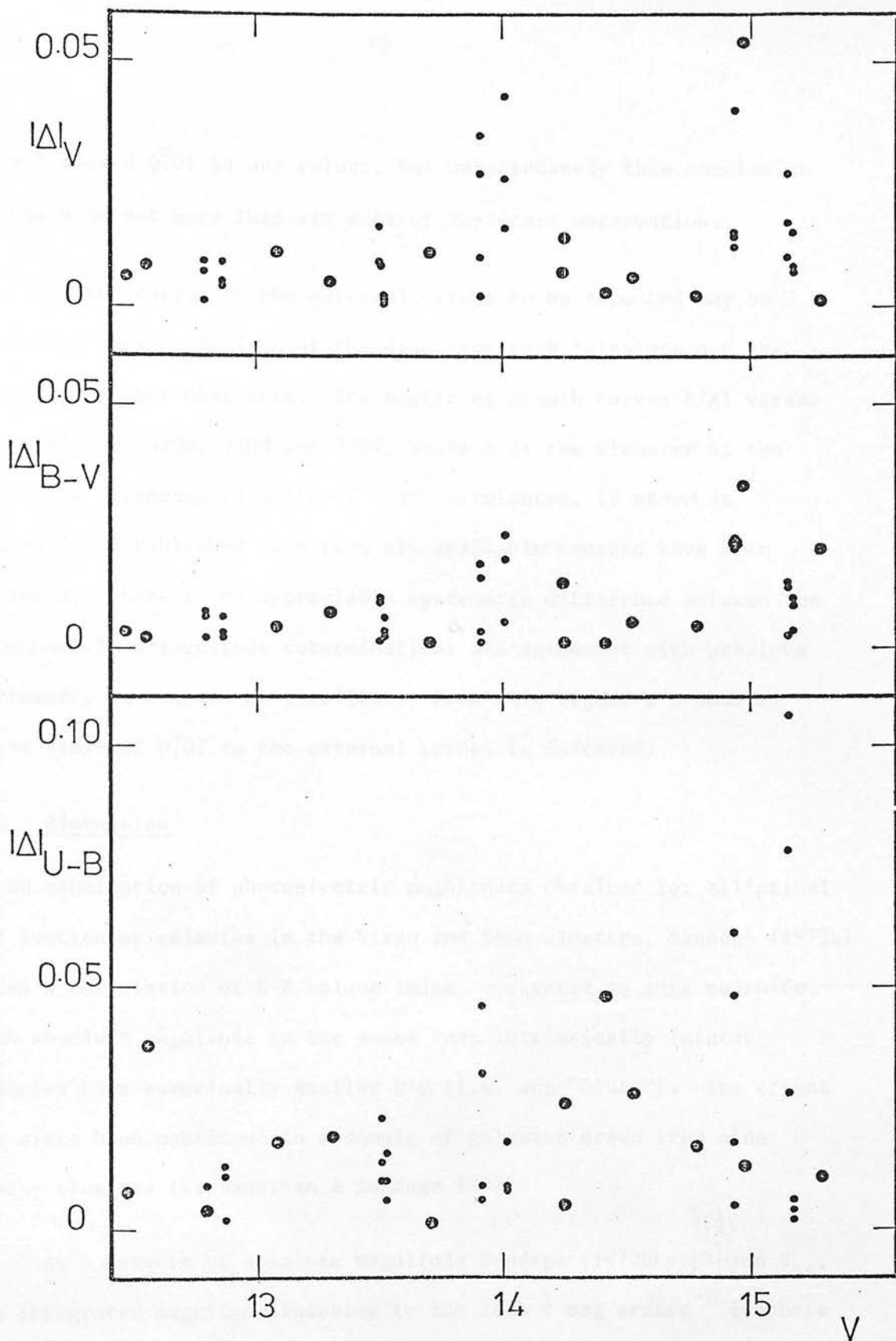


Figure 4.1

The dependence on observed V magnitude of the absolute values of the V, B-V, and U-B residuals from among the 1.0m photometry. Small dots are single points, large dots are double points.

do not exceed  $0.01^m$  in any colour, but unfortunately this conclusion is based on not more than six sets of duplicate observations.

An indication of the external errors to be expected may be gained by intercomparison of the data from each telescope and the results of other observers. The magnitude growth curves  $B(A)$  versus  $\log A$  of NGCs 3305, 3308 and 3309, where  $A$  is the diameter of the measuring diaphragm in units of tenth arcminutes, is shown in Figure 4.2. Published data from all available sources have been included. There is no appreciable systematic difference between the 1.0m and 0.5m B-magnitude determinations and agreement with previous photometry is seen to be excellent. From this figure a probable upper limit of  $0.07^m$  to the external errors is inferred.

#### 4.4 Discussion

In an examination of photoelectric magnitudes obtained for elliptical and lenticular galaxies in the Virgo and Coma clusters, Sandage (1972b) noted a correlation of U-B colour index, corrected to zero redshift, with absolute magnitude in the sense that intrinsically fainter galaxies have numerically smaller U-B (i.e. are "bluer"). The effect has since been confirmed in a sample of galaxies drawn from nine nearby clusters (Visvanathan & Sandage 1977).

As a measure of absolute magnitude Sandage (1972b) employed  $V_{26}$ , the integrated magnitude interior to the 26th V mag arcsec<sup>-2</sup> isophote (adjusted to correspond to an object at the distance of the Coma cluster). In the absence of suitable photographic photometry  $V_{26}$  was estimated by means of a standard magnitude growth curve, originally

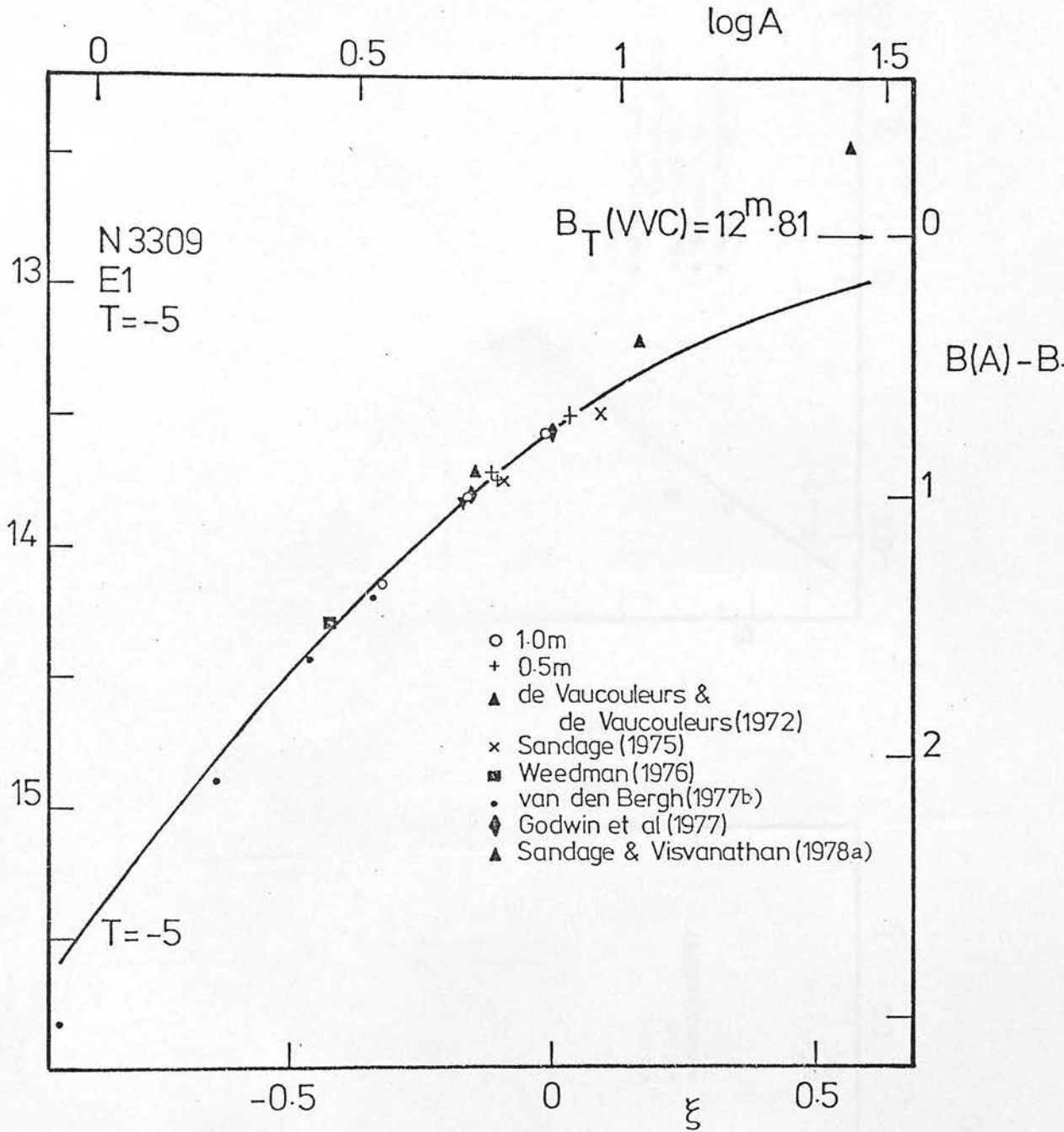


Figure 4.2 (in three parts)

B-magnitude growth curves of three NGC galaxies derived from all available photometry.  $\xi$  is the logarithm of the normalised aperture  $A/A_e$ , where  $A_e$  is that aperture containing half the total B luminosity of each galaxy.

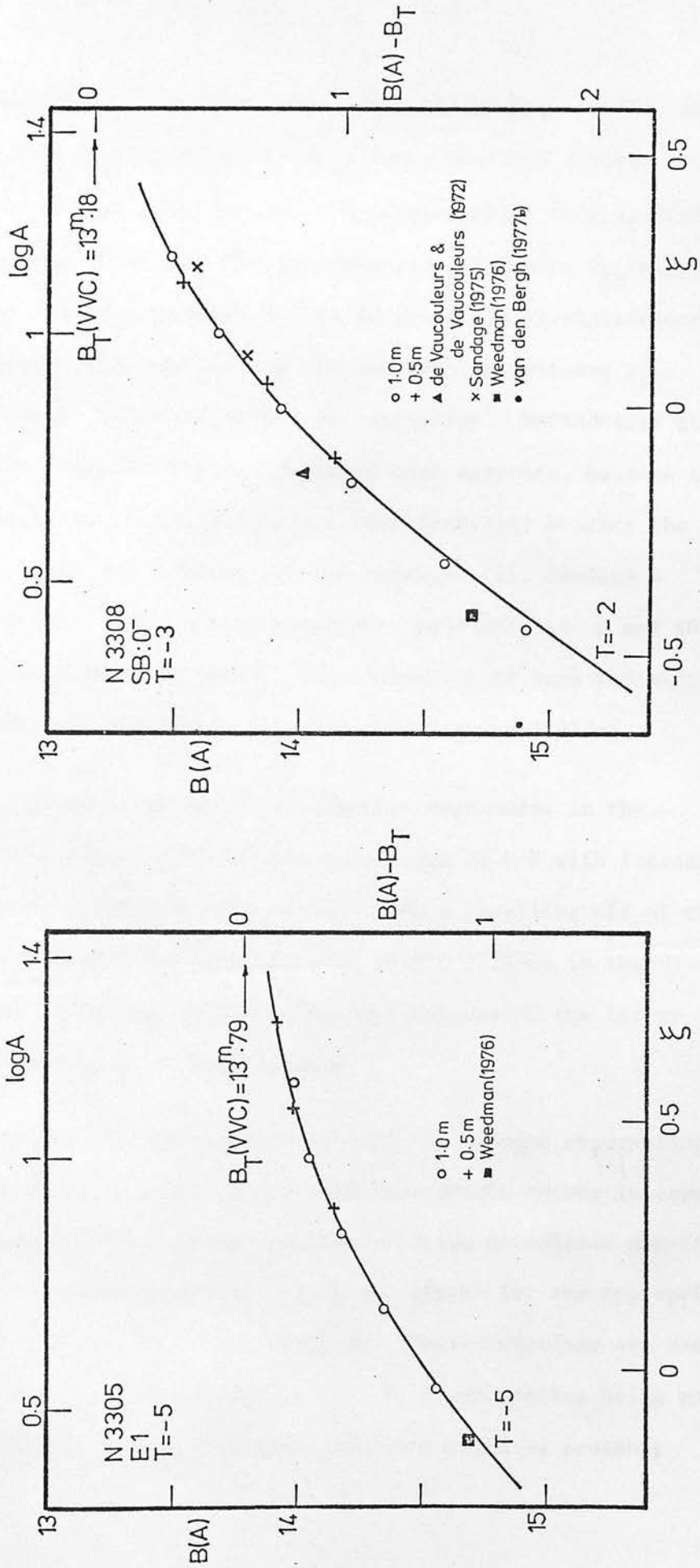


Figure 4.2 (continued)

defined by Humason et al. (1956) and revised by Sandage (1975), which was used to reduce photoelectric V magnitudes measured through diaphragms of various sizes to the values appropriate to a specific fractional galaxy diameter. The determination of metric  $V_{26}$  magnitudes of the galaxies listed in Table 4.3 is in principle straightforward, but it is expedient to use instead the isophotal magnitudes  $V_{25}^c$ , defined in chapter 5, as indicators of luminosity. Furthermore it is permissible to adopt U-B indices averaged over aperture, because the effective dimensions of the diaphragms used generally bracket the face-on major axis semidiameters of the galaxies (cf. Sandage & Visvanathan 1978a), and U-B colour-aperture gradients in E and SO galaxies are comparatively small. The correction to zero redshift (reddening the observed  $\langle U-B \rangle$ ) is taken from Sandage (1972b).

Within the range of the four brightest magnitudes in the Hydra I cluster, Figure 4.3 confirms a decrease of U-B with increasing  $V_{25}$  for galaxies fainter than  $V_{25} \approx 13.5$ , and a levelling off of the relation at the bright end which is also weakly evident in the original study of Sandage (1972b). The implications of the latter result are discussed in a later chapter.

It is of interest to enquire how well the present observations are accommodated by standard galaxy magnitude growth curves in common use. In Figure 4.2 are featured the best-fitting normalised magnitude-aperture relations of de Vaucouleurs et al. (1976) for the appropriate morphological indices, T, of the galaxies. These functions are seen to represent the new observations very well, discrepancies being most apparent among previous photometry of NGC 3309 which is probably

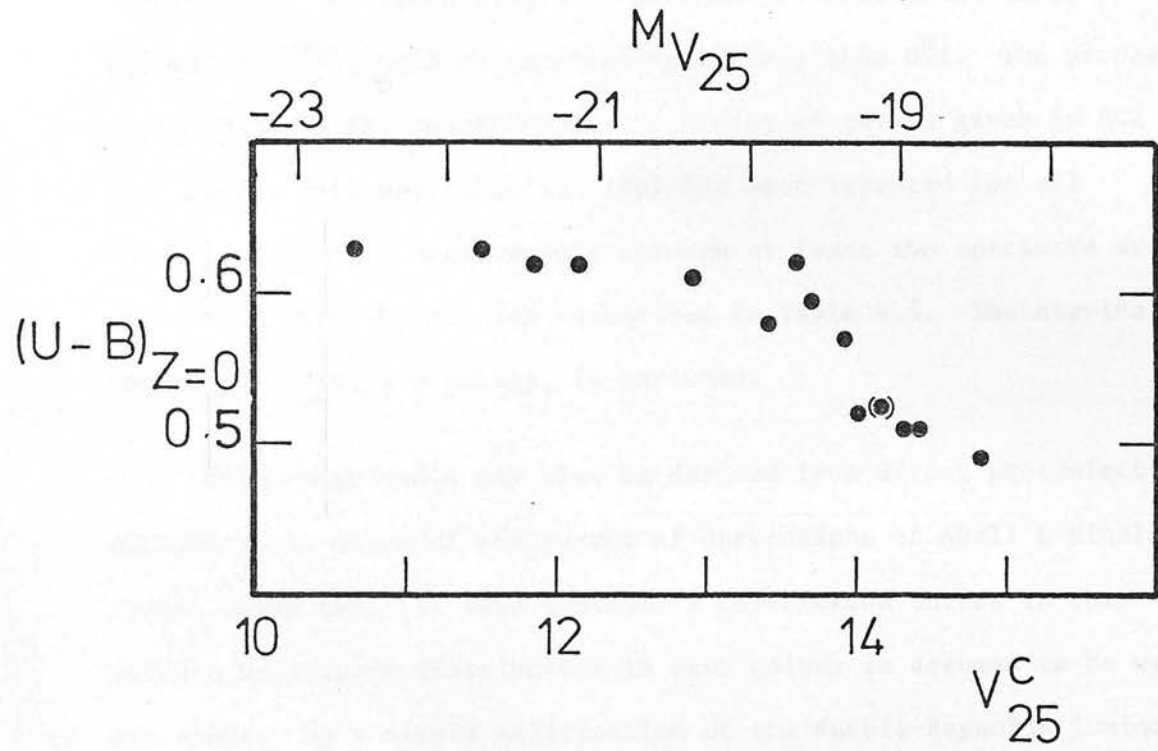


Figure 4.3

The correlation of U-B colour index, corrected to zero redshift, with isophotal magnitude. Parentheses distinguish the S0/a galaxy.



under-luminous at apertures less than 15 arcseconds because of the loss of scattered light, and conversely, affected by contamination of light from the halo of NGC 3311 and a 14th (B) magnitude star at apertures exceeding 56 arcsec. From the fits the "total" or asymptotic B magnitude ( $\log A \rightarrow \infty$ ) of NGC 3309 is found to be  $B_T(\text{VVC}) = 12.^m81$ , with an uncertainty of less than  $0.^m1$ . The process of matching to the observations the family of curves given in RC2 (pp. 26-31, with approximation 17b) has been repeated for all galaxies for which measurements through at least two apertures are available, with the results summarized in Table 4.5. The atypical object NGC 3311, a D galaxy, is excluded.

Total magnitudes may also be derived from direct photoelectric photometry by means of the curves of corrections of Abell & Mihalas (1966, Table IV). In this instance a restriction enters in that the surface brightness distribution in each galaxy is assumed to be well represented by a simple modification of the Hubble-Reynolds luminosity profile; only different degrees of apparent ellipticity are distinguished. The prescription of Abell & Mihalas leads to the extrapolated values denoted by  $B_T(\text{AM})$  in Table 4.5.

Highly satisfactory fits to the present data, with the exception noted in Table 4.5, are obtained by both the above methods. The fundamental  $B_T$  magnitudes agree most closely for elliptical galaxies. The discrepancies of as much as half a magnitude between  $B_T(\text{AM})$  and  $B_T(\text{VVC})$  for the lenticulars and the transition spiral\* reflect the departure of disk galaxies from the luminosity distribution law assumed by Abell & Mihalas. The likelihood that a similar systematic

\* The optical appearance of this object is ambiguous and in fact the observed colours (Fig. 4.3) are consistent with those of a normal S0 or E.

error affects total magnitudes derived by the related procedure of "extrafocal" photographic photometry (Abell & Mihalas 1966) deserves consideration, particularly in studies of the radial distribution of galaxies of different luminosities within a cluster that exhibits marked segregation of galaxy types (cf. Abell 1977).

The reduction procedure employed by Sandage and collaborators to obtain  $V_{26}$  magnitudes suffers from an essentially equivalent drawback: photoelectric data must be fitted to a single magnitude growth curve without regard of the mean surface brightness, intrinsic ellipticity or precise morphological type of a galaxy. For objects of atypical light distribution (e.g. N3311), the procedure is acknowledged to be fallible (Sandage 1975; Schild & Davis 1979) and photographic surface photometry provides a more reliable method of deriving isophotal magnitudes  $V_{25}$  or  $V_{26}$ .

Table 4.1

Photomultiplier tubes and filter combinations.

	<u>StAP blue</u>	<u>PP blue</u>	<u>PP red</u>
photo- multiplier	EMI 6256A S11	EMI 6256A S11	EMI 9659A extended S20
V	2mm OG-515	2mm OG-302	3mm GG-495 +1mm BG-18
B	4mm GG-13 +1mm BG-12	4mm GG-13 +1mm BG-12	2mm OG-1 +2mm OG-23 +1mm BG-12
U	1mm UG-2	1mm UG-2	
R			2mm OG-570

Table 4.2a

Dimensions of diaphragms of the St Andrews Photometer on the 1.0m Elizabeth reflector.

Nominal aperture (arcseconds)	True aperture (mm)	True aperture (arcseconds)
5	0.378	5.1
7	0.637	8.8
10	0.758	10.4
14	1.115	15.3
20	1.509	20.9
28	2.188	30.1
40	3.088	42.6
56	4.329	59.6
80	6.128	84.6

Table 4.2b

Dimensions of single-hole (S) and double-hole (D) diaphragms of the People's Photometer on the 0.5m reflector. The blue (B) and red (R) channels of (D) diaphragms are distinguished. The qualifying colon indicates that measured horizontal and vertical diameter differ by more than two percent.

Nominal aperture (mm)	True aperture (mm)	True aperture (arcseconds)
0.32S	0.294:	7.4
0.63D(B)	0.636:	16.0
(R)	0.664	16.7
0.67D(B)	0.656	16.5
(R)	0.657:	16.5
0.73D(B)	0.652	16.4
(R)	0.639	16.0
1.0S	1.030	25.9
1.0D(B)	0.983	24.7
(R)	0.970	24.3
1.02D(B)	1.043	26.2
(R)	1.019:	25.6
1.141D(B)	1.371	34.4
(R)	1.352	33.9
1.5D(B)	1.404:	35.2
(R)	1.402	35.2
1.56D(B)	1.521	38.2
(R)	1.540	38.7
1.93D(B)	1.912	48.0
(R)	1.911	48.0
2.0S	1.970	49.4
2.0D(B)	2.017	50.6
(R)	2.014	50.6
2.02D(B)	2.018	50.7
(R)	2.024	50.8
3.0S	3.010	75.6
3.0D(B)	3.003	75.4
(R)	2.997	75.2
4.5S	4.486	113
4.5D(B)	4.372:	110
(R)	4.424	111
6.0S	6.017	151
6.0D(B)	6.006	151
(R)	5.961	150
8.0S	8.034	202
8.0D(B)	8.052	202
(R)	8.053	202
10.0S	9.999	251

Table 4.3

Photoelectric observations of galaxies in the Hydra I cluster.  
For a guide to this table see the text.

<u>Galaxy</u>	<u>T</u>	<u>Tel</u>	<u>Ap</u>	<u>V</u>	<u>B-V</u>	<u>U-B</u>	<u>n</u>	<u>V-R</u>	<u>n</u>
1=N3311	-2	1.0	42.6	12.85	1.07	0.64	1		
		0.5	48.0	12.75	1.05	0.61*	2	0.63	1
		1.0	59.6	12.47	1.06	0.61	2		
		0.5	75.2	12.23	1.06		1	0.64	1
2=N3309	-5	1.0	20.9	13.08	1.07	0.65	2		
		1.0	30.1	12.75	1.06	0.58	1		
		0.5	33.9	12.63	1.09		1	0.62	1
		0.5	34.4	12.69	1.06	0.62	1		
		1.0	42.6	12.50	1.07	0.61	1		
		0.5	48.0	12.44	1.06	0.64*	2	0.65	1
4=I2597	-5	1.0	20.9	13.08	1.05	0.62	1		
		1.0	30.1	12.79	1.05	0.59	3		
		0.5	33.9	12.71	1.03		1	0.64	1
		1.0	42.6	12.55	1.04	0.63	2		
		0.5	48.0	12.48	1.04		1	0.63	1
5=N3308	-3	1.0	10.4	14.24	1.06	0.63	1		
		1.0	15.3	13.85	1.05	0.59	1		
		1.0	20.9	13.52	1.06	0.61	3		
		1.0	30.1	13.16	1.04	0.63	1		
		0.5	33.9	13.07	1.06		2	0.61	2
		1.0	42.6	12.86	1.06	0.61	3		
		0.5	48.0	12.82	1.05	0.59*	2	0.63	1
		1.0	59.6	12.62	1.04	0.61	1		
		0.5	75.2	12.50	1.04		1	0.67	1
		1.0	84.6	12.44	1.05	0.61*	1		
12=N3305	-5	1.0	20.9	13.50	1.06	0.60	3		
		1.0	30.1	13.30	1.05	0.59	2		
		1.0	42.6	13.13	1.05	0.59	1		
		0.5	48.0	13.11	1.04		1	0.59	1
		1.0	59.6	13.02	1.03	0.58	1		
		0.5	75.2	12.93	1.05		1	0.64	1
		1.0	84.6	12.95	1.04	0.63	1		
		0.5	111.0	12.89	1.04		1	0.67	1
22	-2	1.0	20.9	13.91	1.06	0.57	4		
		1.0	30.1	13.71	1.05	0.57	2		
		0.5	33.9	13.63	1.06		2	0.64	1
		0.5	48.0	13.51	1.04		2	0.66	1
24=N3315	-2:	1.0	20.9	14.25	1.08	0.59	2		
		1.0	30.1	13.97	1.07	0.66	1		

contd...

Table 4.3 (contd.)

<u>Galaxy</u>	<u>T</u>	<u>Tel</u>	<u>Ap</u>	<u>V</u>	<u>B-V</u>	<u>U-B</u>	<u>n</u>	<u>V-R</u>	<u>n</u>
31	-2	1.0	20.9	14.53	1.05	0.54	2		
		1.0	30.1	14.25	1.05	0.58	2		
		1.0	42.6	14.01	1.04	0.51	1		
38	-2	1.0	20.9	14.01	1.05	0.58	3		
		1.0	30.1	13.83	1.04	0.59	1		
		0.5	33.9	13.78	1.06		1	0.62	1
		0.5	48.0	13.70	1.04		1	0.64	1
56	-3	1.0	20.9	14.78	1.00	0.50	2		
60=N3307	1	1.0	20.9	14.93	1.00	0.44	4		
61	-5	1.0	20.9	14.41	0.99	0.51	2		
66	0:	1.0	20.9	15.15	1.00	0.54	3		
		1.0	30.1	14.97	0.98	0.49	2		
68	-5	1.0	20.9	14.76	1.01	0.50	1		
86	2	1.0	20.9	15.17	0.87	0.28	3		
92	-3	1.0	20.9	15.28	1.00	0.48	2		



Table 4.4

Corrections to the photometry of Cousins (1973, 1974) of stars in the Harvard E-regions. The last column records the number of nights during which each star was observed.

<u>Star</u>	<u>V</u>	<u>B-V</u>	<u>U-B</u>	<u>n</u>
E393	7.175	1.577	1.892	1
E491	6.843	-0.122	-0.527	2
E583	7.974	0.984	0.663	4
E607	8.788	0.120	0.098	2

Table 4.5

Extrapolated total magnitudes  $B_T$  of galaxies in the Hydra cluster. A less reliable determination is enclosed in parentheses.

<u>Galaxy</u>	<u>Type</u>	<u>T</u>	<u><math>B_T</math> (VVC)</u>	<u><math>B_T</math> (AM)</u>
2	E1	-5	12.80	12.65
4	E3	-5	(12.55)	12.70
5	SB:0 <sup>-</sup>	-3	13.20	12.75
12	E1	-5	13.80	13.80
22	SO <sup>0</sup>	-2	14.35	14.15
24	SO:	-2:	14.45	13.90
31	SO	-2	14.60	14.10
38	SO pec	-2	14.60	14.45
66	SO/a?	0?	15.75	15.35

## CHAPTER 5

### PHOTOGRAPHIC PHOTOMETRY

Early photographic photometry of galaxies involved the laborious construction of detailed two-dimensional luminosity profiles of single galaxies from repeated one-dimensional microphotometric scans (e.g. de Vaucouleurs 1948; Holmberg 1950, 1958). The high precision of magnitudes derived by this procedure more than justifies its continuing application (e.g. Fraser 1977, Kormendy 1977a) in studies of modest numbers of galaxies.

The relatively great speed and labour-saving potential of the new generation of rapid two-dimensional digitising microphotometers with auxiliary computers was nevertheless recognised by Jones et al. (1967), who describe numerical mapping techniques to extract, from data generated by these machines, useful photographic parameters of galaxies. Godwin (1976) has published abridged versions of these procedures which facilitate the processing by computer of the large numbers of images involved in studies of clusters. Independent data reduction schemes (e.g. Oemler 1976; Strom & Strom 1977, 1978a; MacGillivray & Dodd 1979) as well as adaptations of the Jones et al. or Godwin procedures (e.g. Okamura 1977, Carter & Dixon 1978, Green 1978) are now in widespread use.

In this chapter the derivation of isophotal magnitudes of several thousand galaxies in the field of A1060 by the analysis of measurements made by the COSMOS machine is described.

### 5.1 Plate Material

The photographic plates of the Hydra I cluster were taken with the UK 1.2m Schmidt telescope by members of the Schmidt Unit during 1975 and 1977. Details of the relevant plates are given in Table 5.1.

The emulsion-filter combinations are those routinely employed by UKSTU to imitate photographically the Johnson B and V photoelectric wavebands (Johnson & Morgan 1951, 1953). The resulting photographic responses cannot be expected to match precisely the effective wavebands of the photoelectric photometry secured for plate calibration, and such mismatching of bands could introduce errors of the photographic zeropoints which correlate with the colours of the calibrators involved. Those employed in this work (Table 4.3) represent too narrow a range of colour to indicate whether the effect is present, so negligible disparity of photoelectric and photographic wavebands is assumed by default.

All plates listed in Table 5.1 were measured on the COSMOS machine in either or both of two measurement modes (Table 5.2). A second B and a U-band plate were also measured, but were not processed to obtain magnitudes. Furthermore it proved unnecessary to calibrate the R plate photoelectrically. Calibration is considered in section 5.4.

COSMOS job number 254 was a pilot study involving the detailed examination of a small number of galaxies; the remaining jobs run under mapping mode each generated four magnetic tapes of data. One tape from job 301 was unfortunately corrupt and the remaining three

were used only to check the photometry of plate V3324.

## 5.2 Processing Techniques

The photographic photometry of galaxies undertaken in this study is of two types. The primary photometry is based on the application of modified Godwin (1976) routines to data generated by COSMOS in mapping-mode operation; the secondary photometry, of lower precision, is obtained from coarse-mode parameters output by the COSMOS hardware pattern analyser and is tied to the primary photometric system. The primary B and V systems are calibrated by photoelectric observations of standard galaxies (Chapter 4). Only V-band magnitudes are defined by the secondary photometry.

The principal virtue of the primary photometry lies in the ability of the adopted routines to handle most of the instances of overlapping images frequently found in the core regions of dense clusters of galaxies, and more generally within regions of enhanced density of field stars. At the galactic latitude of A1060 ( $b \approx 26^\circ$ ), the isophotes of brighter members are contaminated more often than not by light from neighbouring galaxies or superposed stars; such blends are not generally recognised for what they are by the COSMOS coarse-mode pattern analyser (now de-commissioned), and when assessing the secondary photometry this shortcoming must be accounted for.

During the course of the work several modifications to the processing routines of the primary photometry were introduced, and will be described in the following outline of the procedure.

The output from COSMOS in its mapping mode consists of arrays of up to  $3200 \times 3200$  transmission readings representing an area of plate, and additional arrays of measurements across either or both step wedge calibrators impressed on the plate. This amount of data is obviously too large to be held in entirety in a computer core for random access, so subsamples of the main array are processed in instalments. To exploit more fully the computing power available in the present work, the original routines were uprated by 50% to handle individual grids of  $316 \times 316$  picture elements. A subroutine was incorporated to generate the arrays.

A mesh of  $44 \times 44$  integers, referred to as the "coarse" or "outer" grid, is first extracted to model the variation of the sky background in a chosen locality. The value of each array element represents the median of  $15 \times 15$  transmission readings; the median is a more stable quantity than the mode and is less sensitive than the latter to biases in analogue-to-digital converters, although in the mean it lies slightly above the true sky level on account of the skewness of the distribution of transmissions.

Within this matrix of sky values, the 25% of all points encompassed by a specified interior square are removed. A two-dimensional third-order polynomial is then fitted to the reduced array by a process of repeated rejection of points more than  $\beta$  standard deviations distant from the most recent approximating surface, where on successive cycles  $\beta = 2, 2.25, 2.25$ , and  $2.5$ . Only significant terms of this polynomial are retained after each cycle. Early exit is arranged if any cycle rejects no further points;



typically 12% of the points, mainly in the vicinity of bright images, are eliminated by this means. More than sufficient array elements remain at the conclusion of the process to define the final polynomial adequately, whereas the computational demands of coarse grid construction are more than offset by the greater speed of fitting a polynomial to fewer points (cf. Barbon et al. 1976).

A "fine" or "inner" grid interior to the coarse array and positioned over the effective hole in the latter is next constructed from 316 x 316 raw transmission readings. The relative geometry of coarse and fine grids is shown in Figure 5.1. This is a generalisation of Godwin's (1976) Figure ii in that the fine array is now permitted to take up any enclosed position relative to the coarse square, according to parameter values specified at execution time. The effect of this is to reduce the inaccessible border zone from 20% to 1% of the total area of the COSMOS scan. The relative dimensions of the grids employed satisfy the criterion of stability laid down by Jones et al. (1967) for third-order interpolation across the region of the inner grid, though it should be noted that extrapolation is invoked should one or more boundaries of inner and outer grids coincide. The interpolation procedure ensures that bright images contained entirely within the inner grid do not affect the local background determination.

Each pixel in the fine grid is transformed to a scale of relative intensity according to the Baker characteristic of the plate (section 5.3) and compared with the value of the background polynomial, likewise transformed, at the same position. In this way a grid of intensity



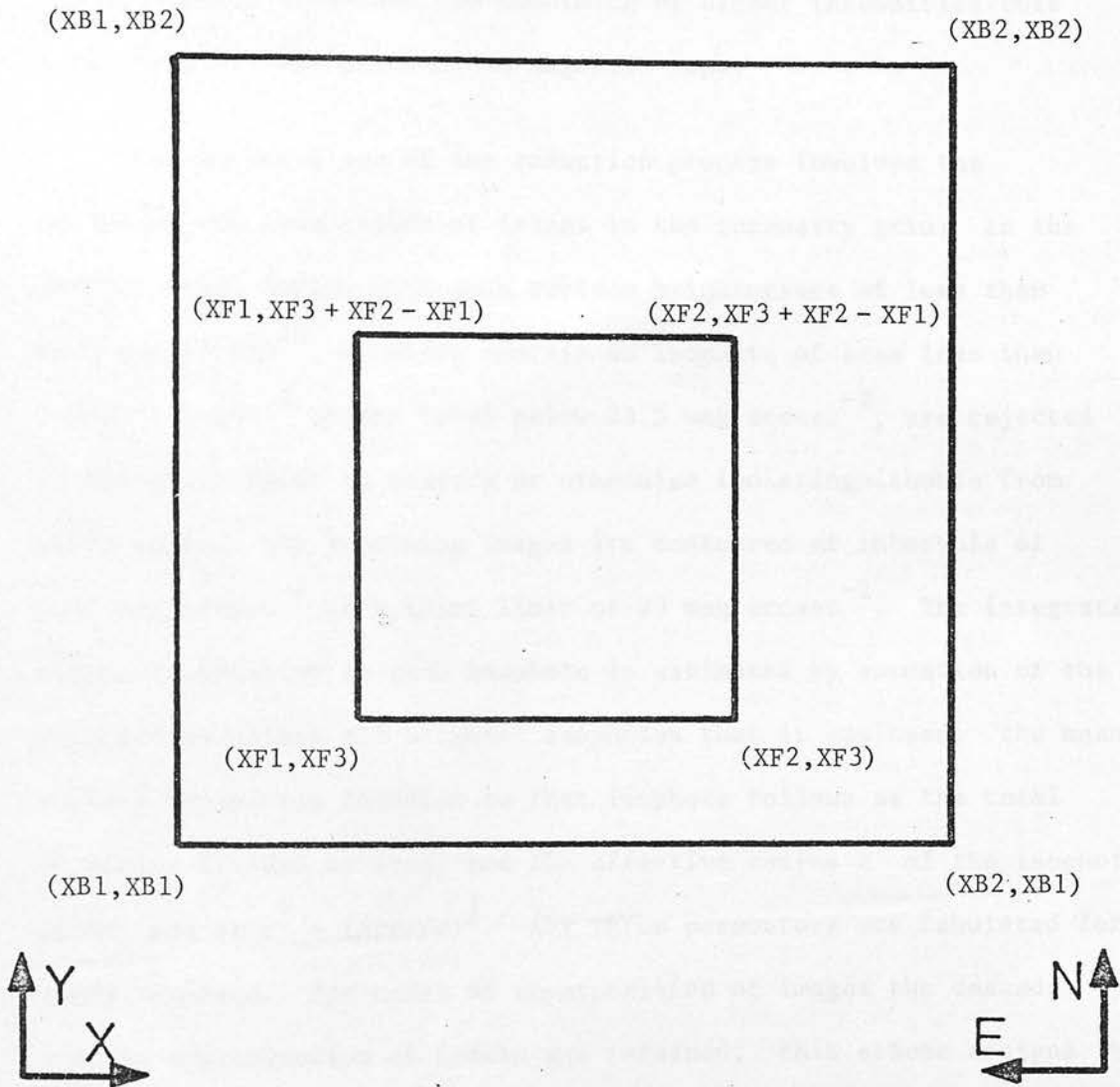


Figure 5.1

The relative geometry of the inner and outer grids.

readings expressed in units of the subtracted night sky background (de Vaucouleurs & Page 1962) is generated. After adjustments to the background level and the smoothing of higher intensities this array is stored temporarily on magnetic tape.

The second stage of the reduction process involves the selection and examination of images in the intensity grid. In the present study images with peak surface brightnesses of less than  $23.5 \text{ mag arcsec}^{-2}$ , or which contain an isophote of area less than  $5.0625\pi \text{ arcsec}^2$  at any level below  $23.5 \text{ mag arcsec}^{-2}$ , are rejected as being too faint to measure or otherwise indistinguishable from plate noise; the remaining images are contoured at intervals of  $0.25 \text{ mag arcsec}^{-2}$  to a faint limit of  $25 \text{ mag arcsec}^{-2}$ . The integrated magnitude interior to each isophote is estimated by summation of the intensities within all brighter isophotes that it encloses; the mean surface brightness interior to that isophote follows as the total intensity divided by area, and the effective radius  $r^*$  of the isophote is defined as  $r^* = (\text{Area}/\pi)^{1/2}$ . All three parameters are tabulated for every isophote. For cases of superposition of images the cascade overlap approximation of Godwin was retained; this scheme assigns the luminosity of succeeding (descending) isophotes to overlapping images in the ratio of the areas of their faintest unconnected isophotes. Since compact objects with steep brightness profiles have relatively large inner isophotes, superposed stars tend to become overluminous at the expense of galaxies in this photometry. Interactive picture processing schemes (e.g. Barbon et al. 1976, Strom & Strom 1977) are better equipped to handle superpositions, but final accuracy, about  $\pm 0.25^m$ , is never high.

Particular modifications were applied as follows. The calibration scheme was reworked to express the relationship between transmission and relative intensity in terms of a Baker characteristic (section 5.3). The coupling between the inner and outer grids was altered. The original method whereby the two sets of measurements were reduced to a common system by comparing transmissions round the perimeter of the inner grid with the values at corresponding positions in the outer grid was discarded as untenable because image contributions were present in one grid but not in the other. Instead a correction was applied to shift the mode of the frequency distribution of intensities in the inner grid to the zero level after the background was subtracted. Repeated checks showed that the variable skewness of the distribution caused by the presence of bright images within the inner grid was inconsequential compared to statistical uncertainties in the position of the mode arising from the quantisation of the transmission measurements. This, approximately 1% at the sky intensity, effectively limits the precision with which individual corrections can be determined, but by comparing the reductions of many grids the necessary adjustment was reckoned as +0.5% on average for the IIaD emulsions and +1.0% for the IIaO. The use of median transmission values to approximate the sky level thus results in a slight overestimate, as expected. These corrections were adopted for later reductions, and previous processing was repeated in cases of gross error. During reductions of plate B3028, every frame was examined individually before final zeropoint adjustments were made.

Numerical smoothing was effected in a contrary sense, to high intensities rather than to low. This reduced the number of cases in

which transmission fluctuations in the high-intensity domain produced spurious discrete peaks within the cores of authentic images; similar treatment of lower intensities, to reduce noise, was discontinued when surface brightness profiles were found to be altered.

Further necessary but trivial modifications included redefined functions and graphical instructions for the construction of isophotal maps. In addition a library of related programs was developed for the display and inspection of photometric data, descriptions of which are omitted.

### 5.3 Transformation to Intensities

On plates taken with the UK Schmidt telescope, guidance on the transformation of photometric measurements on some instrumental system is presently provided by a system of two step wedges. The prototype wedge consists of translucent metal film deposited on a glass base in seven contiguous steps of successively decreasing transparency; the second wedge is a photographic duplicate. In the telescope illuminating lamps leave photographic impressions of the wedges within two unexposed regions near the north and east edges of the plate masked off during the simultaneous sky exposure. Subsequent measurement of any plate by COSMOS in mapping mode operation includes a longitudinal scan of at least one wedge impression.

A regrettable shortcoming of this arrangement that is conspicuous at shorter wavelengths is the faintness of the lamps deployed, in relation to the central surface brightnesses of the galaxies that register on the plate. As a result the impressions of the wedge

calibrators seldom exhibit opacitances as high as those measured in the cores of brighter images, unless the photographic emulsion saturates, as in the case of sky-limited IIIa-J plates. For IIa0, IIaD and 098 plates of representative exposure times, however, this problem is overshadowed by the effective upper limit to the dynamic range of the COSMOS measuring machine.

The relative intensity values corresponding to the steps of each wedge have been measured photoelectrically by Tritton (1977, UKSTU memorandum). The conflict between his figures and those of the less direct derivations of Dawe (1978 revision, privately circulated) is considered below.

Probably the most elegant expression empirically to relate the intensity of illumination to the density induced in the exposure was proposed by Baker (1925) and independently by de Vaucouleurs (1968). If RI denotes relative intensity and D optical density, the empirical Baker function or "generalised Sampson" transform is

$$RI = A(10^{D/B} - 1)^n$$

where A, B and n are constants. This relation also obtains as a limiting case of the Kaiser function (Honeycutt & Chaldu 1970) in the absence of saturation. It frequently offers a linear representation of the characteristic curve of a photographic emulsion, enabling calibrations to be constructed on the basis of few data points, but it is nevertheless fallible, as Barbon et al. (1976) have illustrated. The most satisfactory representation of a general characteristic curve may in fact involve a seven-parameter function (Tsubaki &



Engvold 1975). In particular, the breakdown of the linearity of the Baker function of densely exposed plates has led to the use of less wieldy alternatives (e.g. Schweizer 1976, Kormendy 1977a, Okamura 1977, Carter & Dixon 1978, Carter & Godwin 1979, Dufor et al. 1979), frequently polynomials in the opacitance  $\omega = 10^{D/B} - 1$ .

The expression for the opacitance appropriate to the COSMOS machine is  $\omega = T_0/T - 1$ , where  $T$  is transmission and  $T_0$  represents that of the chemical fog level. The enumeration of four parameters, including the small correction term in transmission that represents the reading for zero incident light intensity, is therefore required to define the Baker characteristic. Due to the established operational practice of adjusting the effective transmission range to ignore photographic intensities fainter than the minimum sky level, the value of  $T_0$  must be estimated by repeating the wedge measurements at different instrumental settings, or by relating the original wedge transmissions to the density readings recorded on, for example, a Joyce-Loebl isodensitometer or Macbeth densitometer. Both methods have been applied here to each determination.

An example of a wedge tracing is shown in Figure 5.2. At each position along the axis the recorded transmission is the mean of 128 samples laterally and may be non-integral. The appreciable non-uniformities of the wedge steps and the consequent difficulty of defining their mean levels constitute the principal limitation of the calibrations. Figure 5.3 shows the Baker functions for this wedge (north; upper) according to the relative intensity values supplied by Tritton and by Dawe. The greater linearity inherent in the former



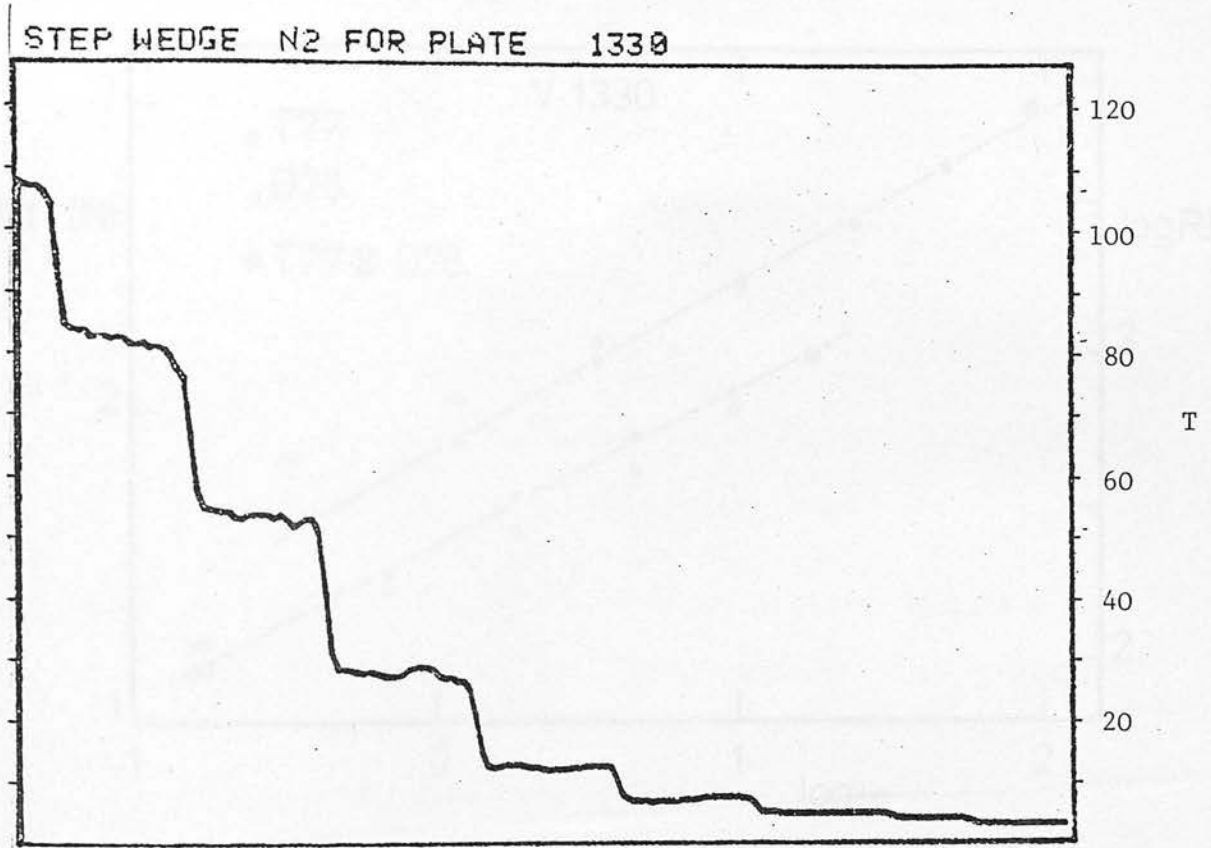


Figure 5.2

Example of a step wedge impression measured by COSMOS. The densest step level is lowermost. T denotes transmission.

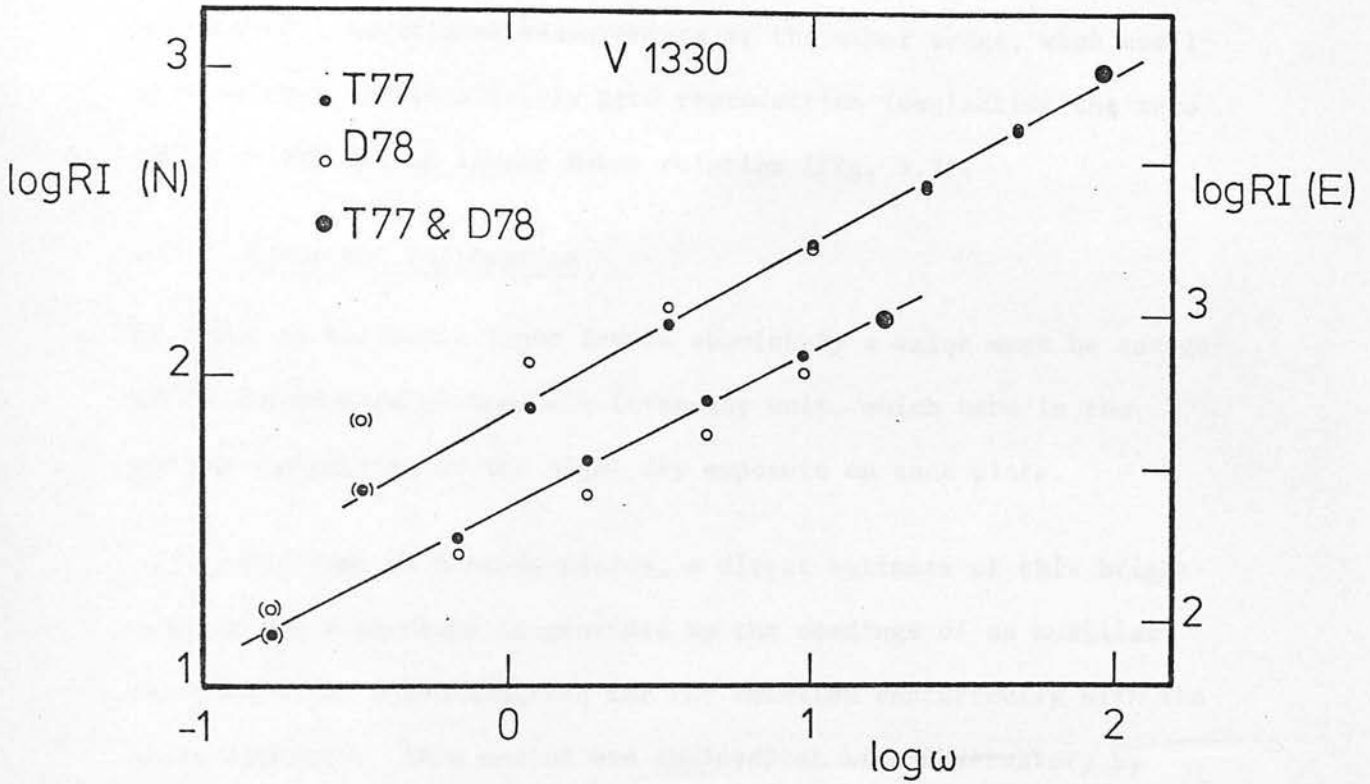


Figure 5.3

The Baker characteristic of plate V1330 as defined by the north step wedge (upper relation, left scale) and by the east step wedge (lower relation, right scale), using the relative intensity values of Tritton (T77) and of Dawe (D78). The lines are fits to Tritton's values; the upper relation was adopted. The zeropoints of the intensity scales are arbitrary.

calibration is also apparent in every other plate/filter combination so far investigated. For this reason the figures of Tritton were adopted provisionally, subject to confirmation\*. Of the two wedges available, preference was normally given to the more densely exposed north wedge in order to maximise the range of opacities thus calibrated; additional measurements of the other wedge, when available, always led to a fairly good reproduction (neglecting the zero point shift) of the linear Baker relation (Fig. 5.3).

#### 5.4 Zeropoint Calibration

In order to calibrate light levels absolutely a value must be assigned to the adopted photometric intensity unit, which here is the surface brightness of the night sky exposure on each plate.

For some UK Schmidt plates, a direct estimate of this brightness in the B waveband is provided by the readings of an auxiliary sky photometer which monitors the sky emission concurrently with the plate exposure. This method was employed at Lick Observatory by Dressler (1978b). Unfortunately the original calibration of the UKST photometer in 1973 was in error by some  $+0.6 \pm .16$  magnitudes and the source is decaying at the rate of about  $0.1 \text{ year}^{-1}$  (MacGillivray, private communication). The photometer does not define sky brightnesses in wavebands other than B.

Photoelectric stellar sequences have been shown (e.g. King & Hinrichs 1967, Kormendy 1973) to be capable of defining both the intensity scale and the zeropoint of a system of photographic measurements to an accuracy of  $\sim 0.1^m$ . One such sequence has been set up in

\* The most recent estimates of the wedge intensities quoted by UKSTU (October 1979) take intermediate values.

the core of the Hydra I cluster by Smith & Weedman (1976). However in the present photometry exploitation of this sequence is seriously hampered by the lack of effective resolution of the spot scanner in the COSMOS machine, the restricted range of surface brightness that is measurable, and most fundamentally, due to the location of virtually all of the standard stars within regions exhibiting luminosity gradients caused by the proximity of very bright images.

The most satisfactory method of estimating the zeropoints has proved to be based on comparisons of integrated photographic magnitude profiles of standard galaxies with reference photoelectric profiles determined from multiaperture photometry. With suitable observations (Table 4.3) profiles can be defined over ranges of several tenths of arcminute, avoiding the effects of resolution and atmospheric seeing (cf. Carter 1979, Schweizer 1979) which dominate the central six arcseconds of images on Schmidt plates.

An example of a magnitude growth curve defined by photoelectric observations is plotted as circles in Figure 5.4. The total photoelectric intensity interior to various radii on a scale of arbitrary zeropoint may be found by interpolation and compared with the corresponding integrated photographic intensity obtained by simulated aperture photometry of the plate image (Fig. 5.5). If  $I^s$  denotes the zeropoint error in the total photographic intensity caused by saturation of transmission readings at the centre of the image, then at any radius  $r$  external to the saturated core of semidiameter  $r_{\text{sat}}$

$$I_{\text{pg}}(\leq r) + I^s = \text{constant} \times I_{\text{pe}}(\leq r), \quad r > r_{\text{sat}}$$

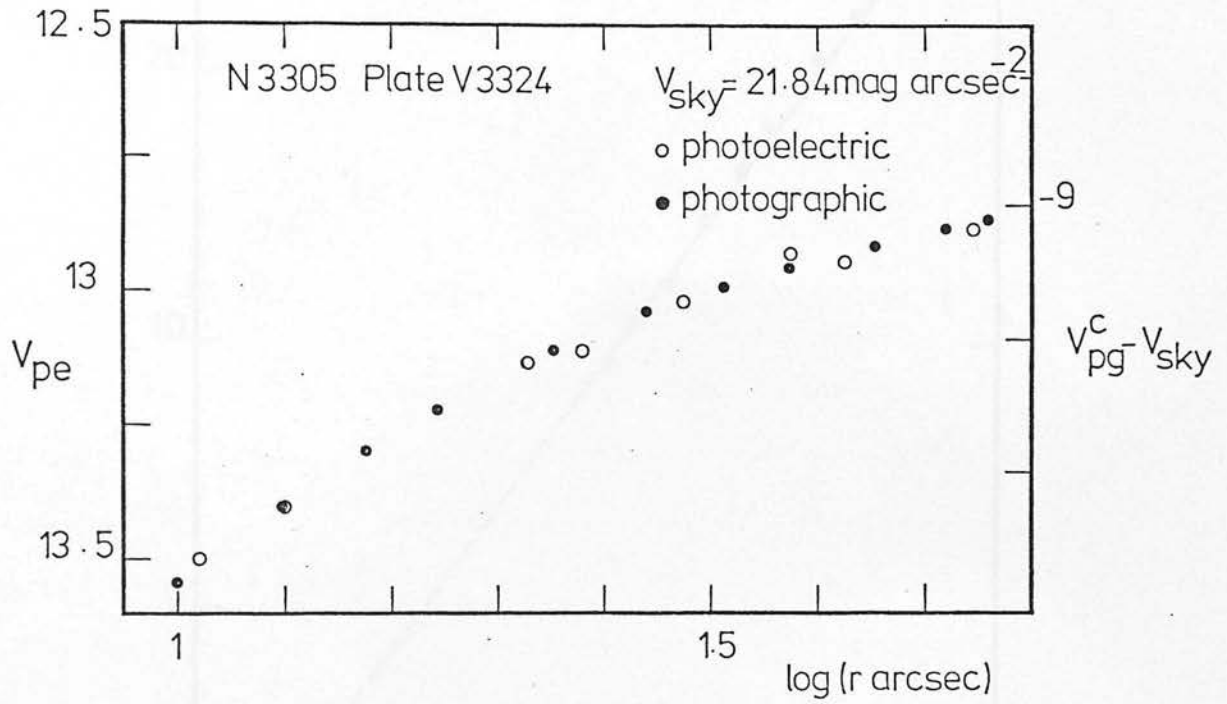


Figure 5.4

The optimally matched photoelectric ( $V_{\text{pe}}$ ) and corrected photographic ( $V_{\text{pg}}^c$ ) magnitude profiles of NGC 3305 in the V band.

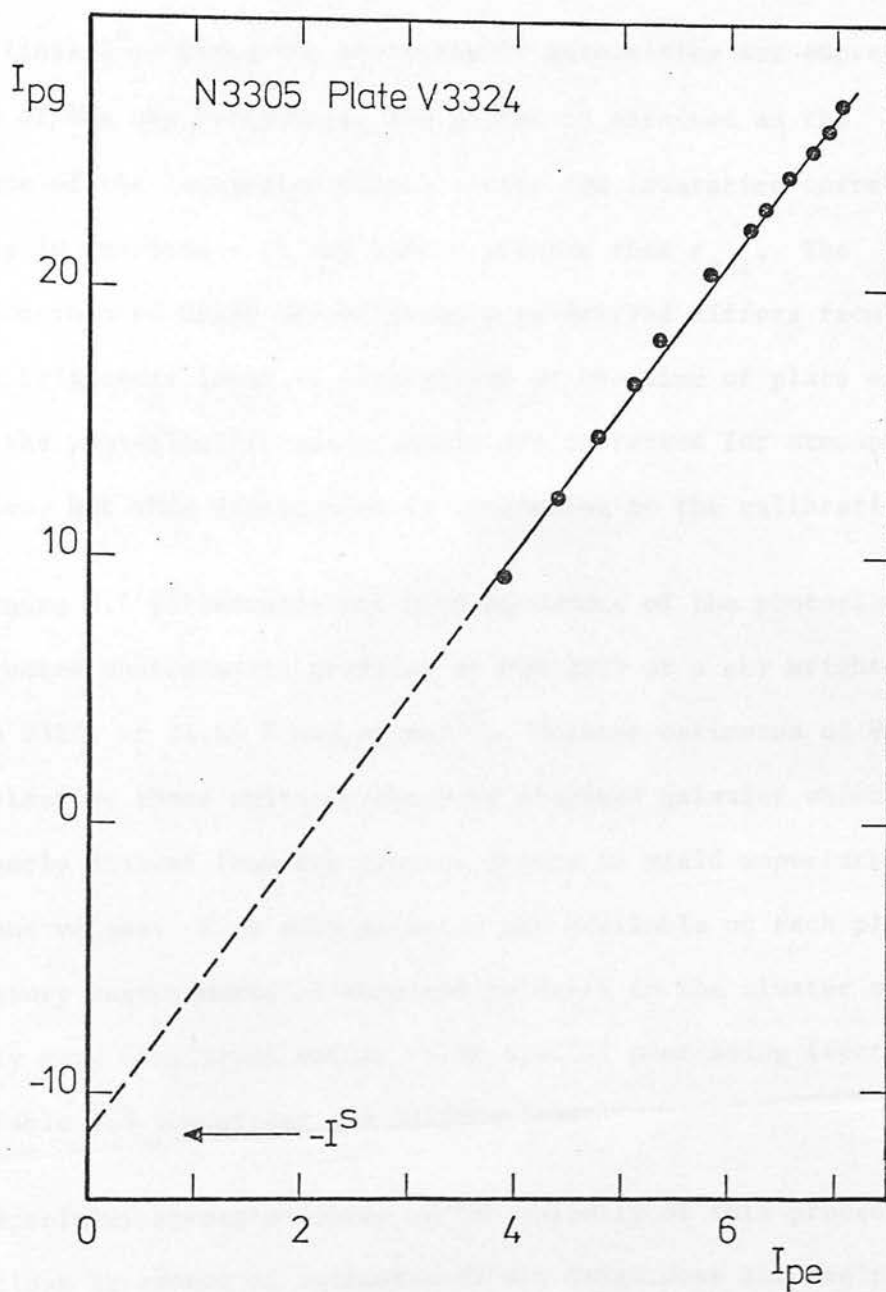


Figure 5.5

Plot of the integrated photographic intensity within simulated apertures centred on a plate image of N3305, versus corresponding photoelectric intensity (both shown in arbitrary units).  $I^S$  is a measure of the saturation of the photographic image.



which defines  $I^S$ . Since the photographic intensities are expressed in units of the sky brightness, the latter is obtained as the difference of the integrated photoelectric and saturation-corrected photographic magnitudes at any radius greater than  $r_{\text{sat}}$ . The extinction-reduced night sky brightness so derived differs from the true sky brightness local to that galaxy at the time of plate exposure because the photoelectric measurements are corrected for atmospheric extinction, but this distinction is immaterial to the calibration.

Figure 5.4 illustrates the good agreement of the photoelectric and corrected photographic profiles of NGC 3305 at a sky brightness on plate V3324 of  $21.84 \text{ V mag arcsec}^{-2}$ . Further estimates of  $V_{\text{sky}}$  are provided by those multiply-observed standard galaxies which are sufficiently distant from the cluster centre to yield unperturbed background values. Five such galaxies are available on each plate; confirmatory measurements of standard galaxies in the cluster core generally gave consistent values after special processing (section 5.5). Table 5.3 summarises the calibrations.

Especially strong evidence of the validity of this procedure is the close agreement of estimates of sky brightness for one plate that follow from independent measurements (Jobs 254 and 301). However the slope,  $m$ , of the relation in Figure 5.4 is not constant for galaxies on any one plate and empirically it is correlated with  $I^S$ , itself dependent on galaxy type (i.e. bulge to disk ratio) and magnitude (cf. Figs. 5.10-11). This would imply systematic errors in the calibrations of relative intensity, and that of plate B3028 in particular.

### 5.5 The Primary Photometry

The large angular extent of A1060 and the desirability of good resolution for the study of small images conspired to guarantee more data from a single plate than could be stored on one magnetic tape. The measurement of each plate was therefore carried out in four installments, each involving one corner of a square region of  $3.65 \text{ deg}^2$ . In order to process the images inevitably split by internal boundaries, routines were devised to enable construction of a fifth tape from various combinations of the original four.

Straightforward computer analysis of four magnetic tapes would have involved the reduction of at least 484 individual frames each of  $316^2$  pixels, too many to contemplate. Accordingly, grids of data from plate V3324 were extracted only at locations which encompassed images suitable for photometry, so disregarding the 65% of unnecessary data. Galaxies were selected for study and later recovered on the basis of coarse-mode identifications (Chapter 2), and resulting photometric parameters were added to the master card file. The limiting magnitude of the photometry is therefore determined by the completeness of the three stages of detection, visual identification and selection of the coarse mode data, and by the surface brightness and magnitude requirements imposed during processing; ideally, the latter would dominate. Eleven objects with magnitudes in the range  $16.2 < V_{25} < 17.9$ , with no corresponding CM identifications, were discovered in the course of the primary photometry. The completeness limit, after including these objects, is considered to be  $V_{25} = 17.0$ .

The practical inconvenience of this piecemeal approach was

avoided in reductions of plate B3028, at some expense in resolution, through the ninefold increase in the areal dimensions of the inner grid afforded by 3 x 3 averaging in intensity. This permitted systematic and complete processing of the region scanned on the B plate. Finally, photographic photometry was also obtained from three data tapes of plate V1330, with and without intensity averaging.

The presence of the bright stars SAO 179041 (=BS 4162,  $V = 4.88$ ) and SAO 179027, which seriously perturb the photographic sky background level local to the core of the cluster, dictated the following approach to image analysis of the central region. The geometry of Figure 5.1 was scaled up by a factor of 128/15 in absolute linear dimensions. Fully contained within the inner grid, the central images no longer affected the external background level; this simultaneously avoided the possible influence of intracluster luminosity concentrated near the cluster core. The luminosity profiles of the two bright stars were determined from diagonal cross-sections which avoided the diffraction spikes. Image analysis was then carried out as before within suitably positioned grids but with all intensity elements referred to the background polynomial determined from the giant external grid, and with both stellar luminosity profiles subtracted out of each matrix.

## 5.6 The Secondary Photometry

The secondary photometry is based on data generated by COSMOS in its coarse analysis mode. Among parameters that are output to describe every image is a measure of area, denoted  $A$ , in units of number of pixels, and a minimum transmission,  $T_{\min}$ . A well-defined empirical

relation exists between the quantity  $\log (A/T_{\min})$  and magnitude (cf. MacGillivray & Dodd 1979).

The original file of coarse mode parameters of galaxies (Chapter 2) refers to optical image characteristics in the R band, whereas V band CM parameters are required in order consistently to relate the secondary photometry to the primary  $V_{25}$  data. Extraction of the V band data was carried out by transforming the coordinates of all galaxies in the file to the system of the V plate. The identification procedure was built around a polynomial transformation routine devised by Mr J. Cooke.

The faintest 20% of the 0.15 million CM images detected on the V plate were attributable to emulsion faults such as grain clumping; all but 5% of the remainder were images of stars. The distribution of spurious images and stars is described by Poissonian statistics and these "images" have mean separation  $\langle d \rangle = (\text{surface density})^{-\frac{1}{2}} = 740\mu\text{m}$ . The root-mean-square difference of the transformed and true positions of 46 faint reference stars distributed over the 22 square degrees was  $11\mu\text{m}$  in each coordinate. The requirement that the maximum discrepancy of position should not exceed  $250\mu\text{m}$  therefore permitted reliable image pairing even for large images with ill-defined centroids; this excludes, however, the very largest objects which did not record coherently during measurement of the V plate. A negligible number of generally very faint galaxies, 4% of the total, which were identified on the R plate were not detected on the V.

In Figure 5.6 the quantity  $\log (A/T_{\min})_V$  derived from V band CM parameters is compared with the corresponding  $V_{25}$  magnitude from the

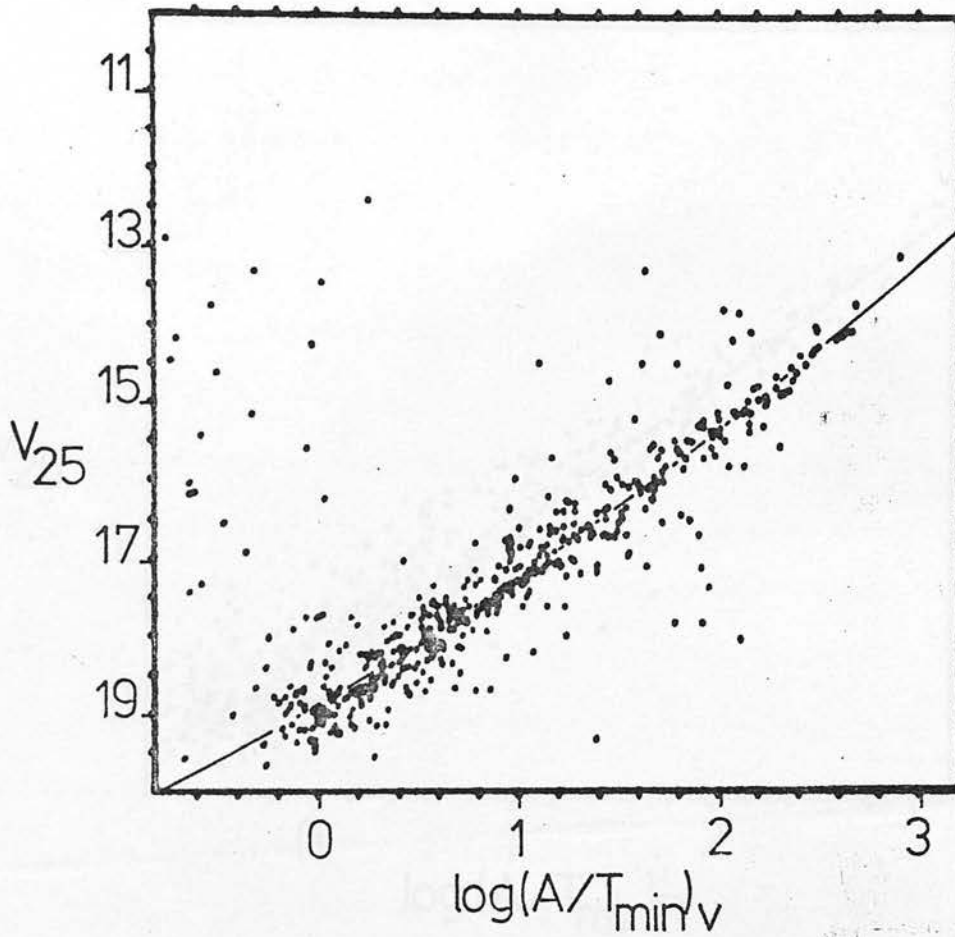


Figure 5.6

The correlation of the quantity  $\log$  (image area/minimum transmission) derived from measurements of plate V1330, with isophotal  $V_{25}$  magnitude from V3324. The curve is the adopted mean relation.



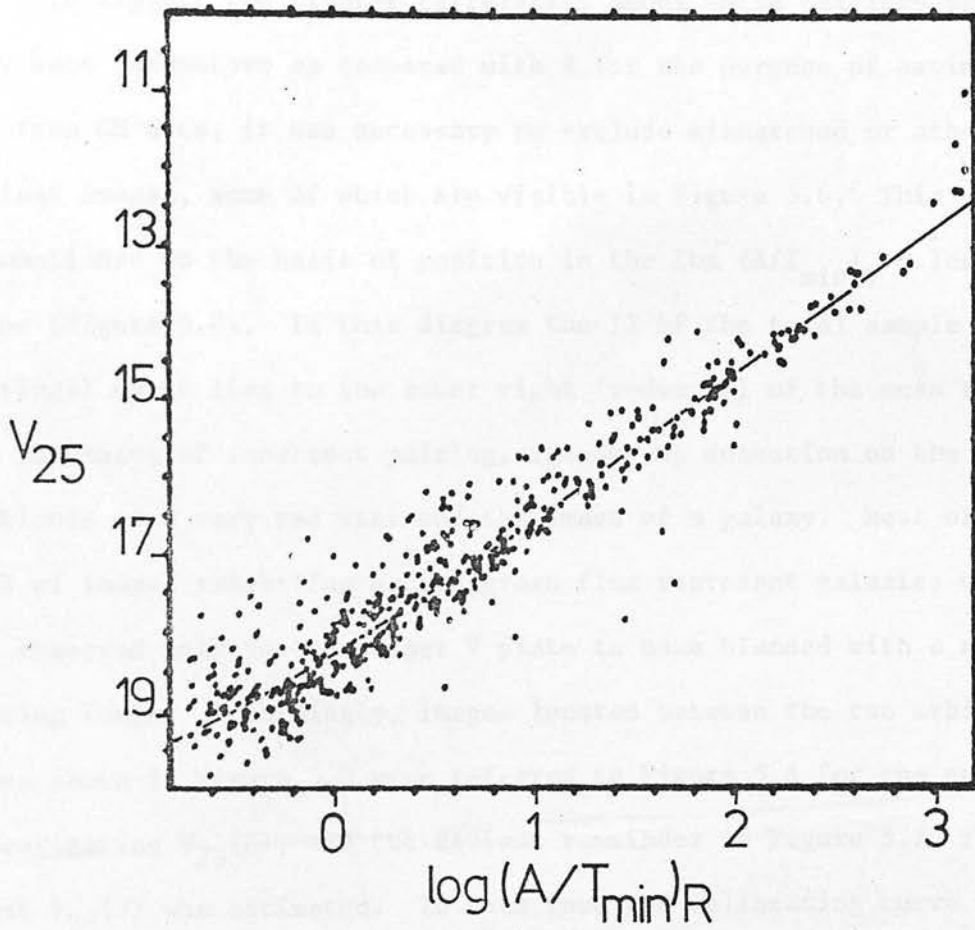


Figure 5.7

The correlation with  $V_{25}$  of the parameter  $\log(A/T_{\min})_R$ , derived from measurements of a R band plate. Compare Figure 5.6.



primary photometry, where available. Good correlation exists among valid image pairings over a range of some 6 mag in  $V_{25}$ . The equivalent diagram based on R band parameters is shown in Figure 5.7.

To exploit the tighter correlation among valid pairings provided by V band parameters as compared with R for the purpose of estimating  $V_{25}$  from CM data, it was necessary to exclude mismatched or otherwise deviant images, some of which are visible in Figure 5.6. This was accomplished on the basis of position in the  $\log(A/T_{\min})_V - \log(A/T_{\min})_R$  plane (Figure 5.8). In this diagram the 1% of the total sample (5287 pairings) which lies to the lower right (redwards) of the mean relation are instances of incorrect pairing, incoherent detection on the V plate, or blends of a very red star and the image of a galaxy. Most of the 3.3% of images exhibiting excess green flux represent galaxies which are observed only on the deeper V plate to have blended with a neighbouring image. Accordingly, images located between the two arbitrary lines shown in Figure 5.8 were referred to Figure 5.6 for the purpose of estimating  $V_{25}^{(2)}$ , and the deviant remainder to Figure 5.7, from which  $V_{25}^{(3)}$  was estimated. In each case the calibrating curve was a polynomial of suitable order in  $\log(A/T_{\min})$  fitted after iterative rejection of discordant points.

In this manner crude  $V_{25}$  magnitudes on the primary system were obtained for all but 17 images for which primary  $V_{25}^{(1)}$  photometry was not already available. The remainder of cases was attributable to hardware processing errors. Visual estimates of the magnitudes of the brighter examples were substituted where necessary and these are listed at the end of Appendix II.

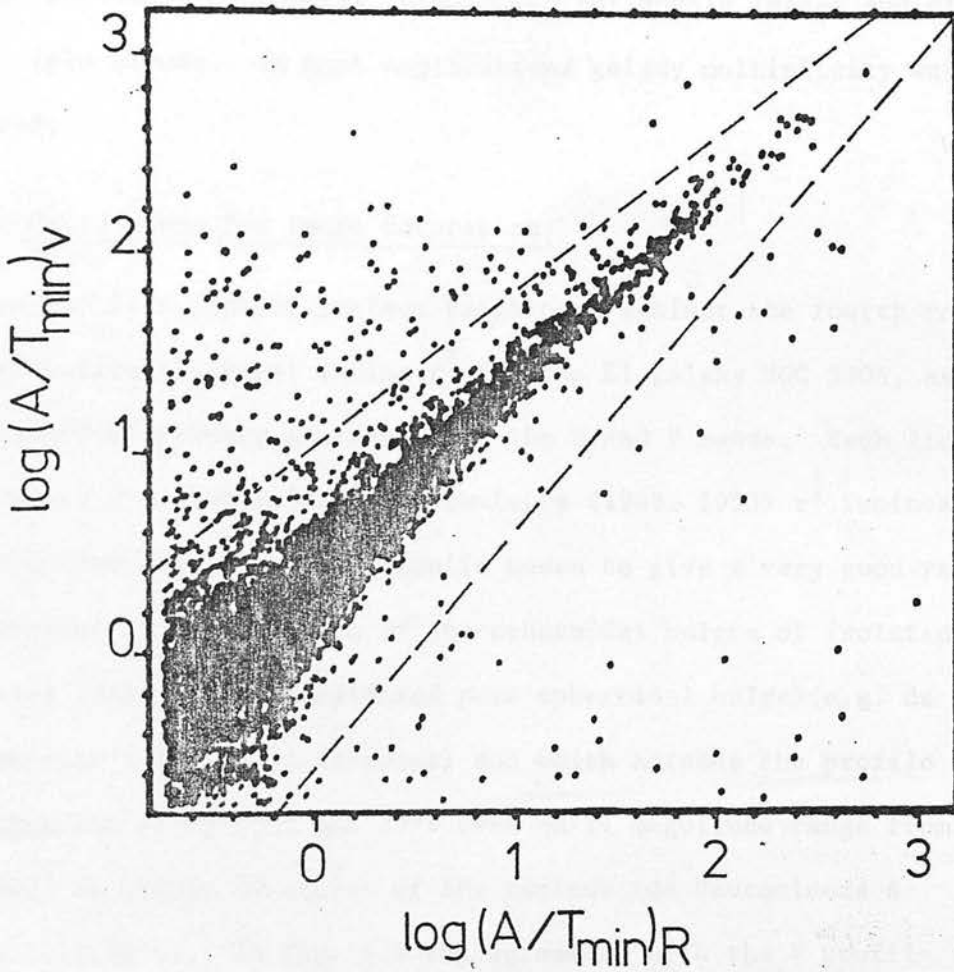


Figure 5.8

A plot of the quantity  $\log(\text{image area}/\text{minimum transmission})$  for galaxies on plate V1330 versus the same quantity measured on the R plate. Strong deviations from the mean sequence are mainly in the vertical direction and correspond to the discrepancies in Figure 5.6. The broken lines are taken to define the arbitrary limits of usefulness of the secondary V photometry.

Lastly, image blending present on both the R and (presumably) the V plate, as opposed to on the V plate only, was taken into account. It was simply assumed that each component of a blend contributed equally to the measured magnitude and a statistical correction applied to the secondary photometry, viz.  $+0.75^m$  for double images and  $+1.19^m$  for triple blends. In most applications galaxy multiplicity was ignored.

### 5.7 Corrections for Image Saturation

Figure 5.9 is a plot of surface brightness against the fourth root of the effective isophotal radius  $r^*$  for the E1 galaxy NGC 3305, as determined by primary photometry in the B and V bands. Each line represents a best-fitting de Vaucouleurs (1948, 1953)  $r^{\frac{1}{4}}$  luminosity distribution law. This relation is known to give a very good representation of the profiles of the spheroidal bulges of isolated galaxies (ellipticals considered pure spheroidal bulge) (e.g. de Vaucouleurs 1977a for references) and which matches the profile of the standard elliptical NGC 3379 over an 11 magnitude range from  $r \geq 440''$  to within 10 arcsec of the nucleus (de Vaucouleurs & Capaccioli 1979). In Fig. 5.9 the agreement with the V profile is seen to be very good to within  $r^* \approx 9''.6$  after which the core saturates on plate V3324; galaxy number 4 (E3) and others reveal the same effect. On the more lightly exposed plate V1330 (Job no. 301) the threshold surface brightness of saturation is about  $\mu_V = 20.5$  V mag arcsec<sup>-2</sup>, corresponding to a transmission T of 17. On an earlier study of the same plate (Job no. 254) this threshold was T = 33; the improvement was brought about by modifications to the COSMOS spot

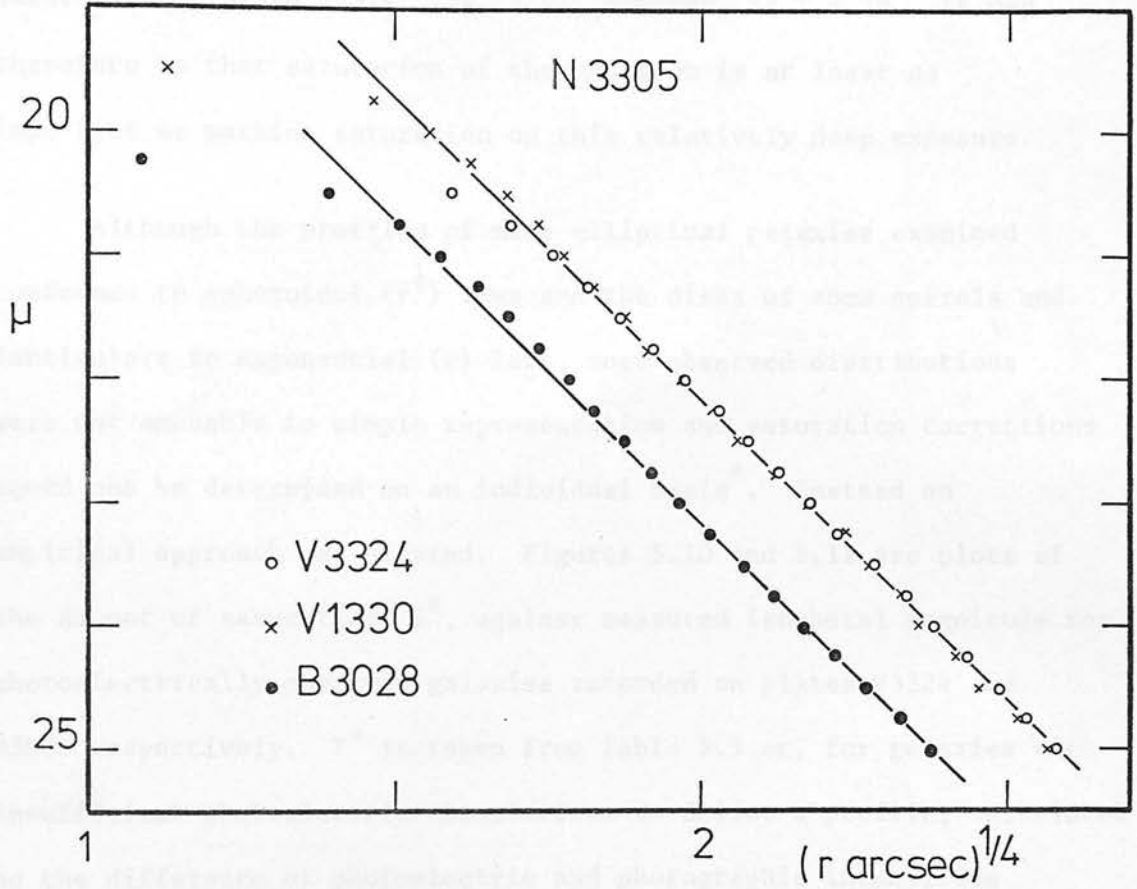


Figure 5.9

Surface brightness profiles of NGC 3305 from primary photographic photometry in the B and V bands. The  $r^{1/4}$  law in  $\mu_V$  represented by the uppermost line is fitted to the profile on plate V3324. The measured profile departs from this law shortwards of  $r^{1/4} \approx 1.76$  ( $r^* \approx 9''.6$ ).

scanning system during the intervening months. The threshold of saturation on plate V3324 (Fig. 5.9), however, is  $T = 29$ . It may therefore be that saturation of the emulsion is at least as important as machine saturation on this relatively deep exposure.

Although the profiles of many elliptical galaxies examined conformed to spheroidal ( $r^{\frac{1}{4}}$ ) laws and the disks of some spirals and lenticulars to exponential ( $r$ ) laws, most observed distributions were not amenable to simple representation and saturation corrections could not be determined on an individual basis\*. Instead an empirical approach was adopted. Figures 5.10 and 5.11 are plots of the amount of saturation,  $I^S$ , against measured isophotal magnitude for photoelectrically observed galaxies recorded on plates V3324 and B3028 respectively.  $I^S$  is taken from Table 5.3 or, for galaxies with insufficient photoelectric observations to define a profile, calculated as the difference of photoelectric and photographic intensities (presuming a mean value of sky brightness as per Table 5.3). The exact magnitude corresponding to the saturation threshold ( $I^S = 0$ ) for V3324 was taken as the median  $V_{25}$  ( $= 15.75^m$ ) of all 22 galaxies with peak surface brightnesses in the range  $20.75 \leq \mu_V \leq 21.0$ , likely examples of incipient saturation. The photoelectric photometry, unfortunately, does not go faint enough to confirm this estimate. Note the correlation of the amount of saturation with relative central concentration of light, i.e. galaxy type (bulge-to-disk ratio), at fixed magnitude. The empirical curves outlined in the figures are defined, without regard to galaxy type, in polynomial form suitable

\* The  $r^{\frac{1}{4}}$  law which best fits the profile of N3305 (Fig. 5.9) leads to metric magnitudes which consistently exceed photoelectric values by  $\sim 0.2^m$ .



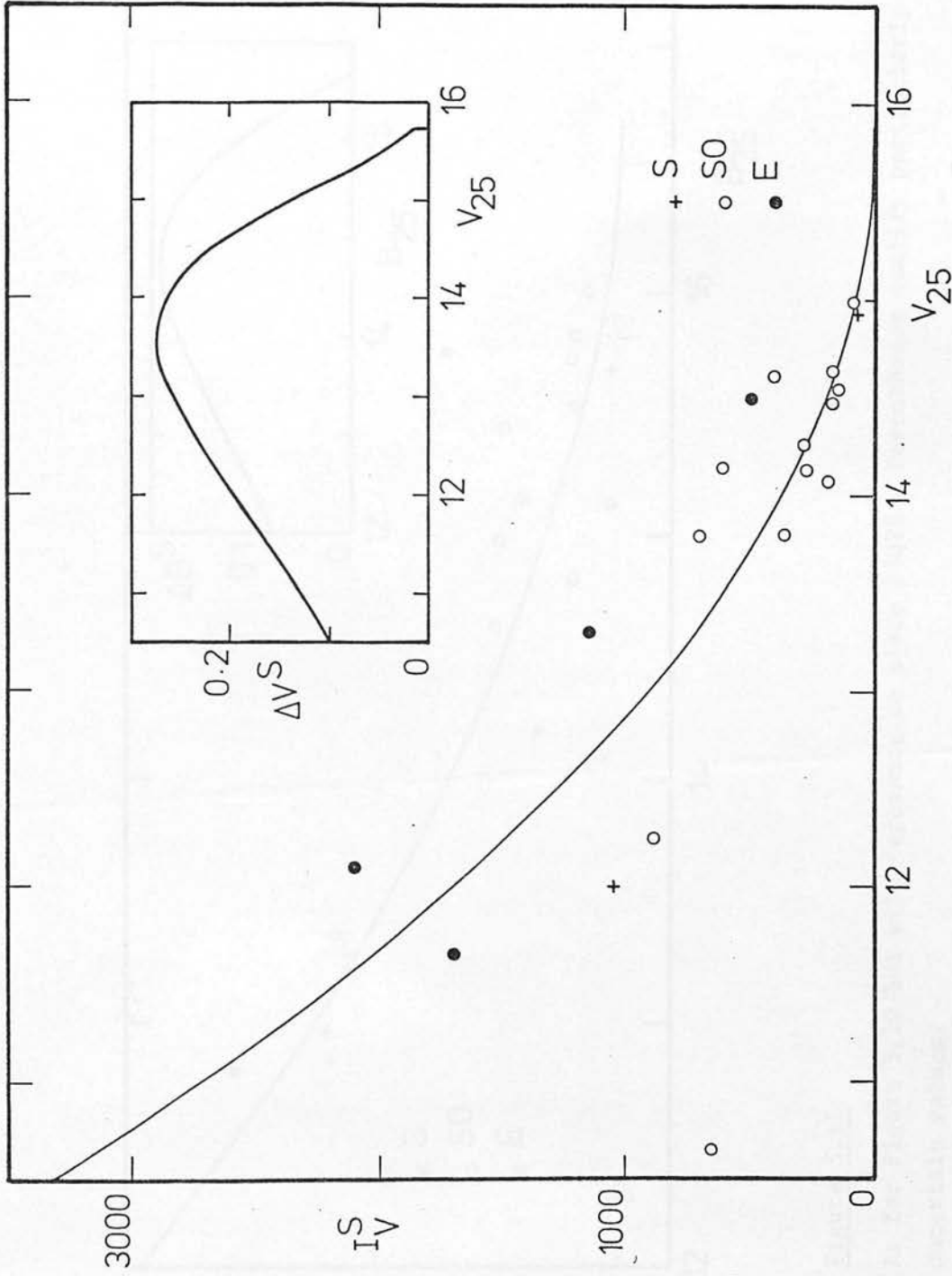


Figure 5.10

The relation between the amount of saturation in units of intensity,  $I_V^S$ , and isophotal magnitude as measured on plate V3324. Lenticular galaxies with relatively large bulges tend to lie above the mean line. The inset shows the curve of corrections to  $V_{25}$  magnitudes necessary to adjust  $V_{25}$  for saturation.



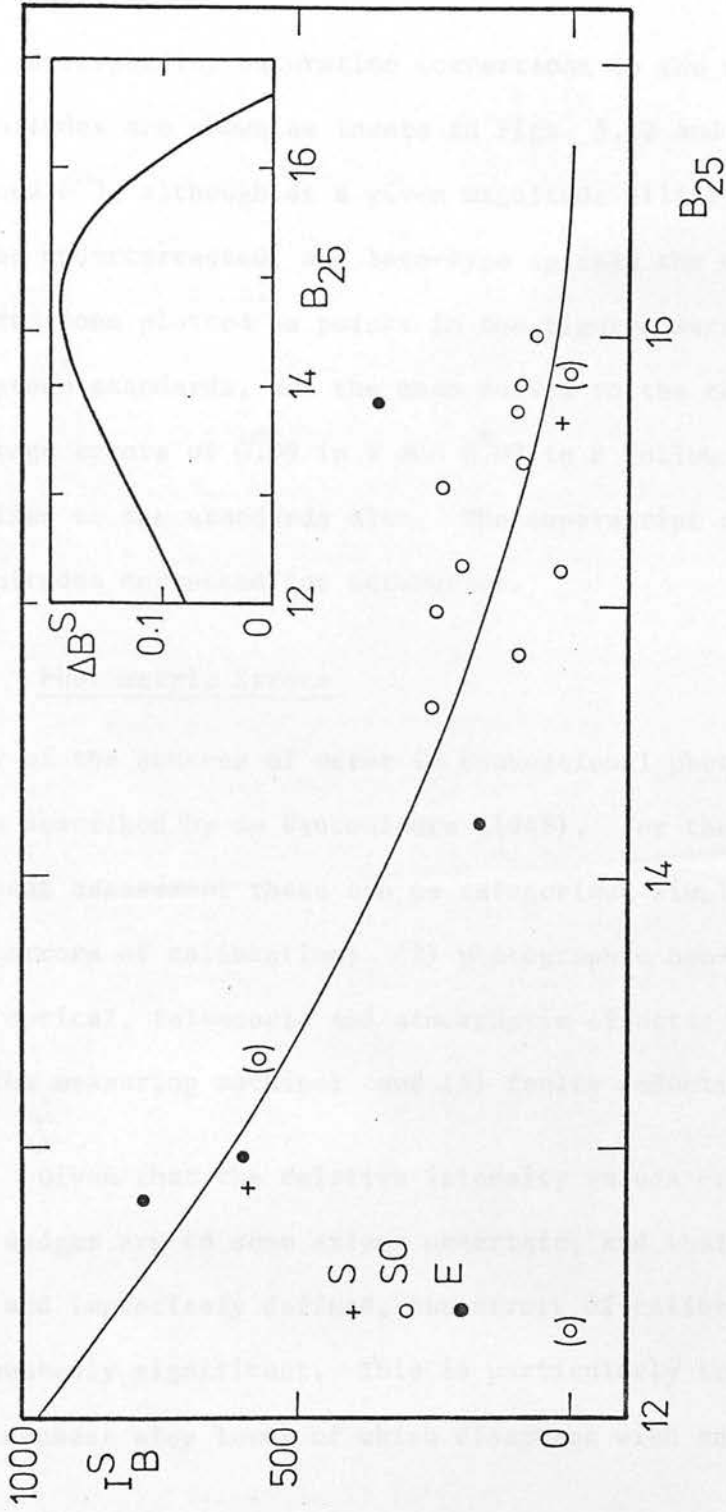


Figure 5.11

As for Figure 5.10 but with reference to plate B3028. Parentheses qualify particularly uncertain values.

for computer application:

$$I_V^s = 29884 - 3794.6 V_{25} + 120.481 V_{25}^2$$

$$I_B^s = 11437 - 1345.8 B_{25} + 39.577 B_{25}^2 .$$

The corresponding saturation corrections to the measured isophotal magnitudes are shown as insets in Figs. 5.10 and 5.11. These nowhere exceed  $0.^m3$ , although at a given magnitude elliptical galaxies tend to be undercorrected, and late-type spirals the opposite. The true corrections plotted as points in the figures were applied to the nineteen standards, and the mean curves to the rest of the photometry; average errors of  $0.^m09$  in V and  $0.^m07$  in B follow if the curves are applied to the standards also. The superscript c, as in  $V_{25}^c$ , denotes magnitudes corrected for saturation.

### 5.8 Photometric Errors

Many of the sources of error in conventional photographic photometry were described by de Vaucouleurs (1948). For the purpose of the present assessment these can be categorized simply as follows:

- (1) errors of calibration; (2) photographic non-uniformities;
- (3) optical, telescopic and atmospheric effects; (4) instabilities of the measuring machine; and (5) faulty reductions.

Given that the relative intensity values of the step levels of the wedges are to some extent uncertain, and that those levels are few and imprecisely defined, the errors of calibration involved are undoubtedly significant. This is particularly true of plate B3028, the densest step level of which disagrees with an extrapolation of

the linear Baker function defined by the rest of the wedge. The consequences are visible in Figure 5.9 as a systematic deviation of the B band radial surface brightness profile from an  $r^{\frac{1}{4}}$  law at high intensities. Unfortunately there is no independent check of the B photometry.

Emulsion non-uniformities may be manifested as variations of the natural emulsion fog across a plate, or similar variation of the transfer characteristic. IIa0 and IIaD emulsions are normally much more uniform than IIIa-J (cf. Carter & Dixon 1978). It is assumed here that the Baker functions determined for regions near the edges of the plates also apply within the scan areas, and that the background sky exposures dominate any variations of the emulsion fog.

Effects in category (3) include atmospheric haze and seeing, telescope guiding errors, and scattered light, diffraction, halation and vignetting. The unvignetted area on UK Schmidt plates is nominally  $5^{\circ}.37$  in diameter; all galaxies within 1.5 Mpc of the centre of A1060 are located within this circle on the plates studied, and only the extreme corners of the CM scan area are affected. All plates have antihalation backing and, apart from a ghost ring of radius 340 arcseconds visible at the  $24.5 \text{ V mag arcsec}^{-2}$  level around the bright central star BS4162, optical and telescopic effects can safely be neglected. Atmospheric seeing causes loss of definition and a net outward transfer of illumination in an image, but the resulting diminution of isophotal magnitude becomes appreciable only for very small images below the completeness limit of the current photometry (Hoffman & Crane 1977). Atmospheric haze on these plates was also insignificant.

There is persuasive evidence, however, of large-scale transmission gradients across the north-east and especially south-east quarters of plate V3324. Small discontinuities at boundaries of scan areas tend to confirm that drift of the photometric zeropoint occurred during measurement of this plate. Independent photometry of the worst affected region could not be secured on plate V1330 because of the loss of one tape of data. There is no suggestion of similar drift in the measurements of B3028, and scans of one step wedge before and after measurement of the latter plate reproduce identically.

Figure 5.12 presents the plots of  $\log r_{25}^*$ , the logarithms of the effective radius of the 25th V mag arcsec<sup>-2</sup> isophote, against magnitude, for all galaxies with primary photometry, on each of three plates. The scatter of points is greater for plate V1330 than for V3324 principally because of the higher level of noise on the shorter exposure. Random fluctuations about the sky level are the result of emulsion granularity and the statistics of transmission measurement, and to a first approximation the amplitude of this noise can be represented by a gaussian of standard deviation 6% about the sky level on both V1330 and B3028, and 5.5% on V3324. Since the 25 V mag arcsec<sup>-2</sup> isophote is only 5.4% above the sky level on plate V3324 and 4.7% above sky on V1330, the consequences of an error of 1-2% in the estimate of the sky level are serious (cf. Evans 1951, Fraser 1977). In fact, a 2% underestimate was found to result in a brightening of primary  $V_{25}$  magnitudes by 0.<sup>m</sup>25 on average on plate V3324. On B3028, on the other hand, the 25th B mag arcsec<sup>-2</sup> isophote is 12.8% above the sky level, and in Figure 5.12 the scatter is obviously smaller.

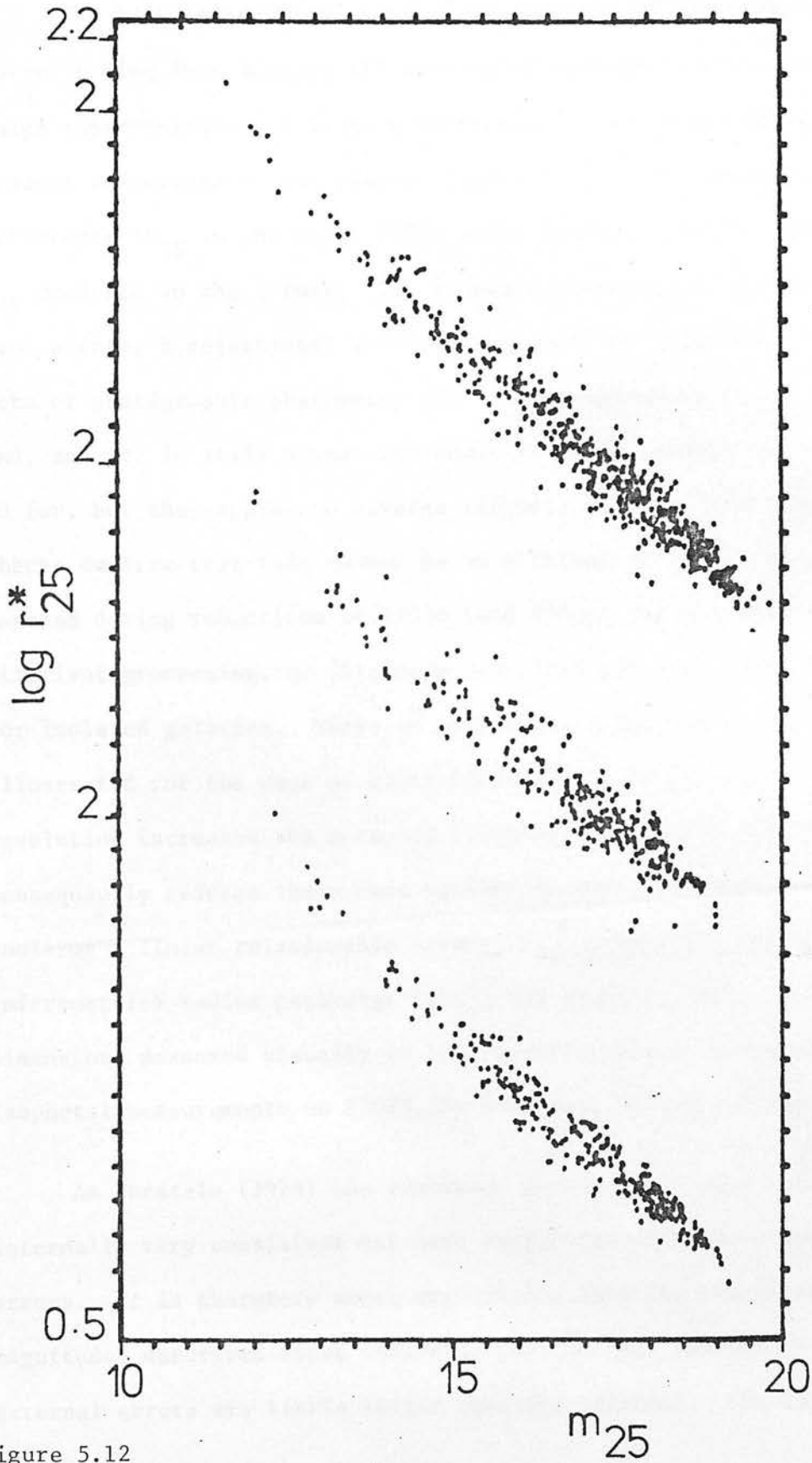


Figure 5.12

Plots of the logarithm of the isophotal radius against isophotal magnitude ( $V_{25}$  or  $B_{25}$  as appropriate) for plate V3324 (top), V1330 (centre) and B3028 (bottom). See text for discussion.

The best indicator of overall reliability of the primary photometry, taking into account all sources of systematic error other than image superposition and image saturation, is the comparison of independent reductions of two plates (Figure 5.13). The magnitude difference  $\Delta V_{25}$  in the sense V3324 minus V1330 is plotted against the  $V_{25}$  measured on the former. The standard deviation of the differences (167 points, 8 rejections) is  $0.^m18$ , dependent on magnitude. The two sets of photographic photometry are in good agreement at the bright end, and are in still better agreement if image saturation is accounted for, but they appear to diverge slightly towards faint magnitudes. Checks confirm that this cannot be an artefact of areal contraction imposed during reductions of V1330 (and B3028) for the sake of efficient processing, as this does not alter the magnitudes derived for isolated galaxies. There is nevertheless the radius effect, illustrated for the case of plate B3028 in Figure 5.14(b): reduced resolution increases the measured isophotal radii of small images and consequently reduces their mean surface brightness. Figure 5.14(a) confirms a linear relationship between  $r_{25}^*$  from V3324 and the (micrometric) radius parameter  $(ab)^{\frac{1}{2}}$ , the geometric mean of semi-axis dimensions measured visually on IIIa-J survey plates (Chapter 2). Isophotal measurements on B3028, by contrast, deviate at small radii.

As Burstein (1979) has remarked, photographic photometry can be internally very consistent but have very large systematic external errors. It is therefore worth emphasising that the comparison of  $V_{25}$  magnitudes described above indicates that in this photometry, the external errors are little larger than the internal. The latter



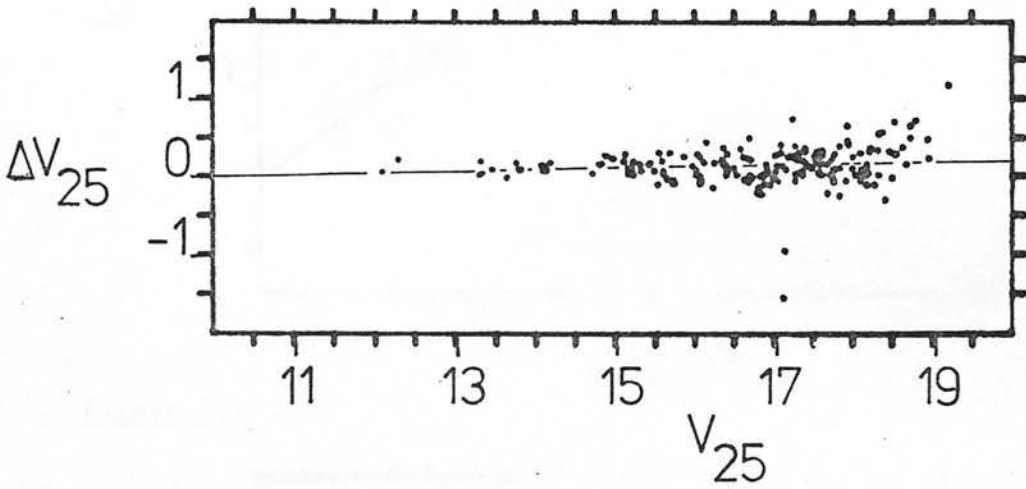


Figure 5.13

The difference of  $V_{25}$  magnitudes of galaxies as determined by primary photometry of two independent plates, plotted as a function of  $V_{25}$  measured on the deeper exposure. The latter is systematically fainter by 0.<sup>m</sup>15 at  $V_{25} = 17$ .

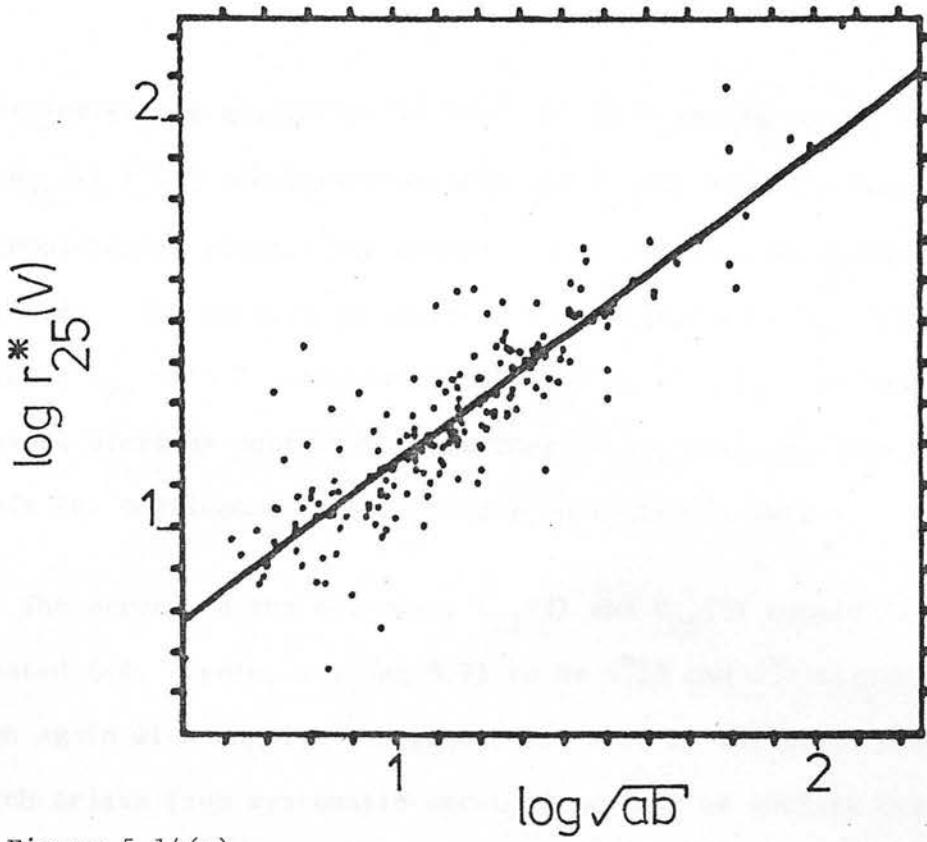


Figure 5.14(a)

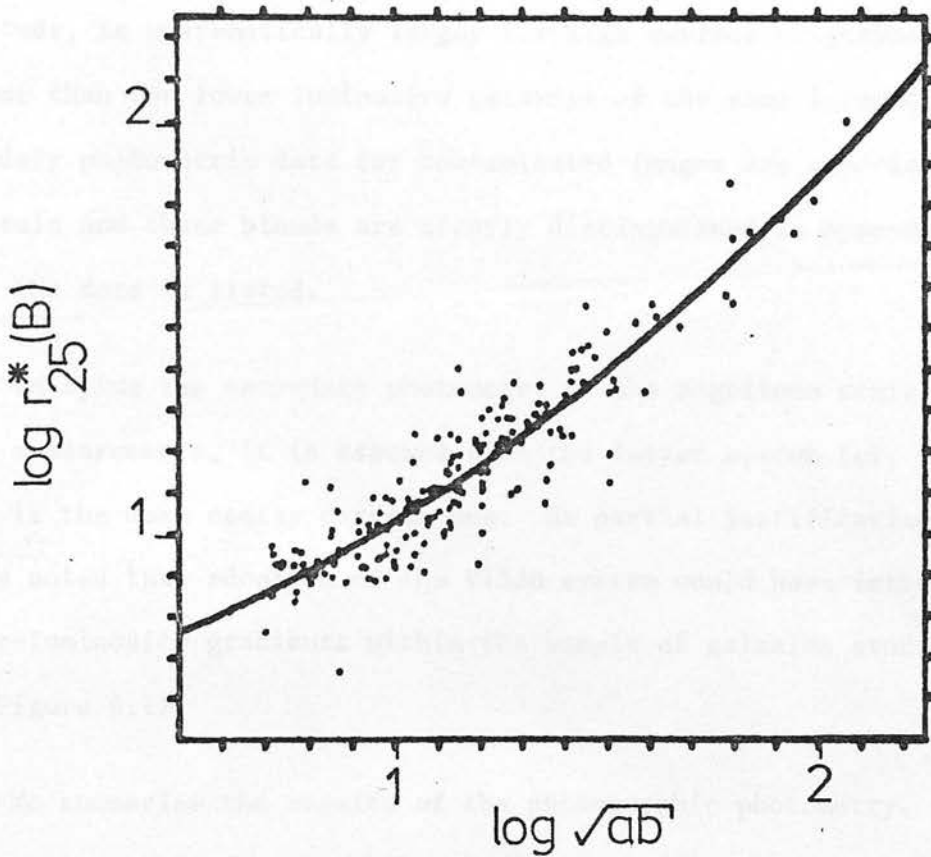


Figure 5.14(b)

Comparisons of photometric (ordinates) and micrometric (abscissae) dimensions of galaxies. Parameters  $a$  and  $b$  denote unadjusted semi-major and semiminor axis lengths in arcseconds respectively. The lower plot has been fitted by a parabola. See text for explanation.

uncertainties are generally  $\pm 0.18/2^{\frac{1}{2}} = 0.13$  in the range  $12 < V_{25}(1) < 19$ , and less than this for  $V_{25}(1)$  brighter than the completeness limit since accuracy progressively deteriorates faintwards. The saturation corrections applied where  $V_{25} \leq 15.75$  and where  $B_{25} < 16.7$  introduce uncertainties of approximately  $\pm 0.1$ , and image overlaps contribute a further  $\pm 0.2$ . Even so, there are grounds for confidence in the primary photographic data.

The errors in the secondary  $V_{25}(2)$  and  $V_{25}(3)$  magnitudes are estimated (cf. Figures 5.6 and 5.7) to be  $0.25$  and  $0.3$  respectively, though again with magnitude dependence. Part of the large scatter in each arises from systematic errors dependent on surface brightness: the measured quantity  $\log(A/T_{\min})$ , and hence the implied magnitude, is systematically larger for high surface brightness systems than for lower-luminosity galaxies of the same isophotal  $m_{25}$ . Secondary photometric data for contaminated images are especially uncertain and these blends are clearly distinguished in Appendix II, where the data is listed.

In tying the secondary photometry to the magnitude scale of the V3324 measurements, it is assumed that the latter system (cf. Figure 5.13) is the more nearly correct one. As partial justification, it may be noted that adoption of the V1330 system would have introduced colour-luminosity gradients within the sample of galaxies studied (cf. Figure 6.1).

To summarise the results of the photographic photometry, primary  $B_{25}$  magnitudes have been obtained for 266 galaxies and

$V_{25}$  (1) magnitudes for 524 galaxies located within a square region of area  $3.65 \text{ degree}^2$  centred on the Hydra I cluster. In addition, secondary estimates of  $V_{25}$  for over 5000 galaxies found in a larger region  $21.8 \text{ degree}^2$  in area are derived. The collated V band photometry is the basis of the first determination of the luminosity function of Hydra I and the surface distribution of luminosity within it, and B magnitudes provide supplementary information on the integrated colours of its members.

Table 1.2.

Plate No.	Exposure	Filter	Area	Aperture	Integration	Area covered
1229	22 May 1974	B	40	15	15	1.85
1230	22 May 1974	B	40	15	15	1.85
1231	22 May 1974	B	40	15	15	1.85
1232	22 May 1974	B	40	15	15	1.85
1233	22 May 1974	B	40	15	15	1.85
1234	22 May 1974	B	40	15	15	1.85
1235	22 May 1974	B	40	15	15	1.85
1236	22 May 1974	B	40	15	15	1.85
1237	22 May 1974	B	40	15	15	1.85

Table 5.1

## Plate Material

Plate No	3028	1330	3324	1347
Date exposed	10-11 April 1977	20-21 March 1975	5-6 July 1977	23-24 March 1975
R.A. (1950)	10 <sup>h</sup> 34 <sup>m</sup>	10 <sup>h</sup> 34 <sup>m</sup>	10 <sup>h</sup> 34 <sup>m</sup>	10 <sup>h</sup> 34 <sup>m</sup> .5
Dec. (1950)	-27° 24'	-27° 30'	-27° 24'	-27° 22'5
Filter	GG385	GG495	GG495	RG630
Emulsion	IIaO	IIaD	IIaD	098-04
Grain type	medium	medium	medium	coarse
Exposure (minutes)	20	12	60	10
Equivalent photo- electric waveband*	B	V	V	R

\* Adopted in the text as prefixes for plate numbers.

Table 5.2

## COSMOS Measurements

Plate No	Date measured	COSMOS Job No	Mode	Spot FWHM(μm)	Increment size(μm)	Area measured (sq. degrees)
3028	22 Mar 1979	333	MM	32	16	3.65
1330	31 Jan 1977	222	CM	25*	16	21.86
1330	17 Aug 1977	254	MM	25*	16	0.91
1330	19 Nov 1978	301	MM	32	16	3.65
3324	30 May 1978	287	MM	32	16	3.65
1347	3 Feb 1977	222	CM	25*	16	21.86

\* Nominally 8μm. This was replaced by a new 8μm spot scanner installed during 1978.

Table 5.3

## Zeropoint Calibrations

<u>Job No</u>	<u>Galaxy</u>	<u>m</u>	<u><math>I^S/I_{\text{sky}}</math></u>	<u><math>m_{\text{sky}}</math></u>	<u><math>\Delta</math></u>
287	4	5.67	2100	21.88	.05
287	12	5.45	1150	21.84	.01
287	22	5.48	690	21.85	.02
287	31	5.02	190	21.77	-.05
287	38	5.29	610	21.81	-.02

Mean sky brightness, plate V3324 = 21.83  $\pm$  .02 (s.e.)

254	12	4.51	780	21.66	-.03
254	22	4.75	480	21.69	.00
254	38	4.75	495	21.69	.00
301	4	5.20	1190	21.79	.10
301	12	4.59	450	21.66	-.03
301	31	4.47	50	21.63	-.06

Mean sky brightness, plate V1330 = 21.69  $\pm$  .02 (s.e.)

333	4	1.41	600	22.87	.10
333	12	1.23	170	22.72	-.05
333	22	1.36	255	22.83	.06
333	31	1.15	20	22.62	-.15
333	38	1.28	245	22.76	-.01

Mean sky brightness, plate B3028 = 22.77  $\pm$  .04 (s.e.)

Column 1: COSMOS job number; col. 2: galaxy designation;  
col. 3: slope in arbitrary units of linear regression between  
integrated photographic and photoelectric intensities (Fig. 5.5);  
col. 4: total photographic saturation in units of sky intensity  
 $\text{arcsec}^{-2}$ ; col. 5: estimated sky brightness; col. 6: the  
deviation, estimated sky less mean sky brightness.



## CHAPTER 6

### PHOTOMETRIC AND OPTICAL DATA ASSESSED

#### 6.1 Correlation of Colour with Morphological Type

Accurate photoelectric aperture photometry of galaxies (of which Table 4.3 is an unrepresentative sample) reveals a clear correlation of integrated colour with morphological type (e.g. de Vaucouleurs 1961a) due to the differing stellar compositions of galaxies of different stages within the revised Hubble classification system. Such a correlation should exist in the present photographic photometry also.

In Figure 6.1 the corrected colour indices  $B_{25}^c - V_{25}^c$  are plotted against V magnitude, for all galaxies in the field of A1060 which have photometry in each waveband and morphological classifications. Strictly, this colour index lacks rigorous physical justification inasmuch as  $B_{25}^c$  and  $V_{25}^c$  refer to light emission interior to different and unrelated isophotes; the correct procedure is to form a colour index in which the magnitude in one waveband is contrasted with the effective magnitude interior to a disk of the same dimensions in the second band (e.g. Austin, Godwin & Peach 1975). This, however, is impracticable here because of the manner of data compilation, and in either case the results are affected by image superpositions and overlaps. Figure 6.1 does reveal that on the crude definition of B-V colour index, spiral and irregular galaxies are on average  $0.25^m$  bluer than lenticulars and ellipticals of

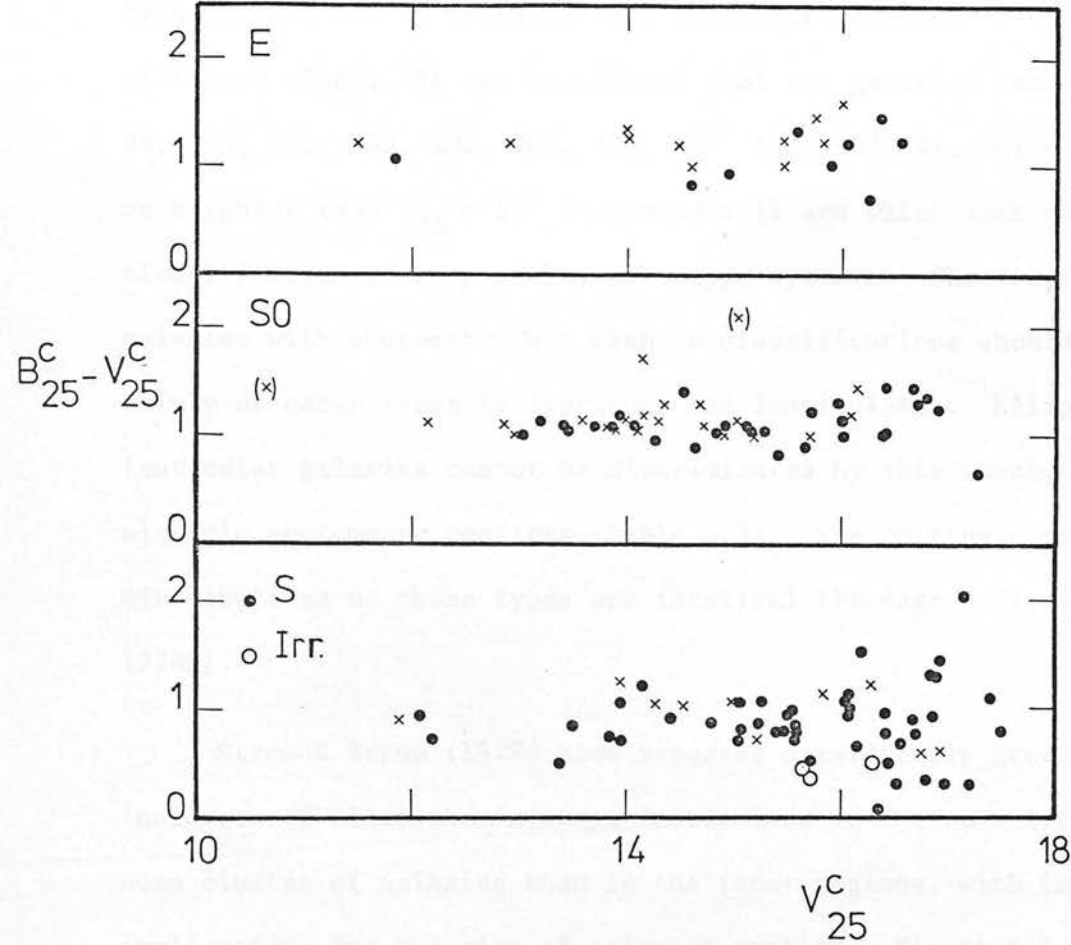


Figure 6.1

Plots of B-V colour index versus  $V_{25}$  magnitude for elliptical (E), lenticular (S0), spiral (S) and irregular (Irr) galaxies in the field of A1060. Crosses represent systems within 0.25 Mpc of the adopted cluster centre, and filled circles those beyond 0.25 Mpc, as seen in projection. Galaxies nearer to the cluster centre tend to be less blue and S and Irr galaxies more blue than other types.

Parentheses enclose less certain values.

comparable  $V_{25}$ . The latter types comprise 76% of all galaxies with measured  $B-V \geq 1.0$ , whereas 81% of those systems with  $B-V < 1.0$  are spiral or irregular (late-type\*) galaxies. Thus, in spite of significant intermingling of colour indices of the two groups, late-type systems can be distinguished with fair reliability on the basis of colour alone. It can be deduced that the galaxies numbers 14, 33, 84, 104, 115, 132, 202, 207, 324, 427, 448, and 481, which are listed as brighter than  $V_{25} = 17^m$  in Appendix II and which lack visual classifications, are probably late-type systems. The remaining galaxies with photometry but with no classifications should consist mainly of early types (ellipticals and lenticulars). Elliptical and lenticular galaxies cannot be discriminated by this means, as photo-electric photometry confirms (Table 4.3). The intrinsic colour distributions of these types are identical (Sandage & Visvanathan 1978b).

Strom & Strom (1979b) have reported a relatively greater incidence of bluer-than-average lenticulars in the outskirts of the Coma cluster of galaxies than in the inner regions, with important implications for theories of galaxy formation. Figure 6.1 suggests that the same trend exists among lenticular galaxies identified, in projection, near the Hydra I cluster. The same inference, however, would also be drawn regarding E and S systems; it is more likely that the V magnitudes determined for galaxies in the cluster core are brighter than they should be, either because of incomplete subtraction of light scattered from the two red field stars, SA0179027

\* A misappropriated term as used in the Hubble and the de Vaucouleurs classification systems. In reality, young stellar populations characterise spiral and S and Irr galaxies; cf. the Yerkes system (Morgan 1958).

and SA0179041, or because of contamination of luminosity from the halo of NGC 3311, or both. The blue magnitudes would have been less affected since they refer to brighter isophotal levels, so accounting for the net reddening of B-V. The demonstration (or otherwise) of radial segregation of lenticular galaxies in A1060 according to colour probably will be possible only if a further programme of direct photoelectric photometry, which avoids many of the large uncertainties inherent in the photographic method, is carried out.

## 6.2 The Luminosity Function

In constructing the V-band luminosity function of the cluster, attention is restricted to galaxies within the conical volume of cross-sectional diameter 3 Mpc at the assumed distance of Hydra I ( $\Delta = 40$  Mpc). A histogram of these galaxies collected in  $0.2^m$  bins represents the differential luminosity function as contaminated by non-members (Figure 6.2). This sample is dominated by faint galaxies (85% of the measured magnitudes lie in the range  $17.0 < V_{25} \leq 19.8$ ), but interest is concentrated on those brighter than the limit at which appreciable incompleteness sets in. Figure 6.3 illustrates the effect of statistically correcting for non-members assuming a homogeneous distribution of field galaxies corresponding to a surface density of  $19 \text{ galaxies degree}^{-2}$  at  $V_{25} = 17.0^m$  (Chapter 2). The dip at  $V_{25} \approx 15.8$  in the shaded histogram may well be a statistical artefact of the correction; this apart, the essential form of the function is preserved although most galaxies fainter than  $V_{25} = 16$  are assigned to the field.

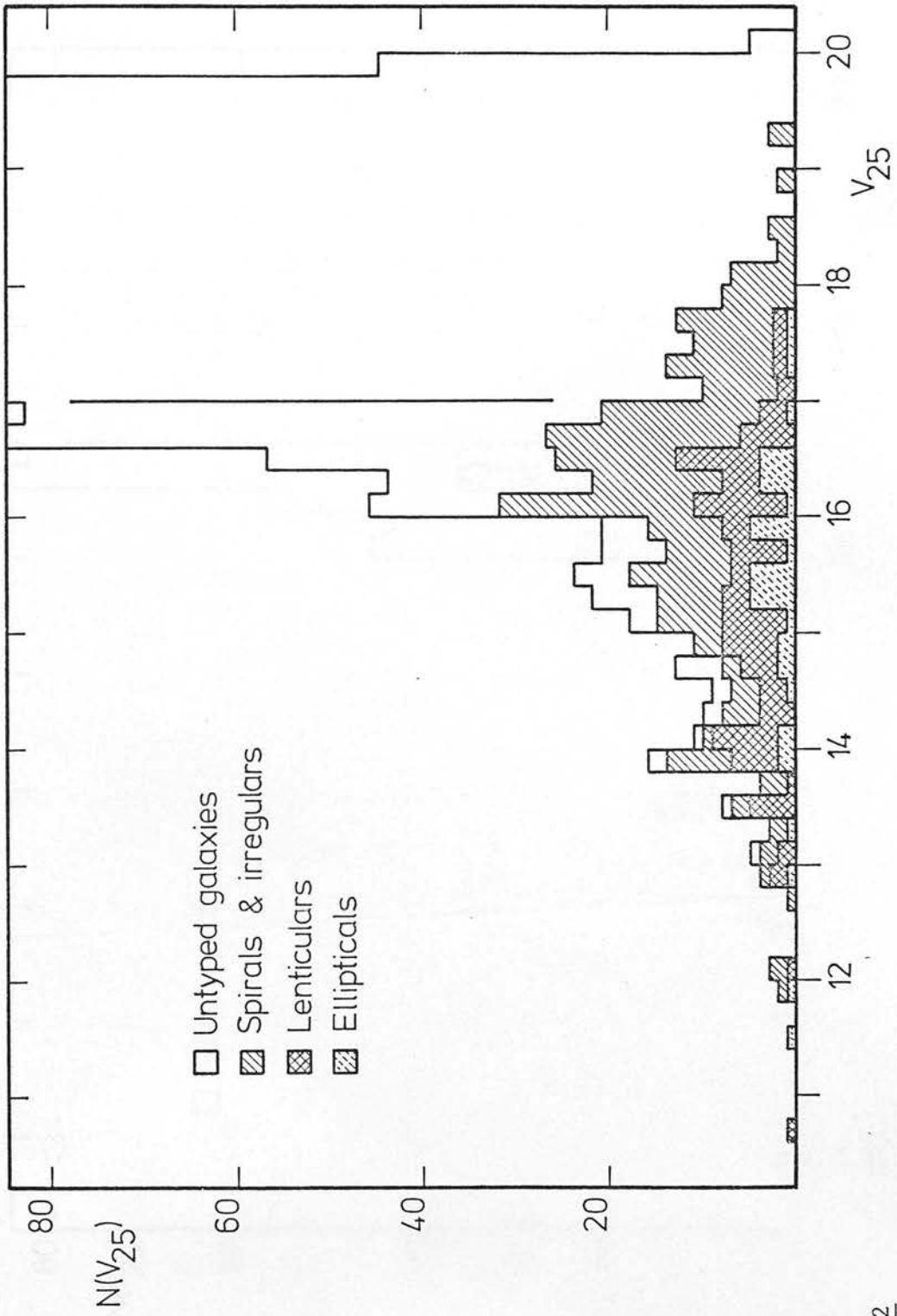


Figure 6.2

The differential luminosity function of the cluster, uncorrected for the field contribution. The vertical line marks the limit of completeness. There are 87 galaxies in  $16.6 < V_{25} \leq 16.8$ .

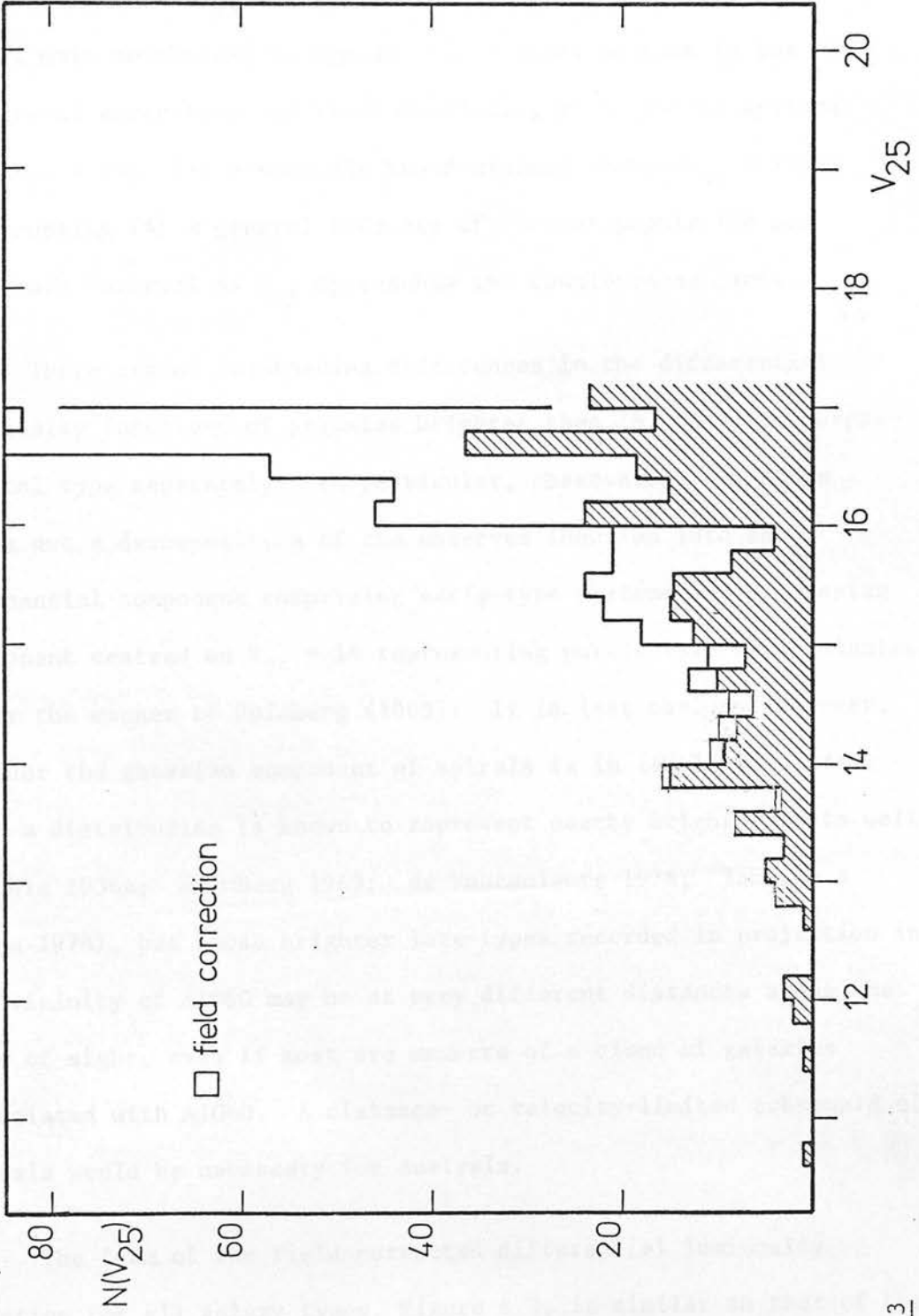


Figure 6.3

The differential luminosity function statistically corrected for the field assuming a homogeneous distribution of non-members corresponding to 19 galaxies degree<sup>-2</sup> at  $V_{25} = 17.0$ .



The basic features of the distribution are (1) a high-luminosity "toe" ( $V_{25} \lesssim 12.2$ ) consisting of bright representatives of all main morphological types; (2) a local maximum in the numbers of early-type and (less convincingly) late-type systems near  $V_{25} = 14$ ; (3) a possible broad minimum around  $V_{25} \approx 14.5$ , interrupting (4) a general increase of cluster population per magnitude interval as  $V_{25}$  approaches the completeness limit.

There are no outstanding differences in the differential luminosity functions of galaxies brighter than  $16^m$  for each morphological type separately. In particular, observation (2) above rules out a decomposition of the observed function into an exponential component comprising early-type systems and a gaussian component centred on  $V_{25} = 14$  representing purely late-type galaxies, after the manner of Holmberg (1969). It is less obvious, however, whether the gaussian component of spirals is in itself precluded. Such a distribution is known to represent nearby bright spirals well (Hubble 1936a; Holmberg 1969; de Vaucouleurs 1974; Tammann & Kraan 1978), but those brighter late-types recorded in projection in the vicinity of A1060 may be at very different distances along the line of sight, even if most are members of a cloud of galaxies associated with A1060. A distance- or velocity-limited subsample of spirals would be necessary for analysis.

The form of the field-corrected differential luminosity function for all galaxy types, Figure 6.3, is similar to that of the Coma cluster (Godwin & Peach 1977), which exhibits a pronounced minimum following a local peak in numbers of galaxies of intermediate

luminosity. As Godwin & Peach point out, these features cannot be accounted for on any smoothly rising analytic representation of the function. The correspondence between the cluster functions appears to be close despite the relative deficiency of elliptical galaxies in Hydra I and the much higher proportion of spirals.

The logarithmic cumulative luminosity function of A1060 before and after adjustment for the inclusion of non-members is shown in Figure 6.4. Except at the bright end, the corrected data is well represented by two straight lines of slopes  $0.60 \pm 0.05$  and  $0.235 \pm 0.005$ . The intersection at  $V_{25}^* = 14.08$  corresponds to the local maximum in the differential luminosity function.

Abell (1962, 1965, 1975) has obtained adequate bilinear representations of the luminosity functions of several clusters in the manner of Figure 6.4, and has suggested (Abell 1972; Bautz & Abell 1973) that the discontinuity or "knee" in the functions occurs at a similar absolute magnitude in all clusters. Austin, Godwin & Peach (1975) quote a mean absolute dispersion of only  $0.^m15$  in the position of the knee in the functions of nine clusters, which supports the impression of universality. The observed  $V_{25}^*$  of Hydra I is readily placed on their system by the prescription  $H_0 = 50 \text{ km s}^{-1} \text{ Mpc}^{-1}$ ,  $q_0 = +1$ , K-correction  $K_V = 0.^m02$  (cf. Sandage 1973b), galactic absorption  $A_V = 0.^m13$  (Sandage 1973a), redshift  $z(\text{A1060}) = 0.0113$  (Chapter 3), and a correction of  $-0.^m2$  to convert  $V_{25}^*$  to a total magnitude. The result is  $M_{V_{25}}^* = -20.42$ , which differs markedly from the average  $\langle M_{V_{25}}^* \rangle = -21.16 \pm 0.15$  reported by Austin, Godwin &

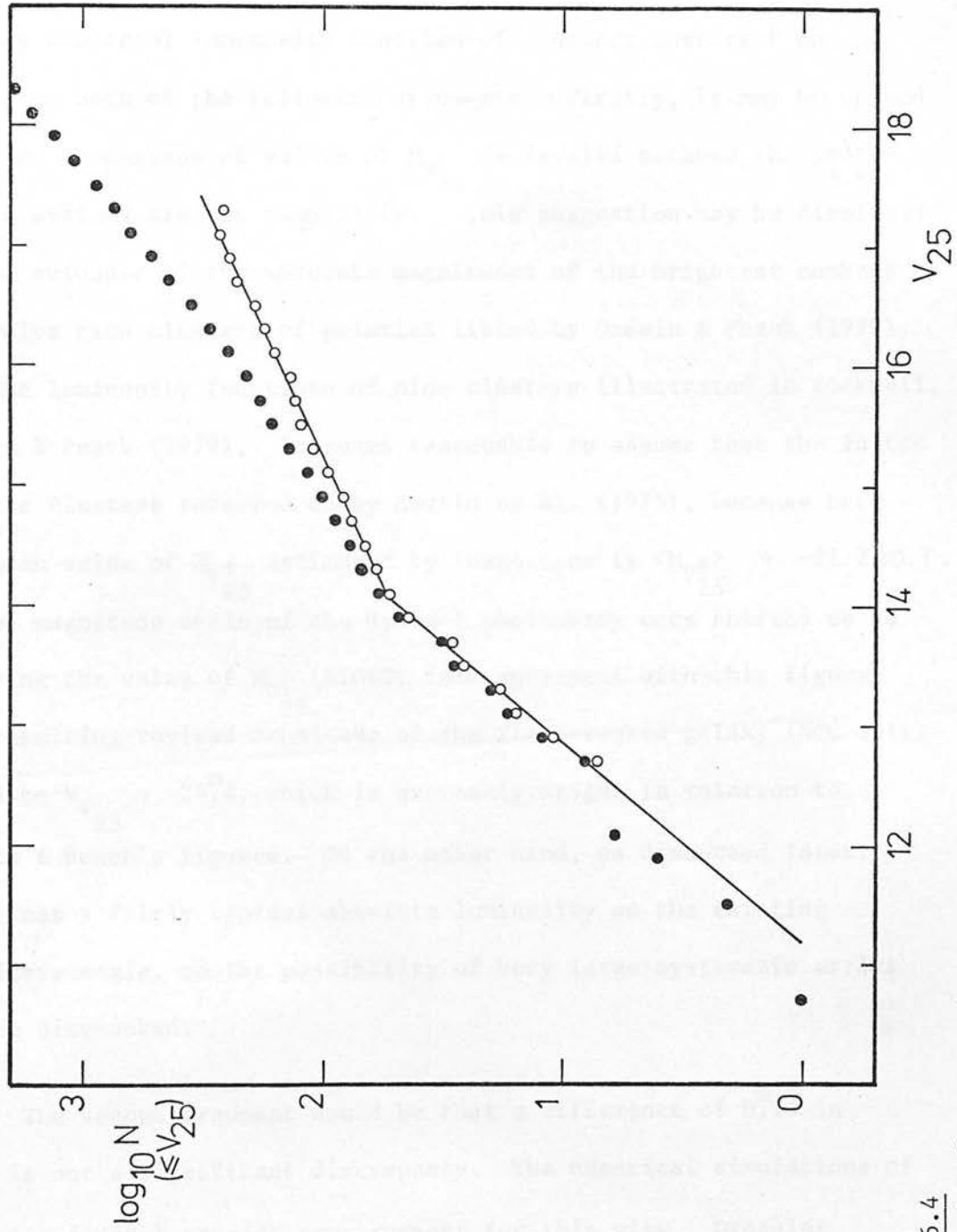


Figure 6.4

The logarithmic cumulative luminosity function of A1060. Filled circles are totals uncorrected for the field, open circles are corrected totals. The two lines which best represent the corrected data faintwards of  $V_{25} = 12.5$  are illustrated.

Peach (1975). An exploratory reduction of the size of the A1060 sample, by changing the radius of the projected boundary from 1.5 Mpc to 0.6 Mpc, did not alter  $V_{25}^*$  or the corrected slopes of the function.

Any attempt to reconcile the above figures with the notion of a truly universal luminosity function of clusters must rest on either or both of the following arguments. Firstly, it may be opined that the comparison of values of  $M_{V_{25}}^*$  is invalid because the photometric systems are not compatible. This suggestion may be dismissed on the evidence of the absolute magnitudes of the brightest members of twelve rich clusters of galaxies listed by Godwin & Peach (1979), and the luminosity functions of nine clusters illustrated in Bucknell, Godwin & Peach (1979). It seems reasonable to assume that the latter are the clusters referred to by Austin et al. (1975), because here the mean value of  $M_{V_{25}}^*$  estimated by inspection is  $\langle M_{V_{25}}^* \rangle = -21.2 \pm 0.1$ . If the magnitude scale of the Hydra I photometry were shifted so as to bring the value of  $M_{V_{25}}^*$  (A1060) into agreement with this figure, the resulting revised magnitude of the first-ranked galaxy (NGC 3311) would be  $M_{V_{25}}^m = -24.4$ , which is extremely bright in relation to Godwin & Peach's figures. On the other hand, as discussed later, NGC 3311 has a fairly typical absolute luminosity on the existing magnitude scale, so the possibility of very large systematic errors can be discounted.

The second argument would be that a difference of  $0.74^m$  in  $M_{V_{25}}^*$  is not a significant discrepancy. The numerical simulations of Dressler (1978a) provide some support for this view. Dressler constructed several thousand model clusters each representing, in

effect, a random selection or "statistical fluctuation" from an infinite population of galaxies conforming exactly to the Schechter (1976) formulation of a universal luminosity function. He found many instances in which the characteristic magnitude  $M^*$  estimated by analysis of a model with the assumed population of an Abell richness class 0 or 1 cluster differed by more than  $0.5^m$  from that used to define the Schechter function on which the model was based. These deviations tended to be much smaller in richer clusters. Although  $M^*$  in Schechter's definition is not equivalent to the discontinuity in Abell's integrated function, statistically they should behave in much the same way. Since A1060 is a richness class 1 cluster whereas six of the aforementioned nine clusters have richnesses  $R \geq 2$ , the value of  $M_{V_{25}}^*$  for A1060 does not contradict the universal-function hypothesis.

A more severe challenge to the idea of universality is the marked inconsistency of the slopes of the bright and faint ends of the luminosity functions constructed by Bucknell et al. (1979) and Dressler (1978b). Among the former sample, there exists at least one cluster which bears only crude resemblance to an Abell function. In the case of A1060 (Figure 6.4) the gradient of the bright-end segment is less steep than the average slope, 0.75, determined by Abell (1975) for four different clusters, and agreement would deteriorate further if the brightest members of Hydra I were afforded due weight in the line fitting. These misgivings acknowledged, it is nonetheless striking how well a linear approximation does fit the data over a range of nearly three magnitudes in  $V_{25}$ , from the position

of the knee at  $14.1^m$  down to the completeness limit of the photographic photometry. The confirmation that the numbers of galaxies in the cluster increase exponentially to this limit is significant, and consideration is given in the next chapter as to whether this relation continues to still fainter limits. Linear extrapolation of Figure 6.4 to  $V_{25} = \infty$  would imply that 90% of the galaxian luminosity contained within the 25th  $V$  mag arcsec $^{-2}$  isophote is represented by systems brighter than  $V_{25} = 17$ .

Introducing the notation  $M_n$  for the absolute  $V_{25}$  magnitude of the  $n$ -th brightest cluster member, and  $M^*$  for that magnitude corresponding to the knee of the integrated function, the difference  $M_1 - M^*$  for A1060 is  $3.3^m$ , which is larger than that of any of the clusters in the sample of Bucknell et al. except possibly A426. The dependence, if any, of this differential on cluster richness and composition is very relevant in relation to the use of both  $M_1$  (strictly, the equivalent metric magnitude) (e.g. Sandage 1973b) and  $M^*$  (e.g. Abell 1978) as standard candles for the determination of distance, and it may also be intimately related to secular changes in the distribution of luminosity in a cluster (e.g. Ostriker 1977). Coincidentally, A426 (the Perseus cluster) is the only other distance class 0 association in the Abell (1958) catalogue, though it is more populous and flattened in outline and possesses two unusual characteristics: an unique radial distribution of members of different morphological types (Melnick & Sargent 1977), and a very luminous first-ranked Seyfert galaxy. This latter fact is partly responsible for the large  $M_1 - M^*$  of A426, and from an uncompromising viewpoint it



would seem surprising that the  $M_1 - M^*$  of A1060 is similar. The statistical uncertainties demonstrated by Dressler (1978a) would, however, be compounded with photometric errors affecting  $M_1$ , rendering  $M_1 - M^*$  liable to large illusory fluctuations from cluster to cluster. For possible underlying trends to be revealed clearly, photographic photometry of still more clusters is required.

### 6.3 Optical Properties of the Brightest Member

The remarkable contrast of the luminosity distributions of the two brightest, centrally located galaxies in Hydra I (e.g. van den Bergh 1977b; Disney & Wall 1977; Fairall 1979) is well shown by an iso-transmission contour plot of the core of the cluster (Figure 6.5). This chart was generated from COSMOS measurements of the CTIO 4m prime focus plate number 995 detailed in Table 7.1.

The dominant galaxy, NGC 3311, forms an optical binary with the second-ranked elliptical NGC 3309 which north precedes it by 100 arcseconds. The luminosity distributions of the two galaxies appear to overlap from a level much brighter than the 25 V mag arcsec<sup>-2</sup> limiting isophote of the photographic photometry, but in view of the lack of serious distortion of N3309, this is thought to be mainly an effect of projection.

Assessment of Figure 6.5 is complicated by the fact that many extraneous sources of luminosity contribute to and influence the visible isophotes (see below). Nonetheless, the relatively great extent of the halo of N3311 is obvious. Inspection also shown that the iso-transmission contours representing successively fainter light levels

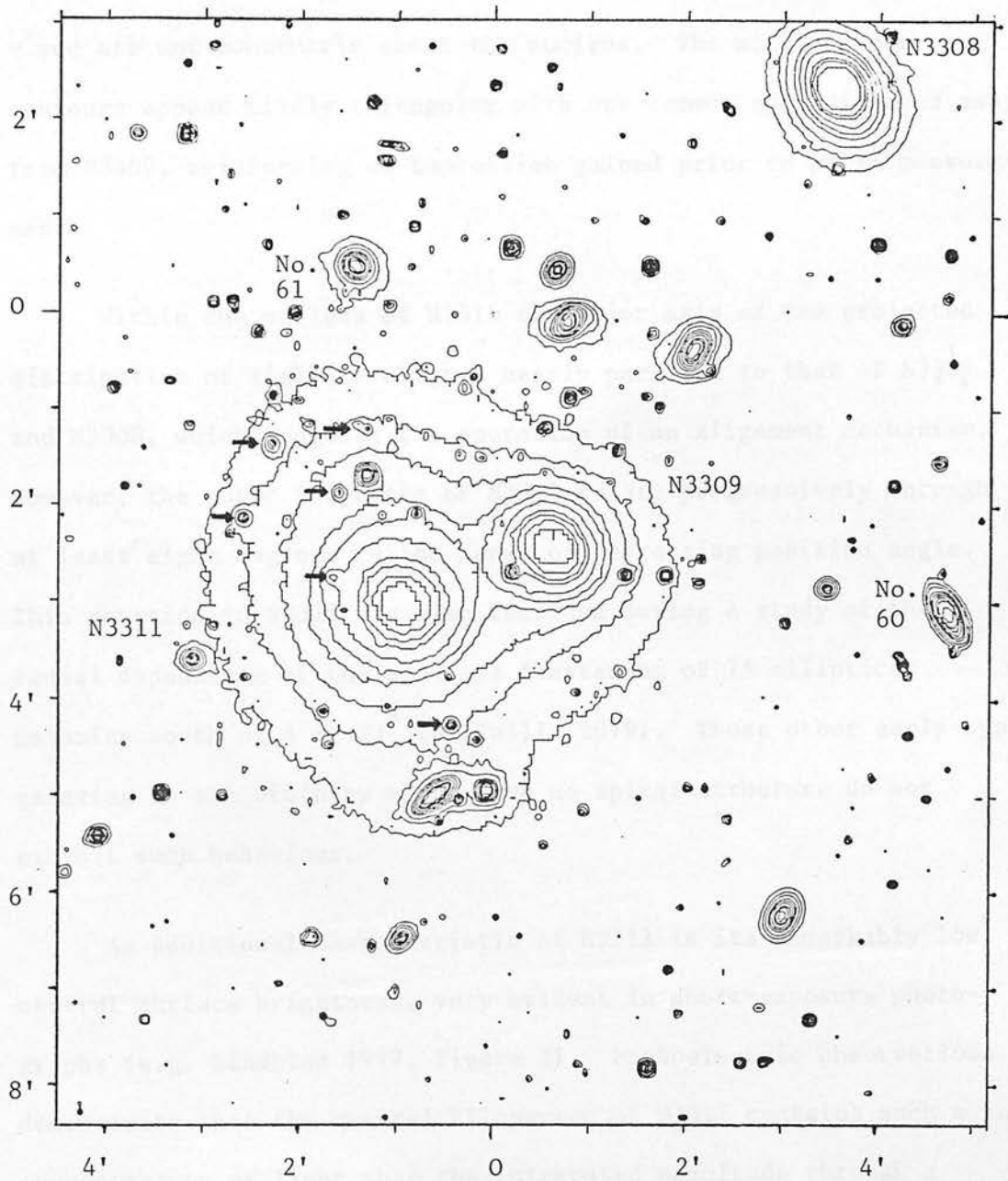


Figure 6.5

An isotransmission contour plot of the core of A1060. N3311 (left) and N3309 form the central binary, and N3308 is at the top right. The horizontal arrows identify six faint galaxies visible in projection in the halo of N3311.

within N3311 do not take the form of simple ellipses - compare N3308 - and are not concentric about the nucleus. The middle-order contours appear mildly triangular with one common apex directed away from N3309, reinforcing an impression gained prior to plate measurement.

Within the nucleus of N3311 the major axis of the projected distribution of light is aligned nearly parallel to that of N3309 and N3308, which suggests the operation of an alignment mechanism. However, the outer isophotes of N3311 rotate progressively through at least eight degrees in the sense of decreasing position angle. This rotation in N3311 was also recorded during a study of the radial dependence of the apparent flattening of 75 elliptical galaxies south of  $\delta = -23^{\circ}$  (di Tullio 1979). Those other early-type galaxies in the vicinity which have no spiral structure do not exhibit such behaviour.

An additional characteristic of N3311 is its remarkably low central surface brightness, very evident in short-exposure photographs (e.g. Lindblad 1977, Figure 2). Photoelectric observations demonstrate that the central kiloparsec of N3311 contains such a low concentration of light that the integrated magnitude through a 5-arcsecond aperture at the nucleus is  $\sim 1.5^m$  fainter than the corresponding nuclear magnitude of N3309 (van den Bergh 1977b), although no doubt part of the effect can be attributed to the absorption lane or dark cloud discovered at the very centre of N3311 by Lindblad (1977) and his collaborators.

The radial surface brightness profile of N3311 in the B and V wavebands has been derived from photographic surface photometry of plates V3324 and B3028 rather than the IIIa-J exposure depicted in part in Figure 6.5. Such 4m telescope plates suffer from severe field non-uniformities (see: Harris & Smith 1976; Rose 1979) which here are impossible to disentangle from the very extensive luminosity gradients associated with the many bright images in the small field. On the more uniform Schmidt plates only image overlaps require treatment. Cross-sections of selected galactic nuclei were taken at intervals of  $\pi/4$  in position angle within grids of intensity values from which both the profiles of stars BS4162 and SA0179027 and the sky background had been subtracted (cf. Section 5.5). The mean B and V profiles of N3311 corrected for superpositions of discrete sources are shown in Figure 6.6. At high intensities (filled circles) the curves are defined by the results of numerical processing, and by cross-sections at fainter levels (open circles). The straight line at the upper right of the figure represents the critical inverse-square decline of projected intensity with radius,  $r$ , that must be attained or exceeded as  $r \rightarrow \infty$  if the integrated luminosity is to converge. The observed gradients are shallower than this out to about 60 kpc; on the crude representation  $\text{Intensity} \propto r^{-\alpha}$ , the exponent  $\alpha$  takes the value 1.55 over the interval  $10 \lesssim r \lesssim 40$  kpc in the V band and 1.85 over  $16 \lesssim r \lesssim 70$  kpc of the B profile. Clearly  $\alpha < 2$ , though the difference between the two values of  $\alpha$  is not significant (photoelectric aperture photometry reveals that it is the V profile which is steeper). The gradient beyond 50 kpc is ill-defined on account of limited sampling, superposed images, deteriorating

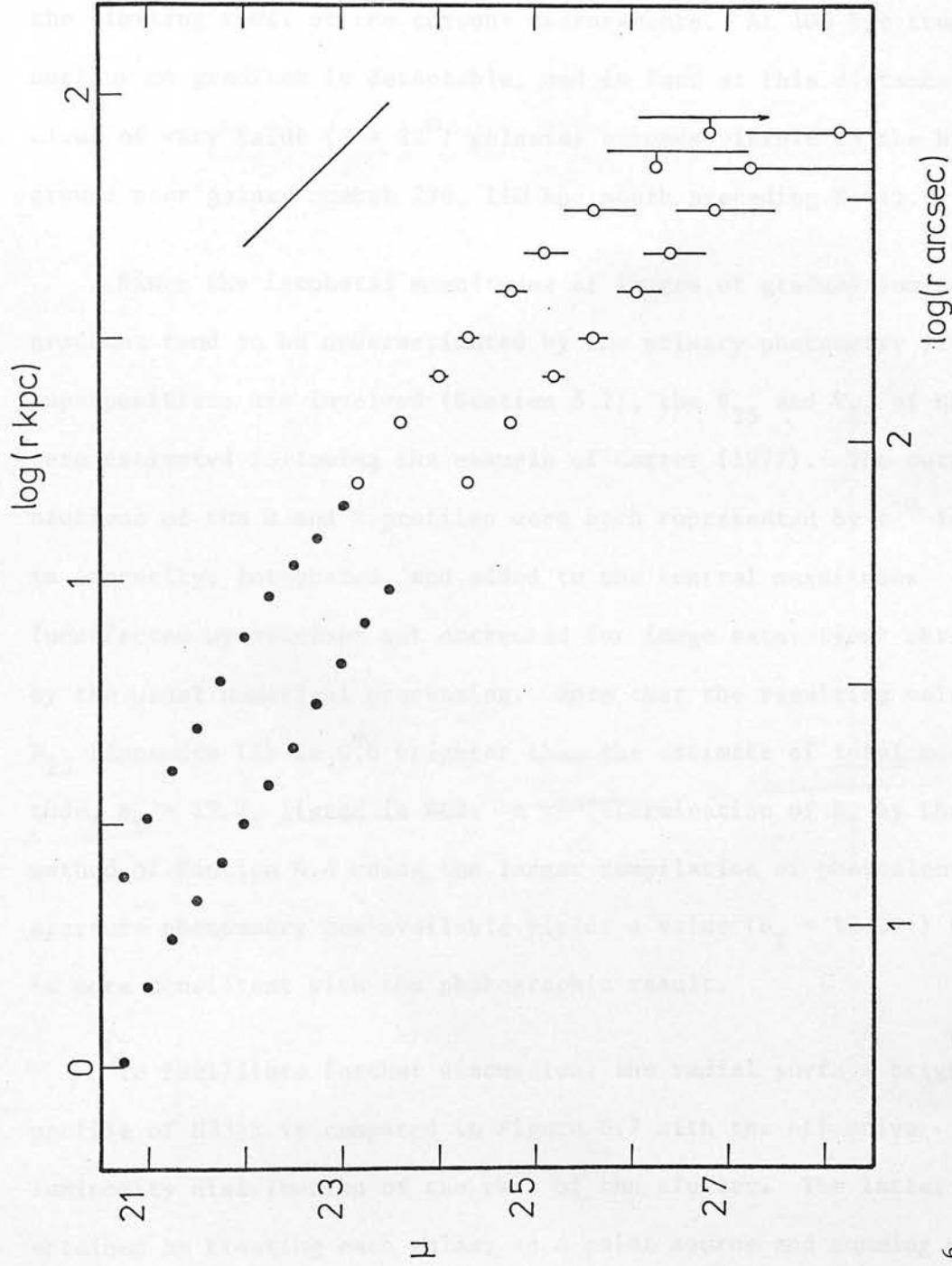


Figure 6.6

The surface brightness profiles of N3311 in the V (upper) and B wavebands. The straight line represents an inverse-square dependence of projected intensity with radius. The error bars correspond to an uncertainty of  $\pm 1\%$  in the subtracted background level.



signal-to-noise ratios and the limitations of the background modelling. Nevertheless, the indications are that the condition for convergence of total luminosity is met - that is,  $\alpha > 2$  - near the limiting radii of the current measurements. At 100 kpc from the nucleus no gradient is detectable, and in fact at this distance a cloud of very faint ( $J \approx 22^m$ ) galaxies becomes visible in the background near galaxy number 296, 110 kpc south preceding N3311.

Since the isophotal magnitudes of images of gradual luminosity gradient tend to be underestimated by the primary photometry if superpositions are involved (Section 5.2), the  $B_{25}$  and  $V_{25}$  of N3311 were estimated following the example of Carter (1977). The outer sections of the B and V profiles were each represented by  $r^{-\alpha}$  laws in intensity, integrated, and added to the central magnitudes (unaffected by overlaps but corrected for image saturation) obtained by the usual numerical processing. Note that the resulting value of  $B_{25}$  (Appendix II) is  $0.6^m$  brighter than the estimate of total magnitude,  $B_T = 12.7$ , listed in RC2. A re-determination of  $B_T$  by the method of Section 4.4 using the larger compilation of photoelectric aperture photometry now available yields a value ( $B_T = 12.06$ ;) that is more consistent with the photographic result.

To facilitate further discussion, the radial surface brightness profile of N3311 is compared in Figure 6.7 with the effective luminosity distribution of the rest of the cluster. The latter was obtained by treating each galaxy as a point source and summing within rings centred on N3311 (itself excluded) the isophotal luminosities of all galaxies brighter than  $V_{25} = 17$ , with appropriate adjustment



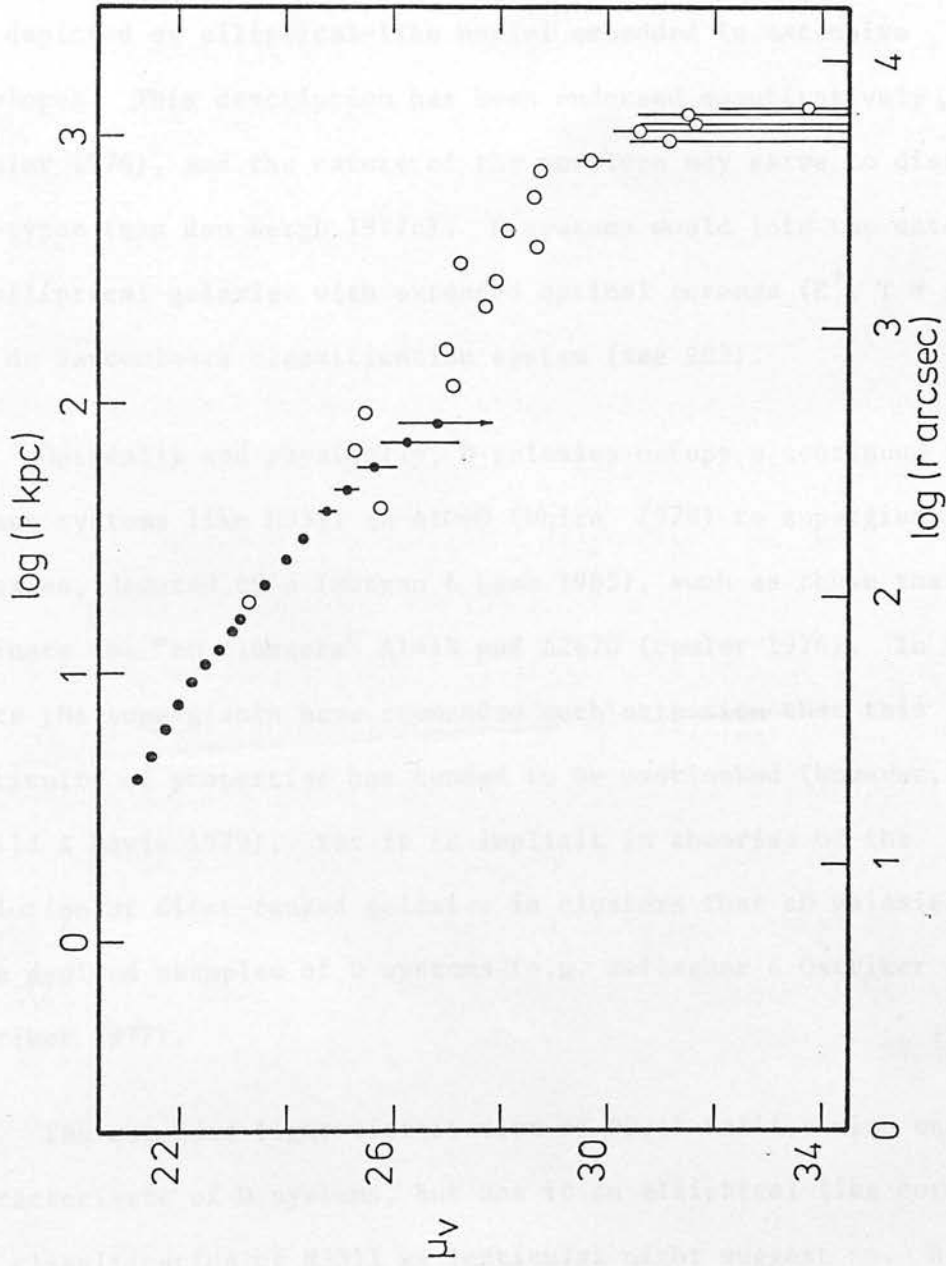


Figure 6.7

The projected radial profile of N3311 (filled circles) compared with the profile of the cluster (open circles).

for the integrated light of the proportion of non-members at each magnitude predicted by the field correction of Section 2.5.

#### 6.4 The Nature and Evolution of N3311

White (1978) has assigned N3311 to the class of objects known as D galaxies (Morgan 1958; Matthews, Morgan & Schmidt 1964). These are depicted as elliptical-like nuclei embedded in extensive envelopes. This description has been endorsed quantitatively (Oemler 1976), and the nature of the envelope may serve to distinguish sub-types (van den Bergh 1977c). D systems would join the category of elliptical galaxies with extended optical coronae ( $E^+$ ,  $T = -4$ ) in the de Vaucouleurs classification system (see RC2).

Optically and physically, D galaxies occupy a continuum from modest systems like N3311 in A1060 (White 1978) to supergiant D galaxies, denoted cD's (Morgan & Lesh 1965), such as those that dominate the "cD clusters" A1413 and A2670 (Oemler 1974). In recent years the supergiants have commanded such attention that this continuity of properties has tended to be overlooked (however, see Schild & Davis 1979). Yet it is implicit in theories of the evolution of first-ranked galaxies in clusters that cD galaxies are more evolved examples of D systems (e.g. Gallagher & Ostriker 1972; Ostriker 1977).

The extended light distribution of N3311 tallies with one characteristic of D systems, but has it an elliptical-like core? The classification of N3311 as lenticular might suggest so. However, its central surface brightness is singularly low among the giant

early-type members of this cluster, which contradicts the proposal (Godwin & Peach 1979) that brighter galaxies in rich clusters exhibit higher central surface brightnesses than fainter systems. Oemler (1976) states that cD galaxies possess exceptionally bright cores. Presumably this refers to the integrated light contained within the core radius, which - although it cannot directly be measured in N3311 due to saturation effects - increases with the scale length of the luminosity distribution, and the latter is large for N3311. The core luminosity of N3311 should be comparable to that of N3309, and neither seems abnormally luminous.

Comparison with other nearby clusters provides a better perspective. Godwin & Peach (1979) have listed the isophotal magnitudes  $M_{V_{25}}$  of the three brightest members of each of twelve rich clusters of galaxies, obtained by a programme of photographic surface photometry (Bucknell, Godwin & Peach 1979). Their data are not corrected for cosmological affects (see, for example, Green 1978), so only the two nearest clusters, Virgo I and A426, may straightforwardly be compared with A1060. The isophotal radii,  $r_{25}^*$  of the brightest members of the respective clusters are 43 kpc (N4486 = M87) and 41 kpc (N1275); adopting the distance scale used by Godwin & Peach ( $H_0 = 50 \text{ km s}^{-1} \text{ Mpc}^{-1}$ ), the equivalent dimensions are 78 kpc for N3311 and 39 kpc for N3309. On the extinction law of Sandage (1973a), the  $M_{V_{25}}$  magnitudes of N3311 and N3309 are  $-23.^m6$  and  $-22.^m8$  respectively, both intermediate relative to those objects of the same rankings in the other two clusters.

These results suggest that the dimensions of N3311 are anomalously large for its luminosity, a fact reflected in its low surface brightness. There is ample additional evidence that supports the picture of N3311 as a relatively extended system. Photoelectric aperture photometry matched to RC2 magnitude growth curves (cf. Chapter 4) indicates that half of the B-band emission from this galaxy comes from a region 131 arcsec in diameter, as compared to only 46 arcsec in the case of N3309. The measurements of di Tullio (1979) imply that at a surface brightness level corresponding to the faintest detectable isophote on ESO (B) films, N3311 is at least twice the size of its bright companion. The photoelectric data of Sandage (1973b) would indicate that the absolute metric V magnitude of N3311 interior to a standard diameter of 86 kpc is not significantly lower than the average for a large sample of brightest galaxies in clusters, contrary to expectation if its mean surface brightness is in fact unusually low. This inference, however, is mistaken, because the metric magnitude of N3311 was derived by Sandage on the erroneous assumption of conformity to universal linear dimensions shared by all first-ranked galaxies in clusters (Sandage 1972a).

It is obvious from the previous section that N3311 is a highly disturbed stellar system. Its bloated appearance and asymmetric light distribution strongly suggest that during the recent (the previous  $\sim 10^9$  year) evolution of the galaxy, strong dynamical interactions occurred between N3311 and other mass concentrations in the cluster. In this possible interpretation,

some of the kinetic energy of the intruding masses has been absorbed in dynamical heating of the stellar velocity distribution in N3311, reducing the negative binding energy of the galaxy and inducing it to expand. This hypothesis would be favoured by the discovery of concentrations of mass in the vicinity of N3311 which show properties that might be associated with gravitational perturbations.

Plausible candidates for such systems are compact galaxies. Fairall (1978b) has suggested that these are formed when the outer regions of normal galaxies are stripped away by tidal forces; and in particular, he speculates (Fairall 1979) that the extremely compact elliptical object number 61 in the Hydra cluster, located near N3311 and to its north (Figure 6.5), is the nuclear remnant of a giant elliptical galaxy that was largely denuded by N3311 during a previous encounter. However, the observed colours of galaxy 61 do not support this hypothesis. On the evidence of colour-aperture gradients in normal ellipticals, an intact and isolated elliptical core should be measurably more red than an entire galaxy of the same total luminosity; there is no indication that this is true of galaxy number 61 (see Figure 4.3).

A more likely explanation for the distension of N3311 is suggested by the presence of several fainter early-type galaxies which appear to be companions of the giant. As indicated in Figure 6.5, five of these small systems are located to the north-east of N3311, superposed on its halo and within 26 kpc of the nucleus as seen in projection. There is no similar clustering of faint galaxies in the halo of N3309. The strongly localised distribution of the



small systems near N3311 must affect its gravitational potential, perturbing stellar orbits and contributing to the global asymmetry. If these companions are actually immersed in the envelope of the giant, dynamical friction will operate, causing the captives to spiral towards the nucleus as they lose energy. The optical properties of the N3311 system are consistent with the theory of cluster evolution by galactic accretion and cannibalism (Ostriker & Tremaine 1975; Ostriker & Hausman 1977; Ostriker 1977; Hausman & Ostriker 1978).

This course of evolution presumes the initial existence of a "seed" galaxy within what was to become the Hydra I cluster, which began to interact with and to pillage its fellow cluster members. Such a system might have been a bright galaxy that initially dominated a subcluster of galaxies of low internal velocity dispersion. All known examples of D and cD galaxies occur within local density enhancements inside rich or poor clusters (White 1978), and within such associations gravitational encounters that might promote the formation of a D galaxy are likely to be relatively frequent. On an alternative scheme (Gallagher & Ostriker 1972), the observed D system was formed as tidal debris, liberated during galaxy collisions and interactions, gradually accumulated in the potential well at the centre of the cluster: no progenitor or "seed" was involved. At least four readily testable predictions follow from the debris hypothesis (e.g. Richstone 1976): the starcloud would have the same relatively blue colour as the outer regions of normal ellipticals; it would have no colour gradients from centre to rim; it should be



centrally located in the cluster, and its radial surface density distribution would be similar to, or somewhat steeper than, the radial luminosity profile of the cluster itself.

Photographic photometry confirms that N3311 satisfies the latter condition (Figure 6.7), and it has already been remarked in Chapter 2 that this galaxy is close to the centre of maximum number density in the cluster. The remaining requirements of the debris hypothesis, however, are clearly contradicted by photoelectric aperture photometry. The face-on intrinsic U-B colour of N3311 extrapolated to infinite aperture  $[(U-B)_T^0]$  is not appreciably bluer than those of the normal elliptical galaxies N3305 and N3309; the figures are  $0.^m48$ ,  $0.^m49$  and  $0.^m50$  for the three systems respectively. Moreover both the B-V and U-B colour indices of N3311 become progressively bluer through increasing apertures, in the same sense and perhaps more markedly than in the case of its companion, N3309. Since N3311 approximates so closely at the present epoch to the stellar composition and population gradients of a normal early-type giant, its formation entirely from tidal debris can be ruled out.

In the cannibalism hypothesis the dominant galaxy evolves towards the properties of a starpile, as distinct from always having existed as one (cf. Gallagher & Ostriker 1972). As captive galaxies disintegrate and lose their identity in the giant system, the integrated colour of the latter should gradually tend towards the mean colour of the galaxies consumed. Numerical simulations (Hausman & Ostriker 1978) indicate that much of the accreted material would have existed as galaxies with magnitudes close to  $M^*$ , the

characteristic magnitude or turnover point in the logarithmic integrated luminosity function (cf. Figure 6.4). These galaxies are less bright than the giant systems and hence should be bluer according to the colour-magnitude relation (Figure 4.3). The fact that N3311 itself is not relatively blue in comparison with N3309 and N3305 implies that N3311 cannot have accreted very many fainter galaxies during its past evolution.

However if it is accepted that N3311 is the only member of A1060 to be growing at the expense of lesser galaxies, a difficulty of interpretation does arise regarding the trends visible in the colour-magnitude diagram, Figure 4.3. Assuming that a linear colour-magnitude relation as observed by Sandage (1972b) applies initially to all unevolved early-type galaxies in clusters, Hausman & Ostriker (1978) predict a gradual flattening of this relation among bright galaxies as they proceed to digest faint companions. Near-equality of the colour of bright galaxies as seen in Figure 4.3 could have resulted, on this hypothesis, only if all of the galaxies on this plateau behave in the manner suggested for N3311.

Recent work by Michard (1979a,b) indicates, nonetheless, that too much should not be made of this apparent failure of the cannibalism hypothesis to account for the presence of a marked horizontal branch in the colour-magnitude diagram for A1060 in the absence of evidence that cannibalism is rife in the cluster. Michard has reported that early-type galaxies in each of many nearby groups and clusters exhibit different intrinsic U-B colours from those found in other groups. Within each cluster, galaxies separate into distinct

sequences and two separate colour-magnitude correlations are observed: bright galaxies populate the nearly horizontal branch; faint galaxies obey a steeper relation. These sequences also segregate in other planes, for example in the total magnitude-surface brightness plane, although insufficient photoelectric observations exist to confirm this in the case of A1060. The subject of the evolution of the colour indices of galaxies is clearly more complex than has hitherto been supposed, and further studies of colour-magnitude correlations are required.

It is tempting to speculate that N3311, uniquely among galaxies in the Hydra I cluster, is evolving towards the characteristics of a supergiant cD system. At present it most likely represents a transitional stage between elliptical galaxies with some indication of tidally distended haloes (Kormendy 1977b), and classic D galaxies such as N4874 in the Coma cluster.

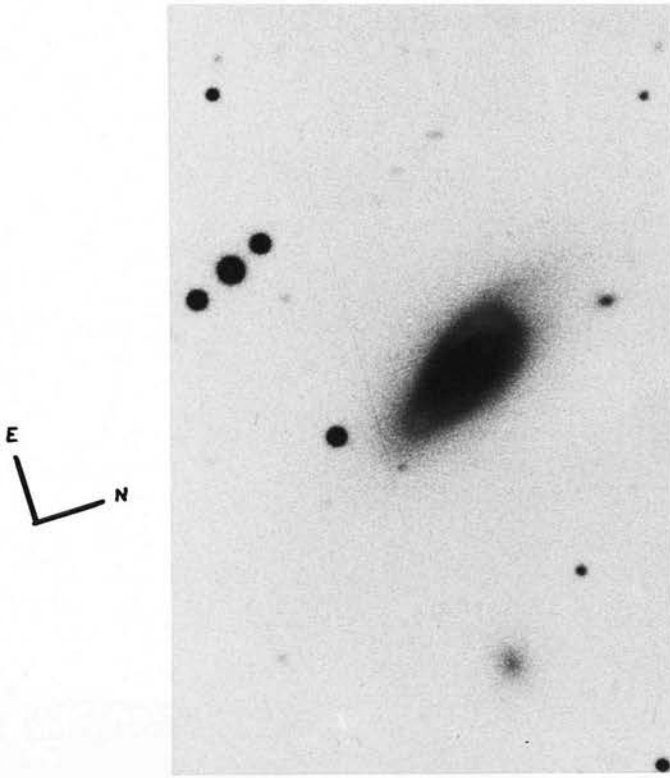
### 6.5 Peculiar Galaxies

NGC 3312 in the Hydra I cluster is a large early-type spiral which appears to be in the process of losing its interstellar dust and gas. It is highly distorted, laced with sharp dust lanes, and there are condensations or filaments to the south-west of the galaxy which may have been stripped from the disk (Gallagher 1978). The asymmetry of the disturbance discourages the idea that internal activity in N3312 is responsible, and the apparently greater stability of the stellar component as compared to the gaseous is likewise difficult to account for on a tidal interaction picture. The most natural explanation for

the appearance of N3312 is that it is being scoured of its interstellar material by the ram-pressure effect of its motion through an intracluster gas (cf. Gunn & Gott 1972).

Gallagher (1978) has remarked that there exist other "anaemic" or gas-poor spirals (van den Bergh 1976) besides N3312 in the core of Hydra I. Inspection of the prime focus plates taken with the CTIO 4m telescope reveals two particularly interesting candidates, both of which show evidence of structural perturbations. The first is the peculiar spiral galaxy number 60 = N3307, located 5 arcminutes west of N3309/N3311. N3307 appears to have narrow, bright and well-defined arms which are strongly warped relative to a plane. On the available 4m plates no Population I tracers (O, B-star complexes, HII regions) are visible in the arms, which suggests that no recent star formation has occurred; this absence is one symptom of gas loss. The optical characteristics of N3307 are those of the smooth-arm spiral galaxies (Strom, Jensen & Strom 1976) whose existence in several other X-ray emitting clusters has been reported by Wilkerson, Strom & Strom (1977).

The second unusual system and possibly a better example is the galaxy number 43, which is also classified as a peculiar spiral. It has a diffuse, low-surface-brightness arm which encircles the nucleus and may formerly have closed in on itself; the presence of a second main arm is uncertain. Apart from knotty condensations on the inside of the first arm, most internal structure in the arm has been obliterated. The inner boundary of this arm is sufficiently sharp, on the south side of the nucleus, to reveal a possible narrow inner (detached?)

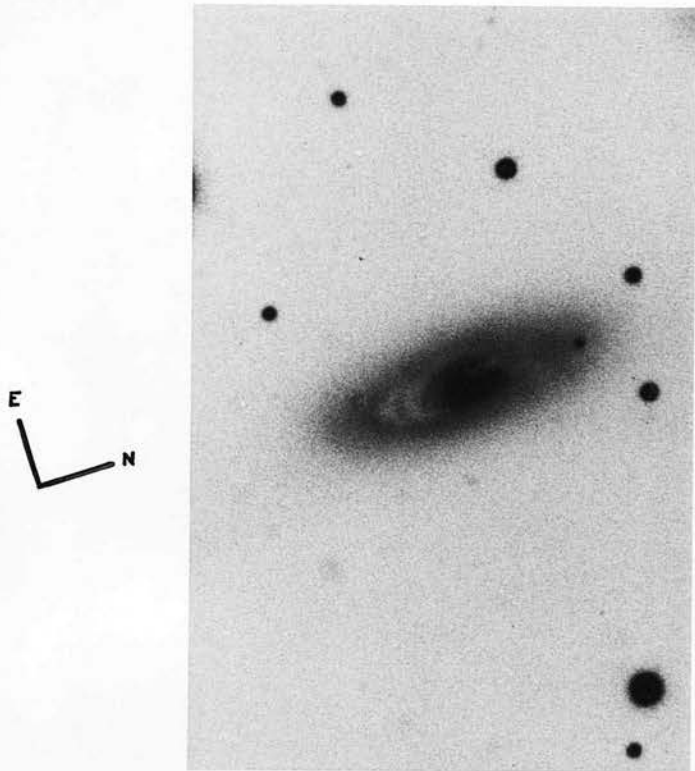


Left: the smooth-arm spiral galaxy no. 60 = N3307. Enlargement of plate no. 641 taken by M.G. Smith at the prime focus of the 4m reflector at CTIO (see page 169 for further details).

Scale  $1.26 \text{ arcsec mm}^{-1}$ .

Seeing  $< 1 \text{ arcsec}$ .

Right: the peculiar system no. 43, as recorded on plate no. 641.





arm of the galaxy; on the north side the arm dissolves chaotically. In addition a prominent absorption lane trails across the northern half of the nuclear region, nearly parallel to the long axis.

If the interpretation of these systems as disturbed spirals is correct, and star-forming events in their disks have ceased, their integrated colours should be anomalously red relative to normal spirals. A comparison of the photoelectric colour indices of N3307 and the normal barred spiral number 86, which is of very similar stage in the revised Hubble classification system, shows that N3307 is indeed redder (Table 4.3), and at the same time less red than a SO. This single observation is, therefore, consistent with the belief that there exist disk galaxies in A1060 which have lost, or are currently losing, a substantial fraction of their gas. Since ram-pressure stripping would be most efficient in the core of a cluster where the mean density of the pervading medium is highest, it is worth noting that the three galaxies mentioned are all within a projected distance of 0.2 Mpc of the centre of A1060.

Only one possible and very dubious instance has been discovered of a spiral system in this cluster in which gas loss has already proceeded to extremes. This is galaxy number 112, a fainter "lenticular" with a small bulge-to-disk ratio and a very smooth envelope. Inspection of CTIO plate 995 leaves a persistent impression of two small, weak arms in the immediate region of the nucleus, in which it bears some resemblance to IC 2951 in A1367 (Strom & Strom 1979b), a smooth-arm spiral in which the arms have almost disappeared and which is otherwise visually indistinguishable from a SO. However,



the arms are not revealed on the shorter-exposure plate number 641; UKST photographs have insufficient resolution to reveal the structure.

Since at least two of the three disturbed spirals - N3312 and number 43 - appear to be currently losing their gas, the stripping of these systems may have begun only recently (a few  $10^9$  years ago). Gas shed by evolving stars in the galaxies would tend to replenish each interstellar medium (e.g. Faber & Gallagher 1976), probably to the extent of rendering such systems invulnerable to stripping over most of their lives (Gisler 1979). The era of gas depletion may have begun when gas injection rates dropped sharply. As the galaxies leave the stripping regime due to their peculiar motions through the cluster, gas-poor character could be maintained through the internal sweeping action of galactic winds driven by supernovae (Matthews & Baker 1971; Johnson & Axford 1971; Bregman 1978).

Whether such systems will eventually evolve into bona fide lenticulars is a question open to dispute. Such transmutations could in principle explain why so many S0s and so few spirals are observed in the central regions of Hydra I; it would also account for the lower proportion of spiral galaxies generally in higher-luminosity X-ray clusters (Bahcall 1977b; Tytler & Vidal 1978) in which the intracluster medium is most likely denser and gas stripping consequently more efficient. However, in the apparent absence (in A1060) of ancient spiral relics which more nearly resemble S0s, it is doubtful whether sufficient time has elapsed for many conversions to have run to completion in A1060. Furthermore, many S0s are situated too far from the centre of any rich cluster for external ram

pressure to have been effective. They are likewise too far from the dense cores of clusters for collisions between the spiral and other galaxies (Spitzer & Baade 1951) to be a viable alternative explanation for the formation of the SOs (see Sandage & Visvanathan 1978b for a review).

To summarise, the distinctive features of the three spiral systems (low rates of star formation in the spiral arms, structural distortion, generally "anaemic" appearance and in the one case measured, reddening of the system) suggests strongly that in the core of A1060, an era of gas stripping has recently begun. The present nature of these candidate "lenticulars" in A1060 could be clarified by HI-line observations of gas abundance and further photo-electric photometry.

## CHAPTER 7

### DWARF GALAXIES IN A1060

#### 7.1 Definition and Review

Dwarf galaxies - a term extended, in the present context, to the most extreme examples of miniature galaxies, including the "pygmy" and "gnome" categories of Zwicky (1971) - represent the faintest and least well delineated region of the spectrum of galaxian luminosities. Like the transition between giant systems and normally bright galaxies, the separation of normal galaxies from dwarfs is to some extent arbitrary; de Vaucouleurs (1975) suggests the definition  $M_{B_T}(\text{dwarf}) \geq -16.0$ . Systems fainter than absolute magnitude  $M_{B_T} = -15$  probably do not contribute significantly to the total mass density of a cluster (Reaves 1967).

The absolute number of these low-luminosity, and presumably low-mass, galaxies is uncertain, even within the nearest  $3 \times 10^3 \text{ Mpc}^3$  (Tammann & Kraan 1978). There exist strong logical arguments (Arp 1965) and observational evidence (e.g. Sandage, Tammann & Yahil 1979) which suggest that at all distances there remain low surface brightness galaxies that are not catalogued. Arp (1965) pointed out that the  $\log L : \log R$  plane, where  $L$  and  $R$  represent the total luminosities and radii of different classes of object, is everywhere populated where it is detectable and that there is no justification for supposing that systems of lower surface brightness do not exist. The omission, as distance increases, of a growing proportion of intrinsically faint galaxies in a sample selected to a limiting

apparent magnitude shifts the observed distribution of absolute magnitudes in the sample towards brighter values (Sandage, Tammann & Yahil 1979).

An important corollary of dynamical relaxation of clusters over long timescales is the radial segregation of galaxies according to mass; for as equipartition of energy is achieved, the less massive galaxies assume a higher velocity dispersion and a relatively broader quasi-equilibrium distribution. Many studies have been pursued which test directly or indirectly for the presence of observable segregation in clusters (Zwicky 1942, 1956; Rudnicki 1963, 1967; Kwast 1966; Rudnicki & Baranowski 1966a,b; Rood 1969; Bahcall 1972, 1973c; Noonan 1971, 1974; MacGillivray et al. 1976b, and others), but most of these have either lacked photometric calibration or involved clusters so distant that genuine dwarfs were unobservable. Moreover the difficulty of distinguishing faint galaxies from intrinsically brighter systems at greater distances has sometimes led to equivocal or even contradictory conclusions (cf. Zwicky 1956, Noonan 1974). It does appear that in rich regular clusters segregation in the expected sense is observed (see Bahcall 1977a), but competing evolutionary effects, which are discussed later, complicate the picture.

The number of moderately rich groups of galaxies that are near enough for their dwarf members to be observationally accessible is limited. Jones & Jones (1978) list sixteen candidates, populous and otherwise. Reaves (1967) has examined the distribution of dwarf galaxies associated with the most amenable systems, namely the Virgo

(Reaves 1956, 1977a) and Fornax (Reaves 1964) clusters; Hodge and his collaborators (Hodge 1959, 1960; Hodge, Pyper & Webb 1965) have also investigated faint objects in the Fornax cluster. The next nearest dense groupings that are comparably rich are Pegasus, Centaurus I and Hydra I.

The structural properties of dwarf galaxies are diverse (e.g. Reaves 1977a,b; Webster et al. 1979), but a dwarf may be one of three general types: dwarf elliptical; dwarf spiral or irregular; or of a species exemplified by IC3475 in the Virgo cluster (Reaves 1977b). Hodge (1971), however, does not recognise the latter as a distinct class, whereas van den Bergh (1959) distinguishes dwarf spirals from Magellanic irregulars. Low surface brightness is the most common visual characteristic of dwarfs, collectively. For example, the Sculptor and Fornax dwarf elliptical members of the Local Group (Shapley 1939) have peak surface brightnesses of only  $24.0 \pm .5$  B mag arcsec<sup>-2</sup>; this figure would be 17 or less for giant galaxies. The surface brightnesses and optical appearances of galaxies vary with distance (van den Bergh notes that IC3475 types are more easily identified when relatively distant than when nearby and completely resolved into stars).

Besides Sculptor and Fornax, other local dwarf ellipticals - N147, N185, N205 and N221=M32 - and the representative dwarf irregulars N6822 and IC1613 (see also Fisher & Tully 1979) are readily visible on Palomar Observatory Sky Survey (POSS) and Southern Sky Survey plates. The detection on POSS plates of the dwarf systems identified in the Coma cluster by Reaves (1966) and



Rood & Baum (1967), each of whom had access to prime focus plates from the 5.1m Hale telescope, poses a greater challenge: as an illustration, M32 at the distance of the Coma cluster ( $\mu \approx 34.1^m$ ) would have an apparent magnitude  $B_T = 19.2^m$  and a standard angular diameter  $D(0)$  of 3 arcseconds. Most of the dwarfs are indistinguishable from stellar images and emulsion defects.

On the new deep IIIa-J surveys, however, dwarf systems in the Virgo I ( $\Delta = 12.5$  Mpc) and Fornax I ( $\Delta = 17$  Mpc, de Vaucouleurs 1977b) clusters are more conspicuous. Thirty-nine of the fifty objects discovered by Hodge et al. (1965) in the latter cluster lie within ESO/SRC Field No 358 on the basis of their epoch 1950 positions, and all but six of those found by Reaves (1964). Inspection of high quality survey film revealed that seven of those objects described by Hodge et al. as similar in appearance to dwarf elliptical galaxies (their designations 6, 9, 11, 22, 37, 38 & 41) are in fact spiral or irregular galaxies; two listed as irregulars (numbers 5 & 39) are probably misclassified, and a further two (29 & 35) are not visible on the film copy. Additional dwarf-like objects exist that are not tabulated by Hodge et al. The survey of Reaves exhibits comparatively greater homogeneity, but includes some late-type galaxies which rather than being members of Fornax I, are probably intrinsically brighter systems in the background. In less than twenty cases do the independent surveys agree on the selection of a given galaxy. The contention of Reaves that the discovery of dwarf galaxies under favourable circumstances is a matter of recognition rather than detection, therefore, receives support; a more striking illustration



is the disclosure (Fisher & Tully 1975) that only 15% of the DDO "dwarf" systems catalogued by van den Bergh (1959, 1966) are true dwarfs on the de Vaucouleurs definition. Identification problems will be greatly alleviated when redshifts surveys of the nearer clusters become available (see Jones & Jones 1978). In the meantime, purely optical criteria must be used.

## 7.2 Selection Procedure

For the purpose of locating dwarf members of the Hydra I cluster, deep IIIa-J reject survey plates were examined for objects which were visually comparable to the Virgo and Fornax dwarfs discovered by Reaves and others. This approach is important, because unlike Reaves (e.g. 1977a), no use is made of the apparent surface distribution of like systems, relative to that of normal galaxies, as an indicator of association of those systems with the cluster. Hence, conclusions regarding the true surface distribution of dwarfs remain valid.

The COSMOS machine was not engaged for the detection and measurement of dwarfs because the relevant regions of the photographic plates were appreciably vignetted; the eye is a more reliable detector within confused regions of multiple luminosity gradients, and moreover the high density of star images ( $\sim 500 \text{ cm}^{-2}$  with  $J \leq 22^m$ ) would have taxed any automatic star sorting routine. The procedure was to insert a metal-on-glass graticule in the design of a labelled 10 x 10 grid behind one eyepiece of a binocular microscope through which each plate was viewed. The position of every recognisable

galaxy was marked on a reproduction of the grid scaled to give a convenient twenty times magnification of the plate ( $3.4 \text{ arcsec mm}^{-1}$ ). Stepping down successive columns while recording galaxies in this manner, a composite chart of a region  $1.8 \text{ degree}^2$  in area and centred on the Hydra cluster was generated. Registration was later restored by matching the individual grids to the coordinate system defined by COSMOS coarse-mode measurements. Candidate dwarfs were emphasized on the charts.

Opportunity to verify these identifications was subsequently provided by two prime focus IIIa-J plates obtained on the 3.8m reflector at CTIO. Their particulars are summarised in Table 7.1. On plate 995, which displays the central  $0.7 \text{ degrees}^2$  of the cluster at superior resolution and to a slightly brighter limiting magnitude than do the sky-limited Schmidt plates, all of the more plausible dwarf galaxy candidates were recovered and five additional faint examples were also found. Such disagreement as there was essentially concerned image structure, as emulsion granularity is relatively more serious at the scale of the Schmidt plates.

The probable dwarf systems discovered in projection within the central region of the Hydra I cluster are listed in Table 7.2. Those galaxies within the central field of plate 995 (group A) are set apart from the systems which are outside this field and visible only on the UKST plate (group B). The least persuasive identifications, i.e., the smallest and faintest images noted, are not included. The systems exhibit much variety of form, and the notes in Table 7.2 supply an abbreviated description of their more obvious optical

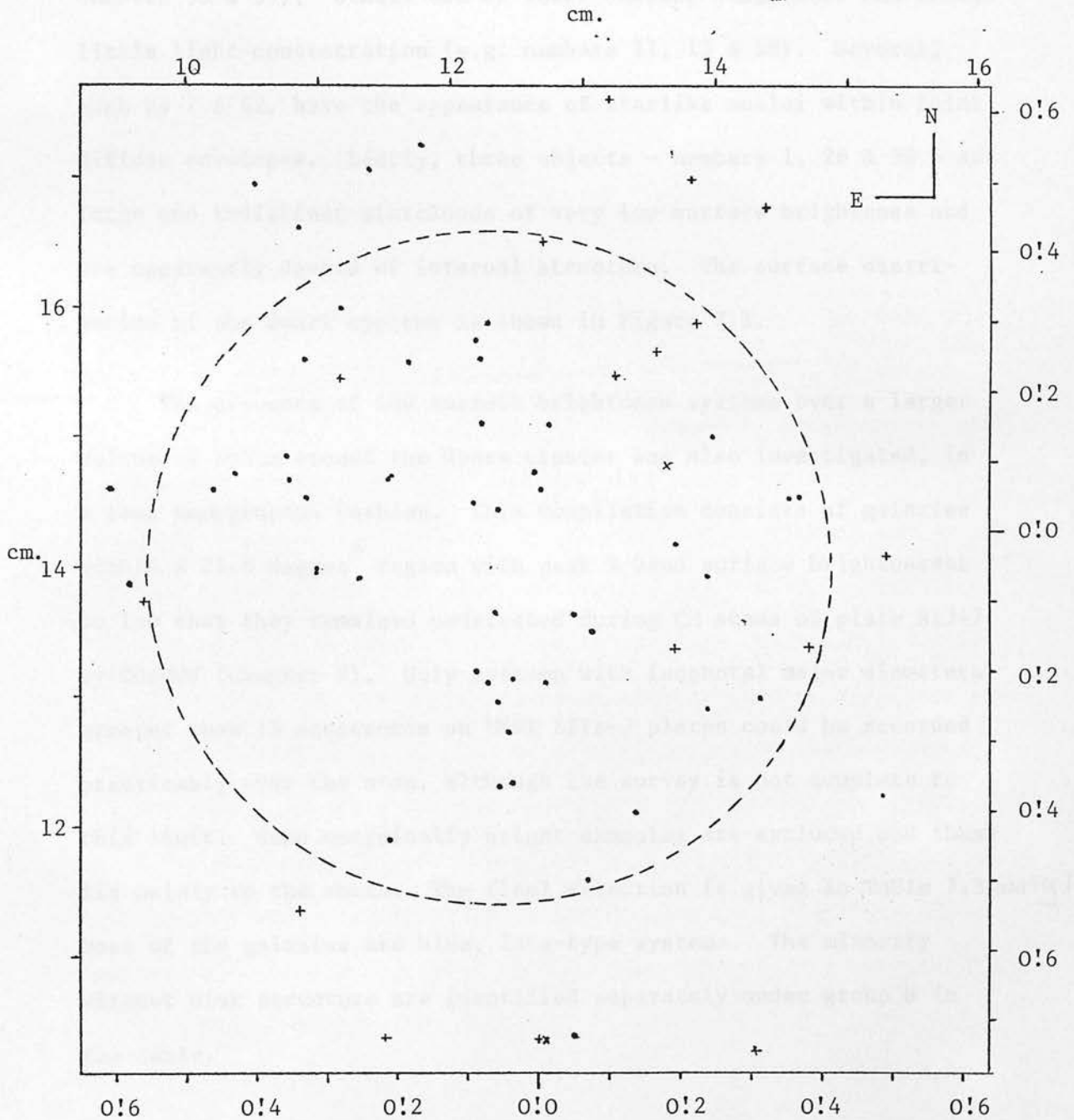


Figure 7.1

The surface distribution of dwarf galaxies (dots: Table 7.2) and low surface brightness systems ('plus' symbols: Table 7.3) near the centre of A1060. Crosses distinguish the two galaxies listed in Table 7.2 which may have spiral structure. The approximate area covered by plate 995 ("central region") is shown by the dotted circle.

characteristics. Some dwarfs are fairly compact, relatively bright and have significant central concentration of luminosity (e.g. numbers 33 & 39); others are of lower surface brightness and reveal little light concentration (e.g. numbers 11, 13 & 38). Several, such as 7 & 42, have the appearance of starlike nuclei within faint diffuse envelopes. Lastly, three objects - numbers 1, 28 & 36 - are large and indistinct starclouds of very low surface brightness and are apparently devoid of internal structure. The surface distribution of the dwarf systems is shown in Figure 7.1.

The presence of low surface brightness systems over a larger volume of space around the Hydra cluster was also investigated, in a less homogeneous fashion. This compilation consists of galaxies within a  $21.8 \text{ degree}^2$  region with peak R-band surface brightnesses so low that they remained undetected during CM scans of plate R1347 by COSMOS (Chapter 2). Only systems with isophotal major diameters greater than 15 arcseconds on UKST IIIa-J plates could be recorded practicably over the area, although the survey is not complete to this limit. Some untypically bright examples are excluded and these lie mainly to the south. The final selection is given in Table 7.3 and Fig. 7.2. Most of the galaxies are blue, late-type systems. The minority without disk structure are identified separately under group B in the table.

### 7.3 Magnitude Estimates

To estimate the apparent magnitudes of the dwarf systems listed in Table 7.2 it is necessary to advance indirect arguments. Firstly,

it can be appreciated immediately that the observation that the brightest members of the system of globular clusters associated with NGC 3311 (cf. Figure 2, Smith & Weedman 1976) are not visible on the Schmidt plates constrains the limiting magnitude to  $J < 23^m.5$ . More usefully, the number counts of galaxies of all types within the  $1.8 \text{ degree}^2$  may be compared with empirical number versus apparent limiting magnitude relations derived from other photometrically calibrated number counts (e.g. Shane 1975, Karachentsev & Kopylov 1977, Tyson & Jarvis 1979, MacGillivray & Dodd 1980). The 2866 galaxies recorded represent a density of  $1570 \text{ degree}^{-2}$ , of which about  $1300 \text{ degree}^{-2}$  are not Hydra cluster members; according to the compilation of Karachentsev & Kopylov, this surface density is typically reached at approximately  $J = 21^m.6$ , equivalent to an apparent magnitude of  $22^m.0$  after extinction ( $A_B = 0^m.38$ , RC2). The relation of Tyson & Jarvis (1979) would indicate the limit  $J \approx 21^m.3$  or  $21^m.8$  on the basis of their empirical extinction law, and Shane's  $J > 22^m$  with extrapolation. Three qualifications of these figures are necessary: density fluctuations (greatly exceeding the statistical) due to largescale clustering of galaxies are liable to vitiate such comparisons; patchy galactic obscuration (e.g. Holmberg 1974, Sandage 1976) is unpredictable to a certain extent; and inconsistencies remain regarding the mean rate of increase of the number of galaxies with magnitude.\* It is also uncertain whether the effective limiting magnitude for detection is the same for diffuse systems as for compact images. For instance, a disk of uniform apparent

\* Shane (1975) and Karachentsev & Kopylov (1977) find  $\log N(<m_J)$  to increase progressively more slowly at faint  $m$ ; Tyson & Jarvis (1979) report a constant  $d \log N(<m_J)/dm_J = 0.41$  in  $17^m < J < 24^m$ , but Kron (1977) quotes a constant Euclidean slope of 0.6 in the blue to  $B = 24^m$ .



surface brightness  $\mu_J = 25 \text{ mag arcsec}^{-2}$  would have an integrated magnitude  $J = 20.^m3$  interior to 10 arcseconds diameter; such an object would probably have been detected. Having regard of all clues, a limiting magnitude of  $21.^m0$  is adopted. While obviously uncertain, this figure is insensitive to the allowance made for the local excess of galaxies represented by A1060 itself, because of the steep increase of the number of field galaxies per magnitude interval at this limit.

As for the galaxies of Table 7.3, their effective radii are about 15 arcseconds in the mean at the visual limiting isophote on survey plates. By reference to Figures 5.14(b) and 5.12, a normally luminous system of this size would have a blue magnitude  $B_{25} = 16.4$ , corresponding to  $V_{25} = 15.6$  for late-type galaxies of group A. Since by definition the latter were missed during a survey complete to  $V_{25} \approx 17.^m0$ , these galaxies are at least one magnitude fainter than normal systems of comparable size.

#### 7.4 Results

It is evident from Tables 7.2 and 7.3 that there are relatively few low-luminosity spiral and irregular galaxies in the cluster core and a preponderance in the outer regions. Figure 7.2 presents the surface distribution of the late-type systems over the  $21.8 \text{ degree}^2$  studied. The weak clustering in the general vicinity of the cluster suggests association, but few are found in the core itself. Because of their comparatively large angular dimensions, it seems unlikely that most of the systems can be distant; they probably form part of a cloud of late-type galaxies surrounding the core of earlier types



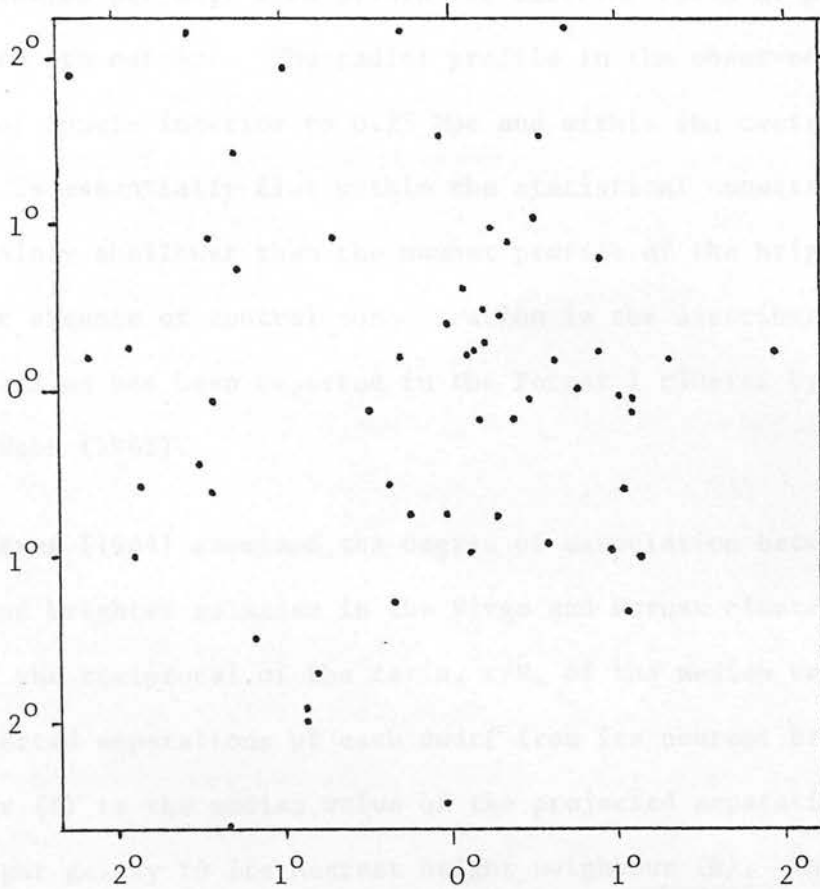


Figure 7.2

The surface distribution of the low surface brightness galaxies listed in Table 7.3. There is some suggestion that these systems are associated with A1060. Compare with Figures 2.7 - 2.9.

(compare with Figures 2.7 - 2.9).

In the inner regions of Hydra I there are appreciably more dwarf galaxies per unit area within the circular field of plate 995 than there are outside. The radial profile in the observed number density of dwarfs interior to 0.25 Mpc and within the central circle, however, is essentially flat within the statistical uncertainties, and certainly shallower than the number profile of the bright members. A similar absence of central concentration in the distribution of dwarf galaxies has been reported in the Fornax I cluster by Hodge, Pyper & Webb (1965).

Reaves (1964) examined the degree of association between dwarfs and brighter galaxies in the Virgo and Fornax clusters by means of the reciprocal of the ratio,  $r/R$ , of the median value of the projected separations of each dwarf from its nearest bright neighbour ( $r$ ) to the median value of the projected separations of each bright galaxy to its nearest bright neighbour ( $R$ ). The dimensionless quantity  $R/r$  is independent of the density (and the distance) of the cluster; if the dwarfs are located preferentially in the vicinity of giant galaxies,  $R/r$  would exceed unity. Adopting the criterion  $M_{V_{25}} < -18.0$  as an arbitrary definition of a bright galaxy, the data here indicate that  $R/r(\text{A1060}) = (43 \text{ kpc})/(45 \text{ kpc}) = 1.0$ . For comparison, this ratio is 1.0 for Virgo I and 0.6 for the Fornax cluster (Reaves 1964). Combined with the result  $R/r(\text{A1656}) = 0.8$  implied by figures given by Rood & Turnrose (1968) pertaining to Reaves' (1966) Coma cluster dwarfs, this is convincing evidence that in general, dwarf systems in the central regions of rich clusters are not satellites of giant galaxies.

The morphology of the faint dwarf systems near the centre of A1060 is noteworthy. The scarcity or absence of dwarf spirals and irregulars parallels the situation in Virgo I (Reaves 1977a,b). The diffuse systems within A1060 with sharp nuclei are similar to some of Reaves' Virgo dwarfs, which may be the same systems as the "nucleated dwarfs" described by Romanishin, Strom & Strom (1977). These types are also visible in the Fornax cluster, yet they have no counterparts in the Local Group. Strom & Strom (1978b) have also reported higher-luminosity versions of nucleated dwarfs in the Perseus cluster (A426), but they find none in A262, A1367 and A1656. A possible inference is that the composition of the local dwarf population varies from cluster to cluster.

The distributions of the apparent ellipticities,  $e = 10(1 - d/D)$ , of the "early-type" and "late-type" dwarfs listed in Tables 7.2 and 7.3(A) respectively are not significantly different at the 20% level in the Kolmogorov-Smirnov test. There are appreciable numbers of elongated ( $e > 3$ ) systems among early as well as late types (although individually, estimates of  $e$  are imprecise). Hodge (1971) notes that dwarf ellipticals in the Local Group have quite high ellipticities; a good example is the Carina dwarf galaxy recently discovered by Cannon et al. (1977). The term applied to such systems by some authors, "dwarf spheroidal", is unfortunate.

Altogether 51 faint galaxies have been identified in the central  $0.73 \text{ degree}^2$  of A1060 in the approximate magnitude range  $17.5 < V_{25} < 21.0 - \langle B_{25} - V_{25} \rangle - A_V$ , i.e.  $17.5 < V_{25} < 20.4$  (Figure 7.1). A linear extrapolation of the Abell-type luminosity function

defined by galaxies brighter than  $V_{25} = 17.0$  (Figure 6.4) would predict an average of 64 galaxies to each  $0.73 \text{ degree}^2$  interior to a radius of 1.5 Mpc. This shortfall of observed numbers relative to prediction is probably significant, for two reasons. Firstly, it would be an exceptional optical survey of intrinsically faint systems that was not to some degree contaminated by normally luminous galaxies; the precedents are not encouraging. Secondly, the actual surface density of galaxies towards the cluster centre ought certainly to be higher than the average predicted by a luminosity function for a much larger area. The force of this argument is only partially diminished by the lack of observable central concentration of the identified dwarfs.

It is possible, nevertheless, that the incorrect conclusion may be drawn, because of the strong bias that favours selection of low surface brightness members of Hydra I. Small compact ellipsoids undoubtedly exist, but are extremely difficult to identify at the distance of A1060 by appearance alone (though the compacts in the halo of N3311 are exceptions, the evidence is circumstantial). The importance of the omission of these types of galaxies cannot be assessed until integrated colours or redshifts are determined, a technically ambitious undertaking that demands the best equipment. The facts presented in this Chapter provide sufficient grounds for questioning whether the numbers and the distribution of poor stellar systems really represent a logical continuation of the properties of the larger systems. There is in fact evidence of a breakdown of a bilinear luminosity function at faint magnitudes in the Coma cluster

(Abell 1977), but the sense of the departure is entirely contrary to that implied here for A1060, and appears to be better determined. The issue is still open, and comparable studies of other nearby clusters would be worthwhile.

Table 7.1

Details of the prime focus plates from the 4m reflector at CTIO

Plate No	641	995
Date exposed	16-17 Feb. 1975	10-11 April 1975
R.A. (1975)	$10^{\text{h}}35^{\text{m}}.7$	$10^{\text{h}}35^{\text{m}}.7$
Dec. (1975)	$-27^{\circ}23'$	$-27^{\circ}23'$
Filter	GG385	GG385
Emulsion	IIIa-J	IIIa-J
Grain type	fine	fine
Exposure (minutes)	45	60
Equivalent photo- electric waveband*	"J"	"J"

\* "J" is a UKSTU designation for a IIIa-J/GG385 combination and is unrelated to the Johnson J photoelectric waveband at wavelength  $1.25\mu\text{m}$ .



Table 7.2

A list of probable dwarf galaxies in A1060. The "central region" is defined by the dotted circle in Figure 7.1.

No	X	Y	$\alpha$ (1950.0)			$\delta$			D	d	Notes
			h	m	s	0	'	"			
A. Galaxies within central region											
1	146400	145400	10	32	38.5	-27	10	12	21	21	F
2	145700	145300	10	32	42.0	-27	10	19	11	9	F,U
3	143500	130050	10	32	52.8	-27	27	24	11	6	U
4	139950	150170	10	33	11.0	-27	04	53	8	7	
5	139500	129200	10	33	13.0	-27	28	22	13	8	F
6	139400	139300	10	33	13.6	-27	17	04	7	6	B
7	137050	141800	10	33	25.5	-27	14	16	11	11	F,pn
8*	136400	147900	10	33	28.7	-27	07	23	19	9	B
9	134000	121150	10	33	40.7	-27	37	24	11	10	
10	130600	135200	10	33	57.9	-27	21	40	11	9	
11	130300	115750	10	33	59.3	-27	43	27	13	6	
12	129200	136400	10	34	05.0	-27	20	20			pn
13	127350	151000	10	34	14.4	-27	03	59	11	6	
14	126900	131200	10	34	16.6	-27	26	09	11	7	F
15	126700	146050	10	34	17.6	-27	09	31	11	9	
16	126300	147350	10	34	19.6	-27	08	04	11	6	
17	124640	123350	10	34	27.9	-27	34	53	30	12	
18	124500	131350	10	34	28.7	-27	25	59	15	9	F,pn
19	124100	127300	10	34	30.7	-27	30	31	9	7	B,pn
20	123520	123080	10	34	33.6	-27	35	11	34	13	
21	123432	129560	10	34	34.1	-27	27	59	18	15	B,pn
22	123350	144400	10	34	34.5	-27	11	22	9	7	B,pn
23	123250	136450	10	34	35.0	-27	20	16	9	6	F
24	122699	158856	10	34	37.8	-26	55	11	11	11	B,pn
25	122600	131100	10	34	38.3	-27	26	16	7	4	
26	122250	151100	10	34	40.0	-27	03	52	10	9	U
27	122150	156150	10	34	40.5	-26	58	13	10	7	
28	121900	157500	10	34	41.8	-26	56	42	26	24	F
29	121750	131900	10	34	42.6	-27	25	22	10	6	B
30	121600	144900	10	34	43.3	-27	10	49	15	6	F

contd....

Table 7.2 (contd.)

<u>No</u>	<u>X</u>	<u>Y</u>	<u><math>\alpha</math></u> (1950.0)			<u><math>\delta</math></u>			<u>D</u>	<u>d</u>	<u>Notes</u>
			<u>h</u>	<u>m</u>	<u>s</u>	<u>0</u>	<u>'</u>	<u>"</u>			
31	116800	155900	10	35	07.4	-26	58	30	9	9	F,pn
32	115440	146880	10	35	14.4	-27	08	32	8	7	
33	115362	119064	10	35	14.9	-27	39	44	10	10	B
34	115130	146640	10	35	15.9	-27	08	48	10	10	
35	112800	138950	10	35	27.6	-27	17	28	14	8	pn
36	111650	159950	10	35	33.2	-26	53	57	23	13	F
37	109700	139650	10	35	43.2	-27	16	40	7	5	
38	109700	139300	10	35	43.2	-27	17	04	13	7	
39	108900	145250	10	35	47.2	-27	10	24	10	8	B
40	108850	156000	10	35	47.5	-26	58	22	11	9	F,U
41	107700	146650	10	35	53.2	-27	08	50	11	6	F
42	107500	148400	10	35	54.2	-27	06	52	10	10	F,pn
43	103500	147000	10	36	14.3	-27	08	26	10	7	
44	102000	145750	10	36	21.9	-27	09	49	10	8	B

## B. Galaxies outside central region

45	152500	122600	10	32	06.9	-27	35	38	20	13	B
46	129200	103900	10	34	04.8	-27	56	43	11	9	F
47*	126600	103700	10	34	18.0	-27	56	56	13	8	
48	117690	172690	10	35	02.8	-26	39	40	20	17	
49	113900	170700	10	35	21.9	-26	41	53	13	11	F
50	108600	166300	10	35	48.6	-26	46	48	19	13	B,U
51	105200	173000	10	36	05.6	-26	39	17	21	10	B
52	95700	138900	10	37	01.7	-27	09	17	11	8	U

Column 1: running number (an asterisk indicates that weak spiral structure in the galaxy is visible); cols. 2 & 3: coordinates in the system of the Appendices, in microns; cols. 4-9: epoch 1950 coordinates; col. 10: micrometric major-axis diameter, arcseconds; col. 11: minor axis diameter; col. 12: optical characteristics (see key).

Key to symbols    B = relatively bright    F = very faint  
                          pn = point nucleus            U = uncertain morphological type

Table 7.3

A list of low surface brightness systems in a region of 21.8 degree<sup>2</sup> centred on A1060.

No	X	Y	$\alpha$ (1950.0)			$\delta$			D	d	Notes
			h	m	s	0	'	"			
A. Disk or irregular systems											
1	232390	155800	10	25	26.7	-26	57	38	26	23	
2	198185	153420	10	28	18.4	-27	00	49	15	13	
3	188720	89520	10	29	01.9	-28	12	20	30	27	
4	186085	140510	10	29	18.7	-27	15	24	27	20	p
5	185960	136810	10	29	19.2	-27	19	33	30	20	
6	184720	72170	10	29	21.4	-28	31	48	23	20	
7	183230	111900	10	29	30.9	-27	47	22	34	27	
8	182380	141880	10	29	37.4	-27	13	55	26	19	F
9	178240	92980	10	29	55.3	-28	08	35	37	23	
10	175860	185900	10	30	12.2	-26	24	39	17	12	
11	175400	156210	10	30	13.01	-26	57	56			
12	168360	143950	10	30	47.4	-27	11	38	27	20	
13	164460	260530	10	31	11.4	-25	01	11	30	17	
14	161140	153770	10	31	24.1	-27	00	42	28	22	
15	158879	94490	10	31	33.8	-28	07	03	37	17	
16	156900	225470	10	31	47.4	-25	40	30	20	13	
17	154510	199590	10	31	59.1	-26	09	28	34	27	
18	152950	141005	10	32	05.1	-27	15	03	27	20	
19	147180	133940	10	32	34.1	-27	22	59	23	23	
20	146215	191244	10	32	40.3	-26	18	51	15	14	
21	143790	167480	10	32	51.7	-26	45	28	27	15	B
22	143100	102600	10	32	54.2	-27	58	04	18	15	
23	141960	195800	10	33	01.3	-26	13	47	40	17	
24	138790	158810	10	33	16.7	-26	55	11	30	13	
25	138190	169860	10	33	19.9	-26	42	49	40	34	p
26	136960	133800	10	33	25.7	-27	23	10	40	34	
27	135680	156690	10	33	32.4	-26	57	33	32	23	F
28	133840	91380	10	33	41.1	-28	10	38	20	13	F
29	132640	154880	10	33	47.7	-26	59	35	20	15	
30	132400	176000	10	33	49.0	-26	35	57	34	27	

contd...

Table 7.3 (contd.)

<u>No</u>	<u>X</u>	<u>Y</u>	$\alpha$ (1950.0)			$\delta$			<u>D</u>	<u>d</u>	<u>Notes</u>
			<u>h</u>	<u>m</u>	<u>s</u>	<u>0</u>	<u>'</u>	<u>"</u>			
31	127140	165120	10	34	15.4	-26	48	08	34	27	
32	126600	103700	10	34	18.0	-27	56	52	15	8	
33	126020	10370	10	34	20.7	-29	41	18	26	10	B
34	124320	225500	10	34	29.6	-25	40	35	27	13	
35	115040	103730	10	35	16.7	-27	56	49	40	34	F
36	112300	259400	10	35	28.9	-25	02	34	20	16	
37	111696	154500	10	35	33.0	-27	00	01	27	20	
38	109720	75930	10	35	44.0	-28	27	55	47	40	p?
39	108540	113610	10	35	49.5	-27	45	45	31	22	
40	101570	137420	10	36	24.4	-27	19	05	37	20	
41	90460	192560	10	37	18.5	-26	17	19	37	34	
42	85000	53190	10	37	50.7	-28	53	15	34	20	B
43	82080	37330	10	38	06.1	-29	10	58	17	13	
44	81210	42160	10	38	10.5	-29	05	33	27	20	
45	74880	247690	10	38	34.2	-25	15	28	27	23	
46	65130	64040	10	39	32.0	-28	40	54	40	17	
47	60000	182270	10	39	51.2	-26	28	33	32	21	
48	59040	220140	10	39	53.8	-25	46	10	23	17	
49	56570	3700	10	40	18.8	-29	48	20	34	10	
50	52340	140670	10	40	32.2	-27	15	00	26	15	
51	51950	111030	10	40	36.0	-27	48	10	34	34	
52	50830	192340	10	40	36.4	-26	17	10	26	21	
53	48140	120530	10	40	54.7	-27	37	29	26	17	B
54	44700	258500	10	41	02.9	-25	03	01	60	47	p
55	29330	112960	10	42	30.3	-27	45	39	17	15	B
56	27420	89960	10	42	41.9	-28	11	21	21	11	B
57	25750	157600	10	42	56.8	-29	34	20	40	34	
58	13910	154140	10	43	44.4	-26	59	17	17	12	
59	8000	245000	10	44	05.3	-25	17	30	67	30	

contd...

Table 7.3 (contd.)

<u>No</u>	<u>X</u>	<u>Y</u>	<u><math>\alpha</math></u> (1950.0)			<u><math>\delta</math></u>			<u>D</u>	<u>d</u>	<u>Notes</u>
			<u>h</u>	<u>m</u>	<u>s</u>	<u>0</u>	<u>'</u>	<u>"</u>			
B. Elliptical or spheroidal systems											
60	144660	54440	10	32	45.45	-28	51	57	28	20	
61	117690	172690	10	35	02.8	-26	39	40	20	17	
62	89520	89140	10	37	26.6	-28	13	03	17	12	
63	85170	42200	10	37	50.2	-29	05	33	19	15	F
64	39890	158040	10	41	33.6	-26	55	23	36	11	

Columns: as for Table 7.2

Key to symbols

B = relatively bright

F = faint

p = peculiar

## CHAPTER 8

### CONCLUDING REMARKS

It is appropriate in this last chapter to review the facts that have been gleaned in this work and elsewhere concerning A1060, and to consider its present state and possible past and future evolution. Before recapitulating, the dynamics of the cluster will be considered in relation to the observed luminous mass.

#### 8.1 The Mass-to-Light Ratio

There are two methods by which the mass of a cluster might be determined, two other approaches to the problem having proved untenable (see Zwicky 1937). The first is to approximate the total mass of a cluster by the sum of the masses of the member galaxies, as estimated from their internal dynamics (rotation curves) or intrinsic luminosities given 'conventional' mass-to-light ratios. The second procedure involves the virial theorem for an ensemble of interacting point masses, and was first applied to the determination of cluster mass by Zwicky (1933).

Derivation of the virial theorem is straightforward: from Newton's laws, the force on the  $i^{\text{th}}$  galaxy whose coordinates with respect to the centre of mass of the system are  $\underline{r}_i$   $\underline{F}_i = m_i \underline{\ddot{r}}_i$ ; taking the scalar product with  $\underline{r}_i$  and summing over  $i$ ,

$$2T + \sum \underline{F}_i \cdot \underline{r}_i = \frac{1}{2} \ddot{I}$$

where  $I$  is the moment of inertia and  $T$  the kinetic energy of the



system. A number of assumptions are then made. Firstly,  $\underline{F}_i$  consists solely of an (internal) gravitational force; thus  $\sum \underline{F}_i \cdot \underline{r}_i = \Omega$ , yielding the theorem. If, secondly, the cluster is assumed to be in a steady state,

$$2\langle T \rangle + \langle \Omega \rangle = 0,$$

the "static virial". Thirdly, a smooth mass distribution is assumed in introducing the following approximations for  $T$  and  $\Omega$ :

$$2T = \sum m_i \dot{r}_i^2 = M\sigma^2$$

$$\Omega = -G \sum_{\text{pairs}} \frac{m_i m_j}{r_{ij}} = \frac{Gm m n(n-1)}{2R'} \approx \frac{-GM^2}{2R'}$$

where  $\sigma^2$  is the variance of the space velocities and  $2R' = R$  is the mean of the individual space separations  $r_{ij}$ , both mass-weighted.

In practice it is necessary to estimate  $R$  in terms of the projected separations (Limber & Matthews 1960) or the distribution of galaxies on the sky (Schwarzschild 1954), and  $\xi$  in terms of the radial velocity dispersion  $S$ :  $\sigma^2 = \xi S^2$ .  $\xi$  may take any value from  $\xi = 1$ , in which near-radial velocities predominate and only galaxies found near the cluster core are sampled, to  $\xi = 3$ , if the velocity field is everywhere isotropic. An additional uncertainty of a factor of two is introduced if the requirement that the cluster must be in a steady state is relaxed to the point where the cluster is on the verge of instability and the total energy  $T + \Omega$  is zero. The (static) virial

theorem thus predicts that

$$M = \xi S^2 R / G$$

where  $\frac{1}{2} \lesssim \xi \lesssim 3$  and  $M$  is the mass of the cluster.

The value of  $R$  for A1060 is uncertain. Strip counts of the distribution of prominent galaxies within an area of  $100 \text{ degree}^2$  that includes A1060 (Section 6.2) lead to the estimate,  $R_m = 4.0 \text{ Mpc}$  by Schwarzschild's (1954) method. This is much larger than the value  $R_m = 1.75 \text{ Mpc}$  ( $H_0 = 85 \text{ km s}^{-1} \text{ Mpc}^{-1}$ ) that has been derived for the Centaurus cluster by Dawe, Dickens & Peterson (1977), although it is less than that of the Coma cluster, for which  $R_m = 5.2 \text{ Mpc}$  on this distance scale (Rood et al. 1972). On the other hand, a fit of a de Vaucouleurs  $r^{1/4}$  luminosity distribution law to the radial profile of the effective surface brightness of the cluster (the latter is shown in Figure 6.7) furnishes an estimate of the effective radius,  $r_e = 0.53 \text{ Mpc}$ , which is very much smaller than  $R_m$ . The  $r_e$  derived in this manner is also applicable to virial theorem estimates of mass, in analogy to mass estimates for individual galaxies (as outlined in Young 1976).

Despite the fact that the fit of the luminosity profile of the cluster to an  $r^{1/4}$  law breaks down beyond  $0.9 \text{ Mpc}$ , it is felt that the  $r_e$  so derived is more representative of the size of Hydra I than is  $R_m$ . Only in the former case are the luminosities of the individual cluster members implicitly taken into account in the fit. Moreover, the core radius of this cluster is unusually small (Section 2.4).

Nevertheless, in view of the uncertainty regarding which value of  $R$  is more appropriate, it is preferable to adopt the compromise of Green (1978). He expresses virial properties in terms of the parameter  $\chi$ , where

$$\chi = \frac{3S^2}{GL} = \frac{M}{L} \frac{1}{R} .$$

Here  $L$  is the optical luminosity of the cluster and the factor, 3, is the adopted value of  $\xi$ .

A fair estimate of  $L$  may be obtained by summation of the luminosities of bright galaxies, and a linear extrapolation of the faint-end slope of the field-corrected luminosity function shown in Figure 6.4 to  $V_{25} = \infty$ . (As suggested in Chapter 7, the numbers of galaxies in the cluster may not continue to increase with magnitude faintwards of  $17^m$  at this exponential rate, but even if the assumption holds they contribute only 10% of the total  $V_{25}$  luminosity of the galaxies). Subtracting  $0.2^m$  from the total magnitude corresponding to the integrated luminosity defined above, to account for the light emission from those regions of the galaxies of surface brightness less than  $25 \text{ V mag arcsec}^{-2}$ , the corrected total luminosity of A1060 (out to a projected radius of 1.5 Mpc) is  $L = 1.26 \times 10^{12} L_{\odot}$ . With  $S = 496 \text{ km s}^{-1}$ ,  $\chi(\text{A1060}) = 136 M_{\odot} L_{\odot}^{-1} \text{ Mpc}^{-1}$ .

Where  $R = 1 \text{ Mpc}$ , the mean mass-to-light ratio of the galaxies in the cluster would be  $136 M_{\odot} L_{\odot}^{-1}$ , roughly an order of magnitude <sub>that</sub> larger than <sub>that</sub> which was assumed for elliptical galaxies for the purpose of mass-weighting of velocities in the determination of  $S$  (Chapter 3). Thus, the longstanding puzzle of the "hidden mass" in groups of

galaxies arises once again. A recent and careful review of current evidence (Faber & Gallagher 1979) favours the existence of such dark matter, although the question of what physical form it takes remains unresolved. If it was divided among the galaxies, more equipartition of energy would be expected than is observed; theoretical time scales for the mutual accretion and merging of galaxies would also become embarrassingly small - as indeed they may be already (cf. the high frequency of multiply-nucleated first-ranked galaxies in clusters that has been reported by Rood & Leir 1979). Conversely, however, if galaxies are only minute perturbations in the overall mass distribution in each cluster, it would be surprising to observe cannibalism proceeding at all.

Nonetheless, it should be emphasised that the magnitude of the apparent mass discrepancy in A1060 is not especially serious. If (i)  $R = r_e \approx 0.5$  Mpc, the value of  $\chi$  is halved; (ii) if near-radial velocities predominate in the cluster ( $\xi \approx 1$ ), a further factor-of-three reduction is appropriate. The central binary might be efficient at throwing cluster galaxies into radial orbits (cf. the binary model of the Coma cluster proposed by Byrd & Valtonen 1979); (iii) if significant gravitational binding energy is locked up in cluster binaries, this would mimic the effect of hidden mass (Wesson & Lermann 1977; Ozernoy & Reinhardt 1979). To investigate the latter two points, more radial velocities of members are required.

## 8.2 The Evolutionary State of A1060

In order to assist rapid evaluation, the main characteristics of the

Hydra I cluster are summarised in Table 8.1.

One of the more surprising results of the previous chapters is the degree of compactness of the core of this cluster; on the Emden sphere model the core radius of the number density distribution of galaxies is less than 100 kpc, one-third - or less - that of many other Abell clusters (Section 2.4). This core is highly symmetric, with an axial ratio of 0.9, and is aligned roughly parallel to the axis defined by the four bright and centrally located NGC galaxies (six if N3314 and N3316 are included). There are no statistically significant secondary maxima in the number density profile within 1.5 Mpc of the centre of the cluster, and the fit to an idealised unbounded isothermal sphere is good (but not necessarily physically significant). There are several recognisable groups of galaxies located within  $5^\circ$  of the centre of A1060, including Klemola 13 at an angular separation of only 0.8 Mpc, but no clear evidence has been revealed that might suggest physical association of Hydra I with other major galaxy groupings that are visible in the same region, namely the Hydra cloud and the Antlia group. To check this finding it may be necessary to investigate, on the basis of redshift data, the spatial structure of a region at least as large as that illustrated in Figure 2.1 (cf. Einasto, Joveer & Saar 1980).

The changes in the relative proportions of different types of galaxies with distance from the centre of A1060 is very marked; a comparable tendency for galaxies of different luminosity to segregate radially has not been confirmed, although the two brightest galaxies are centrally located in the cluster. It is difficult to account for



the radial segregation of spirals and SO galaxies entirely in terms of ram-pressure transmutation of the former type into the latter, since the possible examples of active transformation that have been found appear to represent only the early stages of the process. This finding may favour "independent formation" theories of the creation of lenticular galaxies over those which interpret these systems as spirals, less gas.

The luminosity function of the Hydra I cluster has been defined over a range of more than six magnitudes. The composite function for all galaxy types is similar in form to that of the Coma cluster (A1656, BM type II, RS type B), notwithstanding the differences in the relative proportions of galaxies of each type in the classified samples of members of each cluster. In both cases, the differential luminosity function exhibits an excess of numbers near the bright end followed by a local minimum (A426, as it happens, also exhibits the latter feature). On the other hand, the work of Dressler (1978b) and possibly of Bucknell et al. (1979) suggests significant variation of the overall forms of the luminosity functions of rich clusters.

The cannibalism hypothesis predicts that as a cluster of galaxies evolves, the absolute magnitude ( $M^*$ ) corresponding to the turnover point of the logarithmic cumulative luminosity function will gradually shift faintwards due to the proportionately greater depletion rate for members that have magnitudes close to  $M^*$ . Bucknell et al.'s (1979) data imply that it does not; so the apparently untypical value of  $M^*$  determined here for A1060 is more likely a statistical artefact than a consequence of evolution.



It was argued in Chapter 6 that the shallow slope of the (U - B) colour-magnitude correlation for the brighter members of A1060 could not be a consequence of widespread accretion of galaxies as implied by the Ostriker model. Michard's (1979a,b) work could be interpreted as evidence that the stellar composition of bright galaxies in clusters differs from the composition of the faint, though no explanation is ventured as to why this should be so.

Whilst the inconsistencies as regards A1060 suggests the need for some revision of the theory of evolution via galactic cannibalism, the positive aspects of the model deserve emphasis. Dynamical friction accounts in an attractive manner for the optical properties of N3311, the largest and brightest galaxy. In other, "cD" clusters, the presence of so-called "reverse" segregation (Dressler 1978b, Green 1978) - in which a deficiency of moderately bright galaxies exists in the cluster core, contrary to the expected consequences of dynamical relaxation through two-body encounters - is understandable in terms of galactic accretion, as is the observed tendency of many cD galaxies to align in sympathy with the galaxy distribution in their clusters. The importance of dynamical friction in spherical clusters has been stressed by White (1976a), and by Binney (1977) in the more general case of aspherical clusters.

The apparently normal integrated colours and colour gradients of N3311 imposes a strong limit on the amount of stellar debris that it could in principle have accreted, assuming that the accretion hypothesis is valid. On the evidence of this observation, and its optical characteristics, N3311 is regarded as a D system in the early stages of evolution.

The existence and properties of three anaemic spiral galaxies in the Hydra I cluster was discussed in Section 6.5. It was proposed that these systems are being depleted of their interstellar gas by the ram-pressure effect of their motion through an intracluster medium (the presence of the latter seems to be implied by observations of diffuse X-ray emission from the cluster). This, too, is believed to represent the early stages of a transitory phase.

On the evidence of the preceding, it is thought that A1060 is relatively young in dynamical terms. It seems entirely possible that N3311 will eventually evolve into a cD supergiant, at the expense of other members of the cluster; the compactness of the cluster core could encourage and accelerate this process. One other consequence of dynamical friction, as A1060 evolves (cf. Binney 1977) should be an increase in the flattening of the distribution of the bright galaxies at the centre of the cluster. Ultimately they may agglomerate into an enormous cD system.

### 8.3 Comparison with Other Clusters

The most striking single difference between A1060 and other nearby clusters for which compositions are quite well determined is the absence of spatial (as opposed to radial) segregation of different morphological types of galaxies. In the Centaurus I cluster (HMS 1247 - 4102: Dawe, Dickens & Peterson 1977; Fairall 1978a) - within which spiral SO and elliptical galaxies are present in proportions similar to those in A1060 - the spiral galaxies form an elongated and irregular distribution which contrasts with the more centrally

concentrated cloud of early types, and these clouds are  $\sim 0.3$  Mpc apart (if  $H_0 = 85 \text{ km s}^{-1} \text{ Mpc}^{-1}$ ).

The Virgo cluster has long been considered the classic example. de Vaucouleurs (1961b; see also de Vaucouleurs & de Vaucouleurs 1973) has interpreted this system as a superposition of a nearer cloud of early-type galaxies and a more distant cloud of spirals. Recent HI-profile observations (Helou, Salpeter & Krumm 1979), however, indicate that these clouds are concentric; in which case Virgo I would more nearly resemble Hydra I, although the latter cluster is more regular. Spiral galaxies located in the Virgo I cluster are already known to be deficient in neutral hydrogen (Davies & Lewis 1973), and it has been suggested that the gas loss is proceeding in the present era (Butcher & Oemler 1978).

Finally, in the Fornax cluster (G53 in the list of de Vaucouleurs, 1975) the most luminous galaxy, N1316, is located in a clump of mostly lenticulars and spiral galaxies; the next four brightest are in or near a larger concentration of mostly early types (Welch, Chincarini & Rood 1975).

The theme will not now be elaborated on because further details regarding two of the clusters, Centaurus and Fornax, are lacking. Comment will be limited to two remarks. Firstly, whilst it may be risky and irresponsible to generalise on the basis of so few examples, the features of the Centaurus and Fornax clusters may appear to favour White's (1976b) model of the formation of rich clusters, in which subclustering of galaxies develops rapidly during the expansion

of the initial cloud and anisotropies persist on recollapse. Secondly, the correspondence of some properties of Virgo I and Hydra I is close enough to warrant further investigation.

#### 8.4 The Way Ahead

The necessity and importance of further work on many aspects of the properties of the Hydra I cluster has been emphasised time and again in the preceding chapters.

The most immediately useful additional data on Hydra I that might be secured are redshifts for a larger sample of galaxies, so as better to elucidate the structure, velocity dispersion, and membership of A1060. It would also appear to be important to obtain HI-line observations of spiral and lenticular galaxies in the cluster. Observational programmes along these lines are already underway (Richter 1979, private communication).

In addition, spectrographic studies of the halo of N3311 would enable the velocity dispersion of stellar orbits in the outer regions of this galaxy to be estimated; a low value would rule out the presence of a massive envelope of luminous matter consisting of material stripped from other cluster members, the existence of which is already doubtful on photometric grounds.

The detailed study of nearby clusters of galaxies is only just beginning. A truly vast amount of research remains to be done, involving many disciplines and aimed towards a common goal: comprehension of the physical universe and our place within it.

Table 8.1

Summary of the properties of the Abell 1060 = Hydra I cluster

<u>Property</u>		<u>Reference</u>
1. Position		
equatorial (1950.0)	$\alpha = 10^{\text{h}}34^{\text{m}}17^{\text{s}} \pm 6^{\text{s}}$	
	$\delta = -27^{\circ}17' \pm 1'$	*
galactic	$l = 269.6$	
	$b = +26.5$	*
supergalactic	$\text{SGL} = 139.3$	
	$\text{SGB} = -37.6$	*
2. Identification as X-ray source		
3rd UHURU catalogue	3U 1044-30	Giacconi et al. (1974)
4th UHURU catalogue	4U 1033-26	Forman et al. (1978)
2nd Ariel V catalogue	2A 1033-270	Cooke et al. (1978)
3. X-ray luminosity	$L_{\text{x}} \approx 2 \times 10^{43} \text{ erg s}^{-1}$	Gursky & Schwartz (1977), Bahcall (1977a)
4. Redshift		
heliocentric	$z_{\text{O}} = 0.01221 \pm 0.00042(\text{s.e.})$	*
cosmological	$z_{\text{R}} = 0.01131 \pm 0.00043(\text{s.e.})$	*
5. Velocity dispersion	$S \equiv \langle v_{\text{G}}^2 \rangle^{\frac{1}{2}} = 496 \text{ km s}^{-1}$ ( = 771 km s <sup>-1</sup>	* adopted by Faber & Dressler (1976))
6. Distance	$\Delta = 40.5(1 \pm 0.10:) \text{ Mpc}$ ( = 40 Mpc adopted) ( = 68 Mpc if $H_0 = 50 \text{ km s}^{-1} \text{ Mpc}^{-1}$ )	de Vaucouleurs (1977b)
geometric distance modulus	$\mu_{\text{O}} = 33.0^{\text{m}} \pm 0.2^{\text{m}}$	de Vaucouleurs (1977b)
galactic absorption	$A_{\text{B}} = 0.38^{\text{m}} (A_{\text{V}} = 0.28^{\text{m}})$ $A_{\text{B}} = 0.16^{\text{m}} (A_{\text{V}} = 0.12^{\text{m}})$	RC2 Sandage (1973a)
K-correction	$K_{\text{V}} = 0.02^{\text{m}}$	Sandage (1973b)



Table 8.1 (contd.)

Property	Reference
7. Size	
core radius	$R_c = 0.08 \text{ Mpc}$ $( = 0.14 \text{ Mpc if}$ $H_o = 50 \text{ km s}^{-1} \text{ Mpc}^{-1})$ *
effective radius $r_e$	$= 0.53: \text{ Mpc}$ * $( = 0.90: \text{ Mpc if } H_o = 50)$
visible extent	$= 2 - 3.5 \text{ Mpc}$ * $( = 3.5 - 6 \text{ Mpc if } H_o = 50)$
8. Shape	
axial ratio	$0.87 \pm 0.02$ *
at position angle	(galaxies $V_{25} < 17$ ) $121^\circ$ *
9. Composition	
<1.5 Mpc radius,	$V_{25} < 15: \text{ E:S0:S \& Irr} = 15\%:46\%:39\%$ *
	(14 unclassified, total 102)
<3.5 Mpc radius,	$\text{E:S0:S \& Irr} = 6\%:13\%:81\%$ *
	(87 unclassified, total 459)
10. Luminosity distribution	
isophotal V magnitude at discontinuity in log. integrated luminosity	function $M_{V_{25}}^* \equiv M^* = -19.2$ ( $V_{25}^* = 14.08$ ) *
isophotal V magnitude of first-ranked galaxy (N3311)	$M_1 = -22.6$ ( $V_{25,1} = 10.65$ ) *
magnitude difference between first and second-ranked galaxies:	isophotal, V-band $ M_1 - M_2  = 0.85$ *
	total, B-band $ B_{T,1} - B_{T,2}  = 0.74:$ *



Table 8.1 (contd.)

<u>Property</u>		<u>Reference</u>
11. Mass-to-light ratio	$\chi \equiv \frac{M}{L} \frac{1}{R} = 136 M_{\odot} L_{\odot}^{-1} \text{Mpc}^{-1}$	*
12. Classification		
Bautz-Morgan	BM = III	Bautz & Morgan (1970)
Rood-Sastry	RS = C	Rood & Sastry (1971)
Bahcall	Intermediate	*
Other	spiral cluster	Melnick & Sargent (1977)

\* denotes present study

APPENDIX IFINDING CHARTS OF GALAXIES

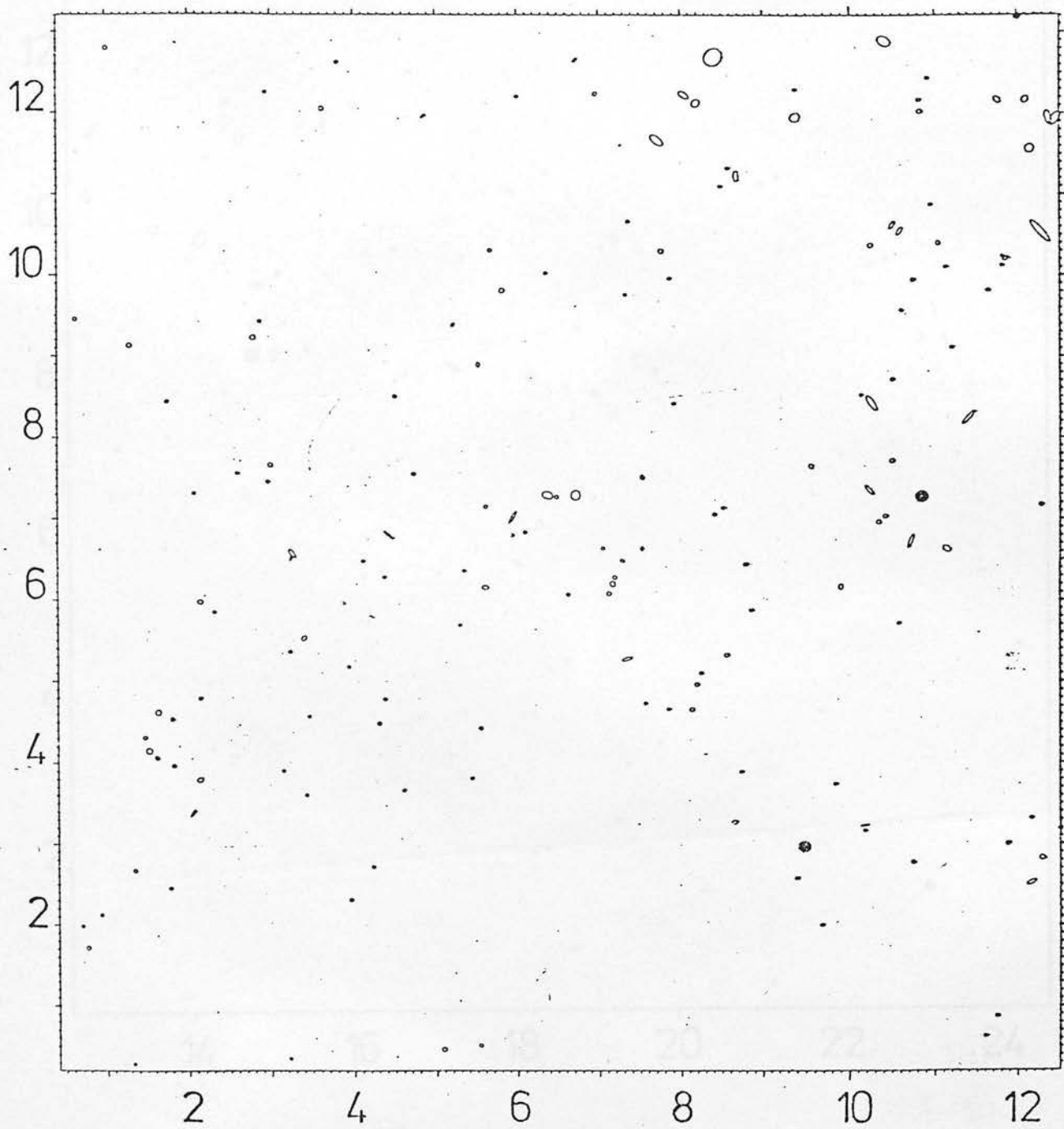
The first four of the following charts record all galaxies near A1060 of  $V_{25}^c < 17.0$  (i.e., those galaxies with photographic magnitudes interior to the 25th V mag arcsec<sup>-2</sup> isophote, corrected for saturation effects and for image overlaps, brighter than 17<sup>m</sup>). Each chart represents a different quarter of a region of total area 21.8 degree<sup>2</sup> approximately).

The fifth chart depicts systems brighter than  $V_{25} = 17.5$  in the central zone of the cluster. The area shown is approximately 0.87 degree<sup>2</sup>.

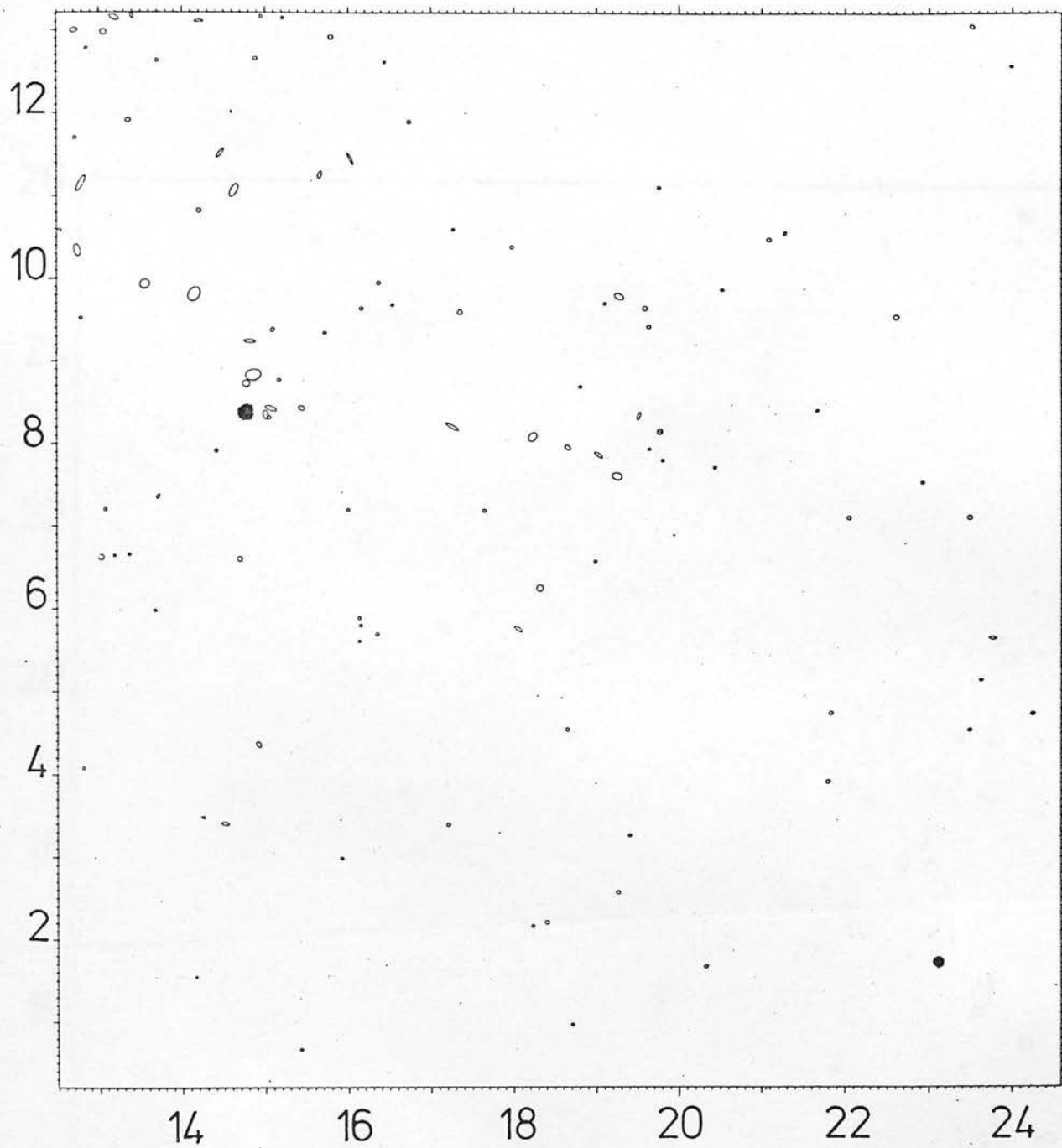
Where the orientation and apparent (micrometric) axis dimensions of a galaxy are known, these are indicated; otherwise each galaxy is represented by a circular disk of area proportional to the luminosity corresponding to the  $V_{25}^c$ . Galaxies listed in the Appendix II Supplement are represented by filled circles (not to scale).

Axes are in the rectangular coordinate system of Appendix II and are labelled in centimetres (1 cm = 671.4 arcseconds).

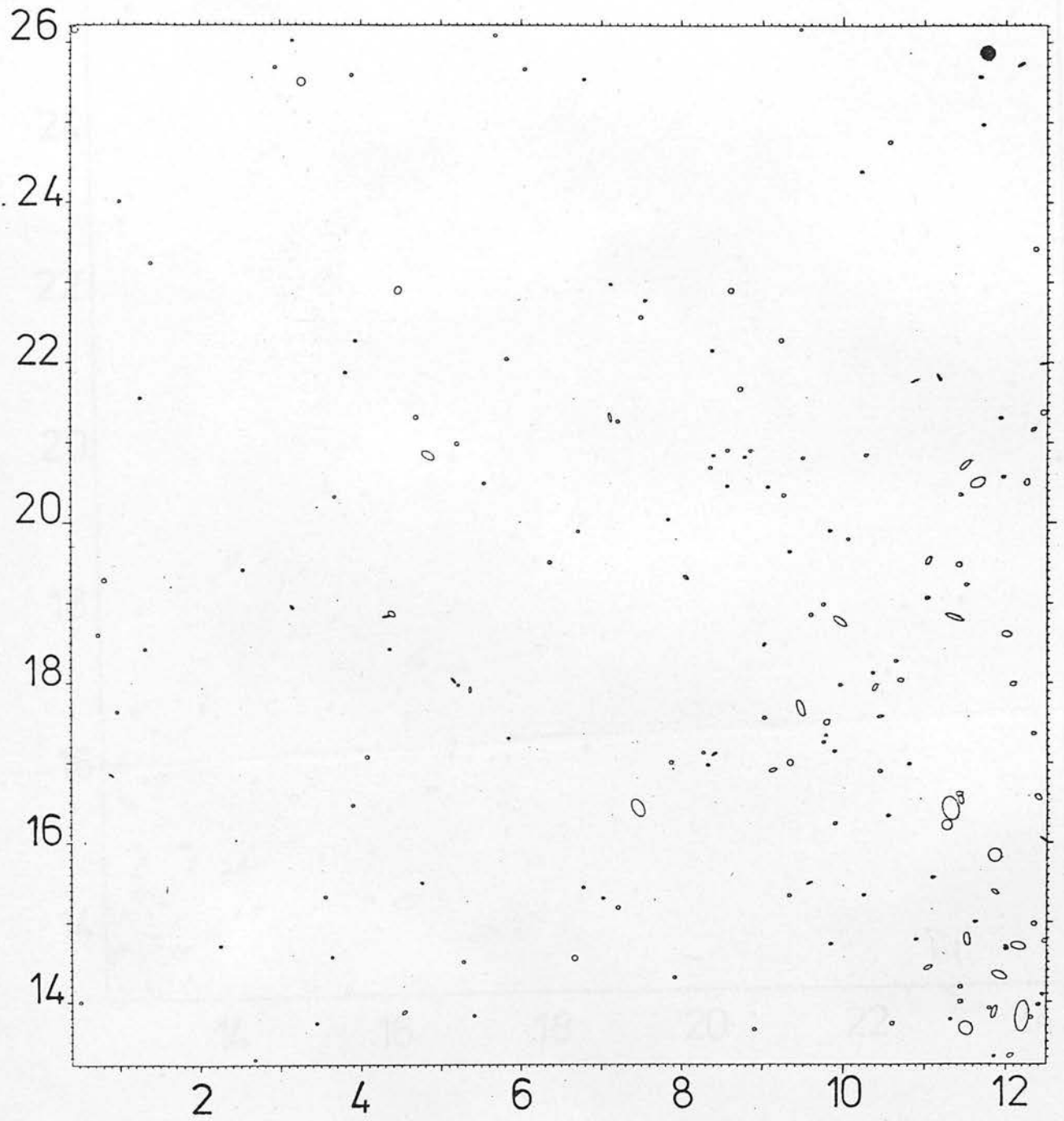
North is to the top, east to the left.



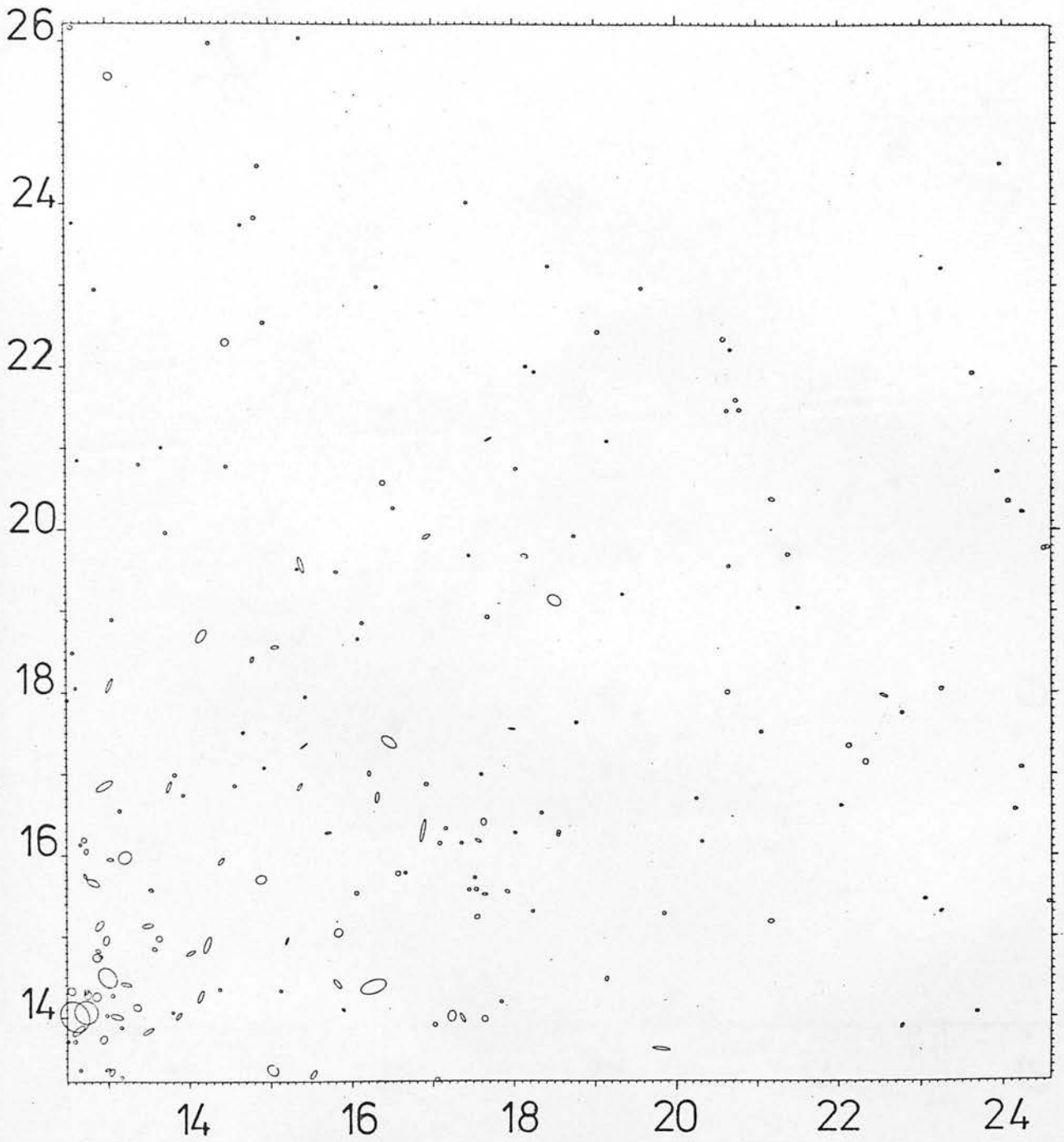
South east;  $V_{25} < 17.0$



South west;  $V_{25} < 17.0$

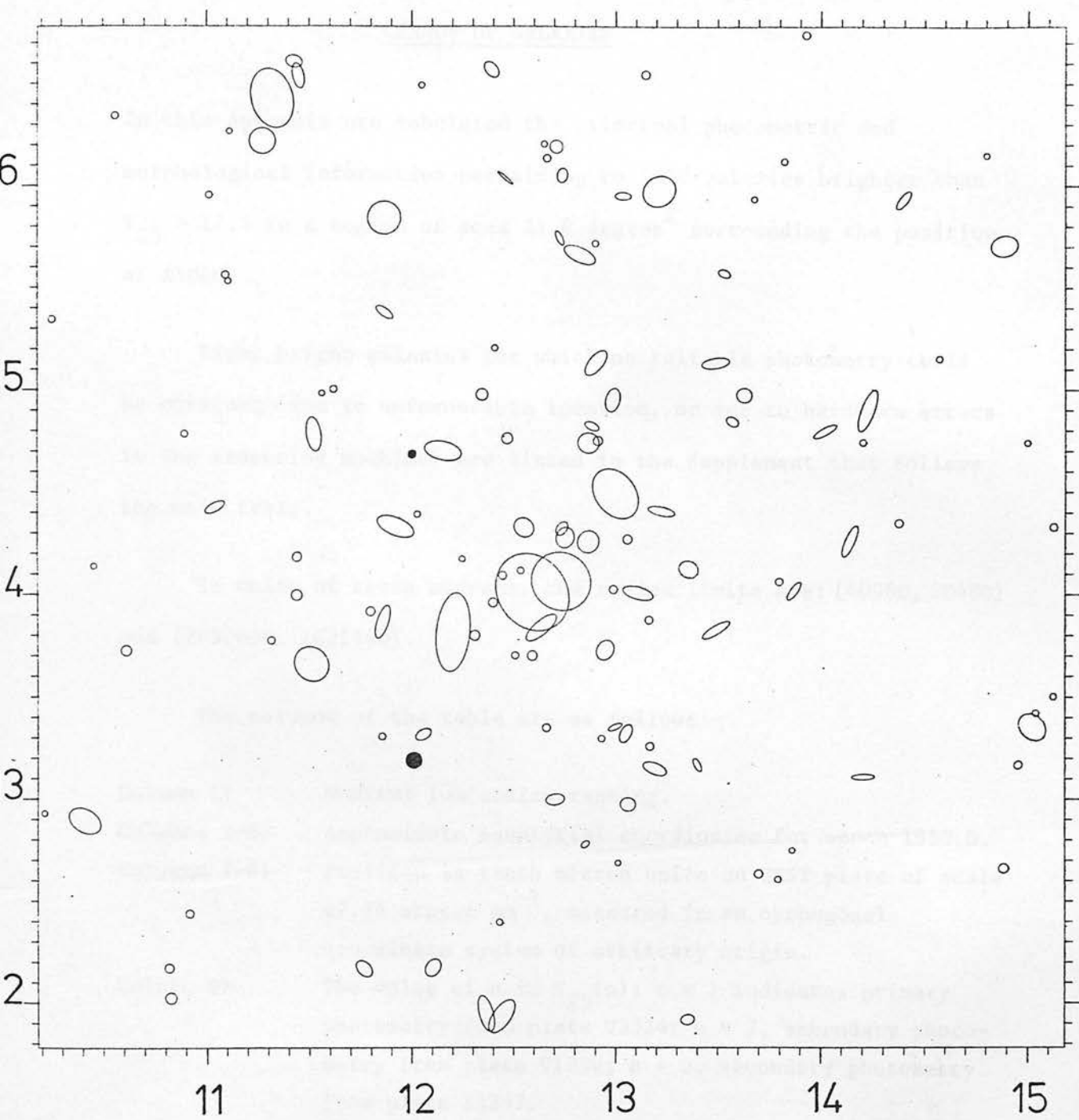


North east;  $V_{25} < 17.0$



North west;  $V_{25} < 17.0$





Centre;  $V_{25} < 17.5$

10'

APPENDIX IICENSUS OF GALAXIES

In this Appendix are tabulated the principal photometric and morphological information pertaining to 1000 galaxies brighter than  $V_{25} = 17.5$  in a region of area  $21.8 \text{ degree}^2$  surrounding the position of A1060.

Eight bright galaxies for which no reliable photometry could be obtained (due to unfavourable location, or due to hardware errors in the measuring machine) are listed in the Supplement that follows the main table.

In units of tenth microns, the region limits are: (40960, 20480) and (2457600, 2621440).

The columns of the table are as follows:-

Column 1)	Nominal luminosity ranking.
Columns 2-6)	Approximate equatorial coordinates for epoch 1950.0.
Columns 7-8)	Position in tenth micron units on UKST plate of scale $67.14 \text{ arcsec mm}^{-1}$ , measured in an orthogonal coordinate system of arbitrary origin.
Column 9)	The value of $n$ in $V_{25}(n)$ : $n = 1$ indicates primary photometry from plate V3324; $n = 2$ , secondary photometry from plate V1330; $n = 3$ , secondary photometry from plate R1347.
Column 10)	$V_{25}^c$ , magnitudes.
Column 11)	$r_{25}^*$ , arcseconds (V band).
Column 12)	$B_{25}^c$ , magnitudes.
Columns 13+)	Morphological classification in the revised Hubble classification system (e.g. de Vaucouleurs 1959). If the image is a blend, this is noted thus: $mGnS = m$ galaxies + $n$ stars (last column).

## Standard designations of brighter galaxies:

<u>Galaxy</u> <u>No</u>	<u>RNGC</u>	<u>IC</u>	<u>MCG</u>	<u>Anon.</u>
1	3311		-4-25-36	
2	3309		-4125-34	
3	3312	629	-4-25-39	
4		2597	-4-25-51	
5	3308		-4-25-32	
6	3285		-4-25-19	
7	3336		-5-25-36	
9	3316		-4-25-46	
12	3305		-4-25-31	
20	3285B		-4-25-22	
24	3315		-4-25-43	A1034-27B; A1034-2712
32			-4-25-25	A1033-27
34	3285A		-4-25-15	
37				A1034-27A; A1034-2718
42			-5-25-32	A1037-27
60	3307		-4-25-29	
70				A1033-2715
A	3313		-4-25-44	
B	3314		-4-25-41	

The MCG gives designations for further galaxies in this field.

RNGC = Revised New General Catalogue of Nonstellar Astronomical Objects (Sulentic & Tifft 1973).

IC = Index Catalogue of Nebulae (Dreyer 1895); Second Index Catalogue of Nebulae and Clusters of Stars (Dreyer 1908).

MCG = Morphological Catalogue of Galaxies (Vorontsov-Velyaminov, Krasnogorskaja & Arkipova 1962-8).

Anon = as in the Second Reference Catalogue of Bright Galaxies (de Vaucouleurs, de Vaucouleurs & Corwin 1976), and Sandage (1978).

Listed below are the saturation-corrected isophotal magnitudes of nineteen galaxies for which photoelectric photometry is available. These photoelectric standards defined the photographic saturation corrections of section 5.7. The magnitudes given here supersede the values in the main table (overleaf).

The magnitudes in the main table have been adjusted on the basis of the mean curves of corrections shown in Figures 5.10 and 5.11.

<u>Galaxy No</u>	<u><math>V_{25}^c</math></u>	<u><math>B_{25}^c</math></u>	$r_{25}^* = 237''$
1	10.65	12.11	
2	11.50	12.73	
3	11.88	12.79	
4	11.83	12.89	
5	12.13	13.25	
12	12.91	14.13	
22	13.41	14.49	
24	13.58	14.75	
31	13.92	15.11	
37	13.91	14.97	
38	13.69	14.79	
46	14.01	15.18	
56	14.28	15.41	
61	13.99	15.27	
63	14.36	15.67	
66	14.16	15.86	
68	14.41	15.57	
86	14.77	15.65	
92	14.82	15.87	

IO NO	R-A	M	S	DEC	X	Y	V25	R25	B25	TYPE	PEC
1	34	10	1	15	1259	1401	1	170.57	12.26	S	PEC
2	34	10	1	15	1273	1390	1	119.56	12.78	S	PEC
3	34	10	1	15	1273	1700	1	85.54	12.89	S	PEC
4	34	10	1	15	1273	3351	1	85.72	13.20	S	PEC
5	34	10	1	15	1273	5100	1	84.40	13.02	S	PEC
6	34	10	1	15	1273	5200	1	71.66	13.92	S	PEC
7	34	10	1	15	1273	4650	1	50.75	12.97	S	PEC
8	34	10	1	15	1273	0000	1	58.01	13.93	S	PEC
9	34	10	1	15	1273	1305	1	50.78	14.08	S	PEC
10	34	10	1	15	1273	3340	1	47.11	13.90	S	PEC
11	34	10	1	15	1273	4600	1	43.77	14.32	S	PEC
12	34	10	1	15	1273	4429	1	44.44	14.32	S	PEC
13	34	10	1	15	1273	5000	1	37.63	13.87	S	PEC
14	34	10	1	15	1273	5200	1	39.20	14.49	S	PEC
15	34	10	1	15	1273	3375	1	38.57	14.33	S	PEC
16	34	10	1	15	1273	5000	1	36.35	14.51	S	PEC
17	34	10	1	15	1273	0000	1	31.97	15.17	S	PEC
18	34	10	1	15	1273	3300	1	33.17	14.69	S	PEC
19	34	10	1	15	1273	5000	1	33.21	14.27	S	PEC
20	34	10	1	15	1273	5200	1	24.62	15.01	S	PEC
21	34	10	1	15	1273	3375	1	32.53	14.91	S	PEC
22	34	10	1	15	1273	5000	1	27.06	14.53	S	PEC
23	34	10	1	15	1273	0000	1	26.85	14.58	S	PEC
24	34	10	1	15	1273	3300	1	22.86	14.93	S	PEC
25	34	10	1	15	1273	5000	1	33.54	14.93	S	PEC
26	34	10	1	15	1273	5200	1	37.58	15.03	S	PEC
27	34	10	1	15	1273	3375	1	25.69	14.86	S	PEC
28	34	10	1	15	1273	5000	1	25.94	14.98	S	PEC
29	34	10	1	15	1273	0000	1	25.94	14.98	S	PEC
30	34	10	1	15	1273	3300	1	25.04	15.01	S	PEC
31	34	10	1	15	1273	5000	1	33.10	14.66	S	PEC
32	34	10	1	15	1273	0000	1	33.75	15.21	S	PEC
33	34	10	1	15	1273	3300	1	35.04	15.33	S	PEC
34	34	10	1	15	1273	5000	1	38.03	15.33	S	PEC
35	34	10	1	15	1273	0000	1	25.29	15.03	S	PEC
36	34	10	1	15	1273	3300	1	30.88	15.16	S	PEC
37	34	10	1	15	1273	5000	1	30.88	15.16	S	PEC
38	34	10	1	15	1273	0000	1	30.88	15.16	S	PEC
39	34	10	1	15	1273	3300	1	30.88	15.16	S	PEC
40	34	10	1	15	1273	5000	1	30.88	15.16	S	PEC



IO NO	R.A.	(1950.0) DEC.	X	Y	V25	R25	B25	TYPE	
1	34	0	0.5	70	13	38.30	15.86	SA (RS)	162S
2	34	0	0.5	70	14	29.76	15.34	SA (RS)	
3	34	0	0.5	70	14			SA	
4	34	0	0.5	70	14			SA	
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99	34	0	0.5	70	14			SA	
100	34	0	0.5	70	14			SA	



IO	ND	R.A.	(1950.0)	DEC.	X	Y	V25	R25	B25	TYPE
101	101	37	33	43	933855	17004535	14.99	24.16	16.04	(R) SB (RS) D
102	102	34	34	45	932575	1415185	14.99	17.13	15.94	SR (R) SP
103	103	34	34	47	933907	1252415	15.00	20.99	15.72	S D
104	104	34	34	49	930533	1656255	15.00	29.21	17.15	SR (R) D
105	105	34	34	51	932233	1329765	15.00	21.50	16.17	SR (R) E3
106	106	34	34	53	932338	1205081	15.00	20.56	16.11	SP SP
107	107	34	34	55	930533	1252415	15.00	18.67	15.86	SR (S) D
108	108	34	34	57	930533	1252415	15.00	25.20	16.21	S D
109	109	34	34	59	930533	1252415	15.00	23.04	15.43	S D
110	110	34	34	61	930533	1252415	15.00	17.57	15.80	S D
111	111	34	34	63	930533	1252415	15.00	23.09	16.13	S D
112	112	34	34	65	930533	1252415	15.00	15.88	16.17	SA (R) D
113	113	34	34	67	930533	1252415	15.00	20.54	16.07	S D
114	114	34	34	69	930533	1252415	15.00	19.18	16.32	S D
115	115	34	34	71	930533	1252415	15.00	19.47	16.32	S D
116	116	34	34	73	930533	1252415	15.00	15.74	16.02	S D
117	117	34	34	75	930533	1252415	15.00	16.67	16.41	S D
118	118	34	34	77	930533	1252415	15.00	18.94	16.41	S D
119	119	34	34	79	930533	1252415	15.00	15.76	16.17	S D
120	120	34	34	81	930533	1252415	15.00	17.06	16.22	S D
121	121	34	34	83	930533	1252415	15.00	17.97	16.48	S D
122	122	34	34	85	930533	1252415	15.00	15.76	16.17	S D
123	123	34	34	87	930533	1252415	15.00	17.06	16.22	S D
124	124	34	34	89	930533	1252415	15.00	17.97	16.48	S D
125	125	34	34	91	930533	1252415	15.00	15.76	16.17	S D
126	126	34	34	93	930533	1252415	15.00	17.06	16.22	S D
127	127	34	34	95	930533	1252415	15.00	17.97	16.48	S D
128	128	34	34	97	930533	1252415	15.00	15.76	16.17	S D
129	129	34	34	99	930533	1252415	15.00	17.06	16.22	S D
130	130	34	34	101	930533	1252415	15.00	17.97	16.48	S D
131	131	34	34	103	930533	1252415	15.00	15.76	16.17	S D
132	132	34	34	105	930533	1252415	15.00	17.06	16.22	S D
133	133	34	34	107	930533	1252415	15.00	17.97	16.48	S D
134	134	34	34	109	930533	1252415	15.00	15.76	16.17	S D
135	135	34	34	111	930533	1252415	15.00	17.06	16.22	S D
136	136	34	34	113	930533	1252415	15.00	17.97	16.48	S D
137	137	34	34	115	930533	1252415	15.00	15.76	16.17	S D
138	138	34	34	117	930533	1252415	15.00	17.06	16.22	S D
139	139	34	34	119	930533	1252415	15.00	17.97	16.48	S D
140	140	34	34	121	930533	1252415	15.00	15.76	16.17	S D
141	141	34	34	123	930533	1252415	15.00	17.06	16.22	S D
142	142	34	34	125	930533	1252415	15.00	17.97	16.48	S D
143	143	34	34	127	930533	1252415	15.00	15.76	16.17	S D
144	144	34	34	129	930533	1252415	15.00	17.06	16.22	S D
145	145	34	34	131	930533	1252415	15.00	17.97	16.48	S D
146	146	34	34	133	930533	1252415	15.00	15.76	16.17	S D
147	147	34	34	135	930533	1252415	15.00	17.06	16.22	S D
148	148	34	34	137	930533	1252415	15.00	17.97	16.48	S D
149	149	34	34	139	930533	1252415	15.00	15.76	16.17	S D
150	150	34	34	141	930533	1252415	15.00	17.06	16.22	S D
151	151	34	34	143	930533	1252415	15.00	17.97	16.48	S D
152	152	34	34	145	930533	1252415	15.00	15.76	16.17	S D
153	153	34	34	147	930533	1252415	15.00	17.06	16.22	S D
154	154	34	34	149	930533	1252415	15.00	17.97	16.48	S D
155	155	34	34	151	930533	1252415	15.00	15.76	16.17	S D
156	156	34	34	153	930533	1252415	15.00	17.06	16.22	S D
157	157	34	34	155	930533	1252415	15.00	17.97	16.48	S D
158	158	34	34	157	930533	1252415	15.00	15.76	16.17	S D
159	159	34	34	159	930533	1252415	15.00	17.06	16.22	S D
160	160	34	34	161	930533	1252415	15.00	17.97	16.48	S D
161	161	34	34	163	930533	1252415	15.00	15.76	16.17	S D
162	162	34	34	165	930533	1252415	15.00	17.06	16.22	S D
163	163	34	34	167	930533	1252415	15.00	17.97	16.48	S D
164	164	34	34	169	930533	1252415	15.00	15.76	16.17	S D
165	165	34	34	171	930533	1252415	15.00	17.06	16.22	S D
166	166	34	34	173	930533	1252415	15.00	17.97	16.48	S D
167	167	34	34	175	930533	1252415	15.00	15.76	16.17	S D
168	168	34	34	177	930533	1252415	15.00	17.06	16.22	S D
169	169	34	34	179	930533	1252415	15.00	17.97	16.48	S D
170	170	34	34	181	930533	1252415	15.00	15.76	16.17	S D
171	171	34	34	183	930533	1252415	15.00	17.06	16.22	S D
172	172	34	34	185	930533	1252415	15.00	17.97	16.48	S D
173	173	34	34	187	930533	1252415	15.00	15.76	16.17	S D
174	174	34	34	189	930533	1252415	15.00	17.06	16.22	S D
175	175	34	34	191	930533	1252415	15.00	17.97	16.48	S D
176	176	34	34	193	930533	1252415	15.00	15.76	16.17	S D
177	177	34	34	195	930533	1252415	15.00	17.06	16.22	S D
178	178	34	34	197	930533	1252415	15.00	17.97	16.48	S D
179	179	34	34	199	930533	1252415	15.00	15.76	16.17	S D
180	180	34	34	201	930533	1252415	15.00	17.06	16.22	S D
181	181	34	34	203	930533	1252415	15.00	17.97	16.48	S D
182	182	34	34	205	930533	1252415	15.00	15.76	16.17	S D
183	183	34	34	207	930533	1252415	15.00	17.06	16.22	S D
184	184	34	34	209	930533	1252415	15.00	17.97	16.48	S D
185	185	34	34	211	930533	1252415	15.00	15.76	16.17	S D
186	186	34	34	213	930533	1252415	15.00	17.06	16.22	S D
187	187	34	34	215	930533	1252415	15.00	17.97	16.48	S D
188	188	34	34	217	930533	1252415	15.00	15.76	16.17	S D
189	189	34	34	219	930533	1252415	15.00	17.06	16.22	S D
190	190	34	34	221	930533	1252415	15.00	17.97	16.48	S D
191	191	34	34	223	930533	1252415	15.00	15.76	16.17	S D
192	192	34	34	225	930533	1252415	15.00	17.06	16.22	S D
193	193	34	34	227	930533	1252415	15.00	17.97	16.48	S D
194	194	34	34	229	930533	1252415	15.00	15.76	16.17	S D
195	195	34	34	231	930533	1252415	15.00	17.06	16.22	S D
196	196	34	34	233	930533	1252415	15.00	17.97	16.48	S D
197	197	34	34	235	930533	1252415	15.00	15.76	16.17	S D
198	198	34	34	237	930533	1252415	15.00	17.06	16.22	S D
199	199	34	34	239	930533	1252415	15.00	17.97	16.48	S D
200	200	34	34	241	930533	1252415	15.00	15.76	16.17	S D

IO NO	R-A	H	S	(1950-01 DEC)	X	Y	V25	R25	B25	TYPE
1	33	4	17	0	138684	15	15	15.23	16.43	E2
2	33	4	16	16	179554	20	15	21.66	16.26	S
3	33	4	14	33	114305	24	15	21.25	16.73	E3
4	33	4	14	33	116041	10	15	19.63	16.49	E3
5	33	4	12	33	138403	5	15	19.09	16.45	SA
6	33	4	12	33	107116	15	15			S
7	33	4	12	33	156010	0	15	13.03	16.51	SR
8	33	4	12	33	174590	5	15			SR
9	33	4	12	33	125950	5	15			PEC
10	33	4	12	33	125950	5	15			SP
11	33	4	12	33	125950	5	15			161S
12	33	4	12	33	125950	5	15			161S
13	33	4	12	33	125950	5	15			161S
14	33	4	12	33	125950	5	15			161S
15	33	4	12	33	125950	5	15			161S
16	33	4	12	33	125950	5	15			161S
17	33	4	12	33	125950	5	15			161S
18	33	4	12	33	125950	5	15			161S
19	33	4	12	33	125950	5	15			161S
20	33	4	12	33	125950	5	15			161S
21	33	4	12	33	125950	5	15			161S
22	33	4	12	33	125950	5	15			161S
23	33	4	12	33	125950	5	15			161S
24	33	4	12	33	125950	5	15			161S
25	33	4	12	33	125950	5	15			161S
26	33	4	12	33	125950	5	15			161S
27	33	4	12	33	125950	5	15			161S
28	33	4	12	33	125950	5	15			161S
29	33	4	12	33	125950	5	15			161S
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32	33	4	12	33	125950	5	15			161S
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35	33	4	12	33	125950	5	15			161S
36	33	4	12	33	125950	5	15			161S
37	33	4	12	33	125950	5	15			161S
38	33	4	12	33	125950	5	15			161S
39	33	4	12	33	125950	5	15			161S
40	33	4	12	33	125950	5	15			161S
41	33	4	12	33	125950	5	15			161S
42	33	4	12	33	125950	5	15			161S
43	33	4	12	33	125950	5	15			161S
44	33	4	12	33	125950	5	15			161S
45	33	4	12	33	125950	5	15			161S
46	33	4	12	33	125950	5	15			161S
47	33	4	12	33	125950	5	15			161S
48	33	4	12	33	125950	5	15			161S
49	33	4	12	33	125950	5	15			161S
50	33	4	12	33	125950	5	15			161S

IO	NO	R-A	H	M	S	(1950-01 DEC)	X	Y	V25	R25	B25	TYPE
1	1	34	10	0	0	0	0	14899785	15.83	16.74	16.97	SR (S)
2	2	35	10	0	0	0	0	14399710	15.83	17.43	16.66	SR (S)
3	3	36	10	0	0	0	0	15647370	15.84			SR (S)
4	4	45	10	0	0	0	0	12710680	15.84			SR (S)
5	5	44	10	0	0	0	0	13907070	15.85	12.15	16.46	SP
6	6	33	10	0	0	0	0	15722655	15.86			SP
7	7	33	10	0	0	0	0	15722655	15.86			SP
8	8	33	10	0	0	0	0	15722655	15.86			SP
9	9	33	10	0	0	0	0	15722655	15.86			SP
10	10	33	10	0	0	0	0	15722655	15.86			SP
11	11	33	10	0	0	0	0	15722655	15.86			SP
12	12	33	10	0	0	0	0	15722655	15.86			SP
13	13	33	10	0	0	0	0	15722655	15.86			SP
14	14	33	10	0	0	0	0	15722655	15.86			SP
15	15	33	10	0	0	0	0	15722655	15.86			SP
16	16	33	10	0	0	0	0	15722655	15.86			SP
17	17	33	10	0	0	0	0	15722655	15.86			SP
18	18	33	10	0	0	0	0	15722655	15.86			SP
19	19	33	10	0	0	0	0	15722655	15.86			SP
20	20	33	10	0	0	0	0	15722655	15.86			SP
21	21	33	10	0	0	0	0	15722655	15.86			SP
22	22	33	10	0	0	0	0	15722655	15.86			SP
23	23	33	10	0	0	0	0	15722655	15.86			SP
24	24	33	10	0	0	0	0	15722655	15.86			SP
25	25	33	10	0	0	0	0	15722655	15.86			SP
26	26	33	10	0	0	0	0	15722655	15.86			SP
27	27	33	10	0	0	0	0	15722655	15.86			SP
28	28	33	10	0	0	0	0	15722655	15.86			SP
29	29	33	10	0	0	0	0	15722655	15.86			SP
30	30	33	10	0	0	0	0	15722655	15.86			SP
31	31	33	10	0	0	0	0	15722655	15.86			SP
32	32	33	10	0	0	0	0	15722655	15.86			SP
33	33	33	10	0	0	0	0	15722655	15.86			SP
34	34	33	10	0	0	0	0	15722655	15.86			SP
35	35	33	10	0	0	0	0	15722655	15.86			SP
36	36	33	10	0	0	0	0	15722655	15.86			SP
37	37	33	10	0	0	0	0	15722655	15.86			SP
38	38	33	10	0	0	0	0	15722655	15.86			SP
39	39	33	10	0	0	0	0	15722655	15.86			SP
40	40	33	10	0	0	0	0	15722655	15.86			SP
41	41	33	10	0	0	0	0	15722655	15.86			SP
42	42	33	10	0	0	0	0	15722655	15.86			SP
43	43	33	10	0	0	0	0	15722655	15.86			SP
44	44	33	10	0	0	0	0	15722655	15.86			SP
45	45	33	10	0	0	0	0	15722655	15.86			SP
46	46	33	10	0	0	0	0	15722655	15.86			SP
47	47	33	10	0	0	0	0	15722655	15.86			SP
48	48	33	10	0	0	0	0	15722655	15.86			SP
49	49	33	10	0	0	0	0	15722655	15.86			SP
50	50	33	10	0	0	0	0	15722655	15.86			SP
51	51	33	10	0	0	0	0	15722655	15.86			SP
52	52	33	10	0	0	0	0	15722655	15.86			SP
53	53	33	10	0	0	0	0	15722655	15.86			SP
54	54	33	10	0	0	0	0	15722655	15.86			SP
55	55	33	10	0	0	0	0	15722655	15.86			SP
56	56	33	10	0	0	0	0	15722655	15.86			SP
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99	99	33	10	0	0	0	0	15722655	15.86			SP
100	100	33	10	0	0	0	0	15722655	15.86			SP



10	ND	R-A	H	M	S	DEC.	X	Y	V25	R25	B25	TYPE
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4	25	3	1	0	0	1950	9448	2703	1			
5	25	3	1	0	0	1950	6438	7270	1			
6	25	3	1	0	0	1950	3393	1378	1			
7	25	3	1	0	0	1950	2722	1743	1			
8	25	3	1	0	0	1950	2079	1574	1			
9	25	3	1	0	0	1950	1571	1627	1			
0	25	3	1	0	0	1950	7521	1625	1			
1	25	3	1	0	0	1950	2356	1803	1			
2	25	3	1	0	0	1950	5668	1030	1			
3	25	3	1	0	0	1950	1478	7131	1			
4	25	3	1	0	0	1950	1768	1055	1			
5	25	3	1	0	0	1950	7682	1420	1			
6	25	3	1	0	0	1950	1738	1141	1			
7	25	3	1	0	0	1950	1147	1555	1			
8	25	3	1	0	0	1950	7773	1440	1			
9	25	3	1	0	0	1950	1503	1377	1			
0	25	3	1	0	0	1950	8390	4901	1			
1	25	3	1	0	0	1950	7390	8658	1			
2	25	3	1	0	0	1950	1830	1933	1			
3	25	3	1	0	0	1950	5053	1875	1			
4	25	3	1	0	0	1950	1085	1471	1			
5	25	3	1	0	0	1950	2483	1875	1			
6	25	3	1	0	0	1950	4283	1710	1			
7	25	3	1	0	0	1950	3353	1922	1			
8	25	3	1	0	0	1950	1325	1889	1			
9	25	3	1	0	0	1950	4370	1697	1			
0	25	3	1	0	0	1950	3743	2547	1			
1	25	3	1	0	0	1950	2550	1907	1			
2	25	3	1	0	0	1950	1377	1742	1			
3	25	3	1	0	0	1950	4772	1899	1			
4	25	3	1	0	0	1950	2595	1630	1			
5	25	3	1	0	0	1950	7747	1890	1			
6	25	3	1	0	0	1950	2588	1630	1			
7	25	3	1	0	0	1950	4772	1890	1			
8	25	3	1	0	0	1950	2595	1630	1			
9	25	3	1	0	0	1950	7747	1890	1			
0	25	3	1	0	0	1950	2588	1630	1			
1	25	3	1	0	0	1950	4772	1890	1			
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3	25	3	1	0	0	1950	7747	1890	1			
4	25	3	1	0	0	1950	2588	1630	1			
5	25	3	1	0	0	1950	4772	1890	1			
6	25	3	1	0	0	1950	2595	1630	1			
7	25	3	1	0	0	1950	7747	1890	1			
8	25	3	1	0	0	1950	2588	1630	1			
9	25	3	1	0	0	1950	4772	1890	1			
0	25	3	1	0	0	1950	2595	1630	1			
1	25	3	1	0	0	1950	7747	1890	1			
2	25	3	1	0	0	1950	2588	1630	1			
3	25	3	1	0	0	1950	4772	1890	1			
4	25	3	1	0	0	1950	2595	1630	1			
5	25	3	1	0	0	1950	7747	1890	1			
6	25	3	1	0	0	1950	2588	1630	1			
7	25	3	1	0	0	1950	4772	1890	1			
8	25	3	1	0	0	1950	2595	1630	1			
9	25	3	1	0	0	1950	7747	1890	1			
0	25	3	1	0	0	1950	2588	1630	1			
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3	25	3	1	0	0	1950	7747	1890	1			
4	25	3	1	0	0	1950	2588	1630	1			
5	25	3	1	0	0	1950	4772	1890	1			
6	25	3	1	0	0	1950	2595	1630	1			
7	25	3	1	0	0	1950	7747	1890	1			
8	25	3	1	0	0	1950	2588	1630	1			
9	25	3	1	0	0	1950	4772	1890	1			
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2	25	3	1	0	0	1950	2588	1630	1			
3	25	3	1	0	0	1950	4772	1890	1			
4	25	3	1	0	0	1950	2595	1630	1			
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6	25	3	1	0	0	1950	2588	1630	1			
7	25	3	1	0	0	1950	4772	1890	1			
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9	25	3	1	0	0	1950	7747	1890	1			
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3	25	3	1	0	0	1950	7747	1890	1			
4	25	3	1	0	0	1950	2588	1630	1			
5	25	3	1	0	0	1950	4772	1890	1			
6	25	3	1	0	0	1950	2595	1630	1			
7	25	3	1	0	0	1950	7747	1890	1			
8	25	3	1	0	0	1950	2588	1630	1			
9	25	3	1	0	0	1950	4772	1890	1			
0	25	3	1	0	0	1950	2595	1630	1			
1	25	3	1	0	0	1950	7747	1890	1			
2	25	3	1	0	0	1950	2588	1630	1			
3	25	3	1	0	0	1950	4772	1890	1			
4	25	3	1	0	0	1950	2595	1630	1			
5	25	3	1	0	0	1950	7747	1890	1			
6	25	3	1	0	0	1950	2588	1630	1			
7	25	3	1	0	0	1950	4772	1890	1			
8	25	3	1	0	0	1950	2595	1630	1			
9	25	3	1	0	0	1950	7747	1890	1			
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2	25	3	1	0	0	1950	2595	1630	1			
3	25	3	1	0	0	1950	7747	1890	1			
4	25	3	1	0	0	1950	2588	1630	1			
5	25	3	1	0	0	1950	4772	1890	1			
6	25	3	1	0	0	1950	2595	1630	1			
7	25	3	1	0	0	1950	7747	1890	1			
8	25	3	1	0	0	1950	2588	1630	1			
9	25	3	1	0	0	1950	4772	1890	1			
0	25	3	1	0	0	1950	2595	1630	1			
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3	25	3	1	0	0	1950	4772	1890	1			
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6	25	3	1	0	0	1950	2588	1630	1			
7	25	3	1	0	0	1950	4772	1890	1			
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9	25	3	1	0	0	1950	7747	1890	1			
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3	25	3	1	0	0	1950	7747	1890	1			
4	25	3	1	0	0	1950	2588	1630	1			
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6	25	3	1	0	0	1950	2595	1630	1			
7	25	3	1	0	0	1950	7747	1890	1			
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6	25	3	1	0	0	1950	2588	1630	1			
7	25	3	1	0	0	1950	4772	1890	1			
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IO	NO	R.A.	(1950.0)	DEC.	X	Y	V25	R25	B25	TYPE
301	1	00	15	15	1984730	1529155	16.28	16.03	16.78	161S
302	2	00	15	15	1426815	1309189	16.28			161S
303	3	00	15	15	1179790	1203535	16.28			
304	4	00	15	15	1144100	2196016	16.28			
305	5	00	15	15	1165945	1133976	16.29			SP
306	6	00	15	15	1720428	2254400	16.29			
307	7	00	15	15	1492185	2872400	16.29	14.44	16.80	IRR
308	8	00	15	15	1055035	1936852	16.31			E4
309	9	00	15	15	6351805	1628525	16.32			E2
310	10	00	15	15	1818205	3994460	16.33			SR (S)
311	11	00	15	15	1074080	4404385	16.33			SR
312	12	00	15	15	1430670	4403265	16.34			D
313	13	00	15	15	7276515	1895305	16.35			D
314	14	00	15	15	3146150	1455050	16.35			
315	15	00	15	15	1864670	1633115	16.35	13.74	17.80	161S
316	16	00	15	15	1715150	1633115	16.35	13.99	17.78	
317	17	00	15	15	1371520	1203630	16.36			
318	18	00	15	15	1341165	1556335	16.37			
319	19	00	15	15	1241335	1547455	16.37			
320	20	00	15	15	1673674	1189660	16.37			
321	21	00	15	15	1636740	1993813	16.37			
322	22	00	15	15	1371520	1968855	16.38			
323	23	00	15	15	1045495	1588855	16.38	11.36	17.17	161S
324	24	00	15	15	2454955	1542060	16.38	8.50	17.14	
325	25	00	15	15	1517240	1876938	16.39	16.89	17.36	
326	26	00	15	15	1672075	405870	16.39	14.12	17.48	
327	27	00	15	15	1797240	1037460	16.39			
328	28	00	15	15	1567700	2508660	16.39			
329	29	00	15	15	3884335	2555727	16.40			
330	30	00	15	15	2022920	1906640	16.40			
331	31	00	15	15	1102920	1606640	16.40	14.57	17.43	SP
332	32	00	15	15	9896955	1624320	16.40	15.88	17.16	
333	33	00	15	15	9326955	1533790	16.40	13.78	17.35	
334	34	00	15	15	9026955	1756845	16.40			
335	35	00	15	15	9026955	1708165	16.41			
336	36	00	15	15	223875	2568835	16.41			
337	37	00	15	15	1744450	1557530	16.42	14.54	17.88	D
338	38	00	15	15	8407955	1711830	16.42	11.09	16.91	
339	39	00	15	15	6722500	1266185	16.43			
340	40	00	15	15	5533500	1444195	16.43			161S
341	41	00	15	15	2126735	1054820	16.44			
342	42	00	15	15	1764185	1719210	16.44			
343	43	00	15	15	1316855	1654640	16.44			
344	44	00	15	15	5533500	1515100	16.44			261S
345	45	00	15	15	2062345	1755895	16.46	12.09	16.63	161S
346	46	00	15	15	1796320	2145589	16.46			





ID	NO	R-A	H	M	S	DEC	X	Y	V25	R25	B25	TYPE
01	01	33	37	29	11	17.55	1316545	1382619	16.59	15.36	17.58	
02	02	33	37	29	11	17.55	1316545	1382619	16.59	17.64	18.29	
03	03	33	37	29	11	17.55	1316545	1382619	16.59			
04	04	33	37	29	11	17.55	1316545	1382619	16.59			
05	05	33	37	29	11	17.55	1316545	1382619	16.59			
06	06	33	37	29	11	17.55	1316545	1382619	16.59			
07	07	33	37	29	11	17.55	1316545	1382619	16.59			
08	08	33	37	29	11	17.55	1316545	1382619	16.59			
09	09	33	37	29	11	17.55	1316545	1382619	16.59			
10	10	33	37	29	11	17.55	1316545	1382619	16.59			
11	11	33	37	29	11	17.55	1316545	1382619	16.59			
12	12	33	37	29	11	17.55	1316545	1382619	16.59			
13	13	33	37	29	11	17.55	1316545	1382619	16.59			
14	14	33	37	29	11	17.55	1316545	1382619	16.59			
15	15	33	37	29	11	17.55	1316545	1382619	16.59			
16	16	33	37	29	11	17.55	1316545	1382619	16.59			
17	17	33	37	29	11	17.55	1316545	1382619	16.59			
18	18	33	37	29	11	17.55	1316545	1382619	16.59			
19	19	33	37	29	11	17.55	1316545	1382619	16.59			
20	20	33	37	29	11	17.55	1316545	1382619	16.59			
21	21	33	37	29	11	17.55	1316545	1382619	16.59			
22	22	33	37	29	11	17.55	1316545	1382619	16.59			
23	23	33	37	29	11	17.55	1316545	1382619	16.59			
24	24	33	37	29	11	17.55	1316545	1382619	16.59			
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26	26	33	37	29	11	17.55	1316545	1382619	16.59			
27	27	33	37	29	11	17.55	1316545	1382619	16.59			
28	28	33	37	29	11	17.55	1316545	1382619	16.59			
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30	30	33	37	29	11	17.55	1316545	1382619	16.59			
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32	32	33	37	29	11	17.55	1316545	1382619	16.59			
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40	40	33	37	29	11	17.55	1316545	1382619	16.59			
41	41	33	37	29	11	17.55	1316545	1382619	16.59			
42	42	33	37	29	11	17.55	1316545	1382619	16.59			
43	43	33	37	29	11	17.55	1316545	1382619	16.59			
44	44	33	37	29	11	17.55	1316545	1382619	16.59			
45	45	33	37	29	11	17.55	1316545	1382619	16.59			
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67	67	33	37	29	11	17.55	1316545	1382619	16.59			
68	68	33	37	29	11	17.55	1316545	1382619	16.59			
69	69	33	37	29	11	17.55	1316545	1382619	16.59			
70	70	33	37	29	11	17.55	1316545	1382619	16.59			
71	71	33	37	29	11	17.55	1316545	1382619	16.59			
72	72	33	37	29	11	17.55	1316545	1382619	16.59			
73	73	33	37	29	11	17.55	1316545	1382619	16.59			
74	74	33	37	29	11	17.55	1316545	1382619	16.59			
75	75	33	37	29	11	17.55	1316545	1382619	16.59			
76	76	33	37	29	11	17.55	1316545	1382619	16.59			
77	77	33	37	29	11	17.55	1316545	1382619	16.59			
78	78	33	37	29	11	17.55	1316545	1382619	16.59			
79	79	33	37	29	11	17.55	1316545	1382619	16.59			
80	80	33	37	29	11	17.55	1316545	1382619	16.59			
81	81	33	37	29	11	17.55	1316545	1382619	16.59			
82	82	33	37	29	11	17.55	1316545	1382619	16.59			
83	83	33	37	29	11	17.55	1316545	1382619	16.59			
84	84	33	37	29	11	17.55	1316545	1382619	16.59			
85	85	33	37	29	11	17.55	1316545	1382619	16.59			
86	86	33	37	29	11	17.55	1316545	1382619	16.59			
87	87	33	37	29	11	17.55	1316545	1382619	16.59			
88	88	33	37	29	11	17.55	1316545	1382619	16.59			
89	89	33	37	29	11	17.55	1316545	1382619	16.59			
90	90	33	37	29	11	17.55	1316545	1382619	16.59			
91	91	33	37	29	11	17.55	1316545	1382619	16.59			
92	92	33	37	29	11	17.55	1316545	1382619	16.59			
93	93	33	37	29	11	17.55	1316545	1382619	16.59			
94	94	33	37	29	11	17.55	1316545	1382619	16.59			
95	95	33	37	29	11	17.55	1316545	1382619	16.59			
96	96	33	37	29	11	17.55	1316545	1382619	16.59			
97	97	33	37	29	11	17.55	1316545	1382619	16.59			
98	98	33	37	29	11	17.55	1316545	1382619	16.59			
99	99	33	37	29	11	17.55	1316545	1382619	16.59			
100	100	33	37	29	11	17.55	1316545	1382619	16.59			

10	NO	R-A	H	S	DEC	X	Y	V25	R25	B25	TYPE
1	1	0	0	1	17	175011565	1722425	16.68			SB
1	2	0	0	1	17	160039531	171800853	16.68			S
1	3	0	0	1	17	193755210	207266330	16.68	11.98	17.43	S
1	4	0	0	1	17	236718700	142980050	16.69			S
1	5	0	0	1	17	151233000	144080100	16.69	11.74	18.21	S
1	6	0	0	1	17	12808864	3181245	16.69			S
1	7	0	0	1	17	112487785	208423110	16.69			SR
1	8	0	0	1	17	110221150	184502930	16.69			E1
1	9	0	0	1	17	233487735	245574650	16.70			
1	10	0	0	1	17	223042220	154574755	16.70			
1	11	0	0	1	17	173730200	121112060	16.70	12.02	18.24	SP
1	12	0	0	1	17	179110250	193010560	16.70	18.63		D
1	13	0	0	1	17	67011540	193010560	16.70	12.76	17.99	S
1	14	0	0	1	17	54115855	138225150	16.70			S
1	15	0	0	1	17	168724600	134783350	16.71	11.23	18.01	SP
1	16	0	0	1	17	101846250	157771755	16.71			
1	17	0	0	1	17	142228250	216345530	16.71			
1	18	0	0	1	17	182651505	161305080	16.72			
1	19	0	0	1	17	125037950	179079980	16.72	10.79	17.80	D
1	20	0	0	1	17	178383305	146509490	16.72	10.99	17.85	
1	21	0	0	1	17	14624850	102334950	16.73	12.15	17.86	
1	22	0	0	1	17	118899750	171667331	16.73	12.10	17.28	
1	23	0	0	1	17	187077506	190667370	16.74			SAR
1	24	0	0	1	17	59290355	190667370	16.74			S
1	25	0	0	1	17	184507350	176337750	16.75			PEC
1	26	0	0	1	17	138073300	160876335	16.75	12.03	17.34	SP
1	27	0	0	1	17	20213650	157700955	16.75			S
1	28	0	0	1	17	18013650	152299205	16.76			
1	29	0	0	1	17	175771100	159033105	16.76			
1	30	0	0	1	17	136807350	130533600	16.76			
1	31	0	0	1	17	126071500	114324550	16.76	14.93	17.95	IRR
1	32	0	0	1	17	7555700	114324550	16.76	11.09	13.03	
1	33	0	0	1	17	6775700	154329950	16.77			SR (S)
1	34	0	0	1	17	23987900	18688320	16.77			
1	35	0	0	1	17	18779700	18688320	16.77			

IO NO	R.A.	(1950.0) DEC.	X	Y	V25	R25	B25	TYPE
1	155555	17.78	16.77	298675	16.77			
2	155555	17.78	16.77	298675	16.77			
3	155555	17.78	16.77	298675	16.77			
4	155555	17.78	16.77	298675	16.77			
5	155555	17.78	16.77	298675	16.77			
6	155555	17.78	16.77	298675	16.77			
7	155555	17.78	16.77	298675	16.77			
8	155555	17.78	16.77	298675	16.77			
9	155555	17.78	16.77	298675	16.77			
10	155555	17.78	16.77	298675	16.77			
11	155555	17.78	16.77	298675	16.77			
12	155555	17.78	16.77	298675	16.77			
13	155555	17.78	16.77	298675	16.77			
14	155555	17.78	16.77	298675	16.77			
15	155555	17.78	16.77	298675	16.77			
16	155555	17.78	16.77	298675	16.77			
17	155555	17.78	16.77	298675	16.77			
18	155555	17.78	16.77	298675	16.77			
19	155555	17.78	16.77	298675	16.77			
20	155555	17.78	16.77	298675	16.77			
21	155555	17.78	16.77	298675	16.77			
22	155555	17.78	16.77	298675	16.77			
23	155555	17.78	16.77	298675	16.77			
24	155555	17.78	16.77	298675	16.77			
25	155555	17.78	16.77	298675	16.77			
26	155555	17.78	16.77	298675	16.77			
27	155555	17.78	16.77	298675	16.77			
28	155555	17.78	16.77	298675	16.77			
29	155555	17.78	16.77	298675	16.77			
30	155555	17.78	16.77	298675	16.77			
31	155555	17.78	16.77	298675	16.77			
32	155555	17.78	16.77	298675	16.77			
33	155555	17.78	16.77	298675	16.77			
34	155555	17.78	16.77	298675	16.77			
35	155555	17.78	16.77	298675	16.77			
36	155555	17.78	16.77	298675	16.77			
37	155555	17.78	16.77	298675	16.77			
38	155555	17.78	16.77	298675	16.77			
39	155555	17.78	16.77	298675	16.77			
40	155555	17.78	16.77	298675	16.77			
41	155555	17.78	16.77	298675	16.77			
42	155555	17.78	16.77	298675	16.77			
43	155555	17.78	16.77	298675	16.77			
44	155555	17.78	16.77	298675	16.77			
45	155555	17.78	16.77	298675	16.77			
46	155555	17.78	16.77	298675	16.77			
47	155555	17.78	16.77	298675	16.77			
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49	155555	17.78	16.77	298675	16.77			
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53	155555	17.78	16.77	298675	16.77			
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57	155555	17.78	16.77	298675	16.77			
58	155555	17.78	16.77	298675	16.77			
59	155555	17.78	16.77	298675	16.77			
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61	155555	17.78	16.77	298675	16.77			
62	155555	17.78	16.77	298675	16.77			
63	155555	17.78	16.77	298675	16.77			
64	155555	17.78	16.77	298675	16.77			
65	155555	17.78	16.77	298675	16.77			
66	155555	17.78	16.77	298675	16.77			
67	155555	17.78	16.77	298675	16.77			
68	155555	17.78	16.77	298675	16.77			
69	155555	17.78	16.77	298675	16.77			
70	155555	17.78	16.77	298675	16.77			
71	155555	17.78	16.77	298675	16.77			
72	155555	17.78	16.77	298675	16.77			
73	155555	17.78	16.77	298675	16.77			
74	155555	17.78	16.77	298675	16.77			
75	155555	17.78	16.77	298675	16.77			
76	155555	17.78	16.77	298675	16.77			
77	155555	17.78	16.77	298675	16.77			
78	155555	17.78	16.77	298675	16.77			
79	155555	17.78	16.77	298675	16.77			
80	155555	17.78	16.77	298675	16.77			
81	155555	17.78	16.77	298675	16.77			
82	155555	17.78	16.77	298675	16.77			
83	155555	17.78	16.77	298675	16.77			
84	155555	17.78	16.77	298675	16.77			
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86	155555	17.78	16.77	298675	16.77			
87	155555	17.78	16.77	298675	16.77			
88	155555	17.78	16.77	298675	16.77			
89	155555	17.78	16.77	298675	16.77			
90	155555	17.78	16.77	298675	16.77			
91	155555	17.78	16.77	298675	16.77			
92	155555	17.78	16.77	298675	16.77			
93	155555	17.78	16.77	298675	16.77			
94	155555	17.78	16.77	298675	16.77			
95	155555	17.78	16.77	298675	16.77			
96	155555	17.78	16.77	298675	16.77			
97	155555	17.78	16.77	298675	16.77			
98	155555	17.78	16.77	298675	16.77			
99	155555	17.78	16.77	298675	16.77			
100	155555	17.78	16.77	298675	16.77			

NO	H	A	S	DEC	X	Y	V25	R25	B25	TYPE
1	1	1	1	1	1	1	16.87	9.51	17.54	
2	1	1	1	1	1	1	16.87	12.16	16.17	S
3	1	1	1	1	1	1	16.87			
4	1	1	1	1	1	1	16.87			
5	1	1	1	1	1	1	16.87			
6	1	1	1	1	1	1	16.87			
7	1	1	1	1	1	1	16.87			
8	1	1	1	1	1	1	16.87			
9	1	1	1	1	1	1	16.87			
10	1	1	1	1	1	1	16.87			
11	1	1	1	1	1	1	16.87			
12	1	1	1	1	1	1	16.87			
13	1	1	1	1	1	1	16.87			
14	1	1	1	1	1	1	16.87			
15	1	1	1	1	1	1	16.87			
16	1	1	1	1	1	1	16.87			
17	1	1	1	1	1	1	16.87			
18	1	1	1	1	1	1	16.87			
19	1	1	1	1	1	1	16.87			
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21	1	1	1	1	1	1	16.87			
22	1	1	1	1	1	1	16.87			
23	1	1	1	1	1	1	16.87			
24	1	1	1	1	1	1	16.87			
25	1	1	1	1	1	1	16.87			
26	1	1	1	1	1	1	16.87			
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29	1	1	1	1	1	1	16.87			
30	1	1	1	1	1	1	16.87			
31	1	1	1	1	1	1	16.87			
32	1	1	1	1	1	1	16.87			
33	1	1	1	1	1	1	16.87			
34	1	1	1	1	1	1	16.87			
35	1	1	1	1	1	1	16.87			
36	1	1	1	1	1	1	16.87			
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38	1	1	1	1	1	1	16.87			
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41	1	1	1	1	1	1	16.87			
42	1	1	1	1	1	1	16.87			
43	1	1	1	1	1	1	16.87			
44	1	1	1	1	1	1	16.87			
45	1	1	1	1	1	1	16.87			
46	1	1	1	1	1	1	16.87			
47	1	1	1	1	1	1	16.87			
48	1	1	1	1	1	1	16.87			
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51	1	1	1	1	1	1	16.87			
52	1	1	1	1	1	1	16.87			
53	1	1	1	1	1	1	16.87			
54	1	1	1	1	1	1	16.87			
55	1	1	1	1	1	1	16.87			
56	1	1	1	1	1	1	16.87			
57	1	1	1	1	1	1	16.87			



IO	NO	R.A.	DEC.	X	Y	V25	R25	B25	TYPE
01	001	10 00 38	55 00 35	710740	2298025	16.96			
02	002	10 00 34	55 01 35	315015	2202215	16.96			
03	003	10 00 34	55 02 13	1250335	22376360	16.97			
04	004	10 00 34	55 03 13	1022345	22439160	16.97			
05	005	10 00 34	55 03 28	364918	22221002	16.97			
06	006	10 00 34	55 04 17	2065905	22221008	16.98			
07	007	10 00 34	55 05 05	1933309	21922068	16.98			
08	008	10 00 34	55 05 33	1617780	20821780	16.98			
09	009	10 00 34	55 06 07	877450	21731380	16.98			
10	010	10 00 34	55 06 33	5211770	21797330	16.98			
11	011	10 00 34	55 07 07	1963331	21797390	16.99			
12	012	10 00 34	55 07 33	1294990	2085080	16.99			
13	013	10 00 34	55 08 07	1094035	20872200	16.99	10.04	18.09	
14	014	10 00 34	55 08 33	394775	21518560	16.99			
15	015	10 00 34	55 09 07	2315720	20872200	16.99			
16	016	10 00 34	55 09 33	2307475	20999885	16.99			
17	017	10 00 34	55 10 07	1820544	21518560	16.99			
18	018	10 00 34	55 10 33	1705769	20999885	16.99			
19	019	10 00 34	55 11 07	1557795	21518560	16.99			
20	020	10 00 34	55 11 33	1435235	20999885	16.99			
21	021	10 00 34	55 12 07	1337880	21518560	16.99			
22	022	10 00 34	55 12 33	1978810	20999885	16.99			
23	023	10 00 34	55 13 07	6252470	21518560	16.99			
24	024	10 00 34	55 13 33	5593215	20999885	16.99			
25	025	10 00 34	55 14 07	1206915	21518560	16.99			
26	026	10 00 34	55 14 33	3023755	20999885	16.99			
27	027	10 00 34	55 15 07	1631095	21518560	16.99			
28	028	10 00 34	55 15 33	2395550	20999885	16.99			
29	029	10 00 34	55 16 07	1895460	21518560	16.99			
30	030	10 00 34	55 16 33	1772195	20999885	16.99			
31	031	10 00 34	55 17 07	1601950	21518560	16.99			
32	032	10 00 34	55 17 33	1585550	20999885	16.99			
33	033	10 00 34	55 18 07	1535520	21518560	16.99			
34	034	10 00 34	55 18 33	1474720	20999885	16.99			
35	035	10 00 34	55 19 07	1232050	21518560	16.99			
36	036	10 00 34	55 19 33	1833365	20999885	16.99			
37	037	10 00 34	55 20 07	505965	21518560	16.99			
38	038	10 00 34	55 20 33	270610	20999885	16.99			
39	039	10 00 34	55 21 07	249455	21518560	16.99			
40	040	10 00 34	55 21 33	150355	20999885	16.99			
41	041	10 00 34	55 22 07	215035	21518560	16.99			
42	042	10 00 34	55 22 33	150355	20999885	16.99			
43	043	10 00 34	55 23 07	215035	21518560	16.99			
44	044	10 00 34	55 23 33	150355	20999885	16.99			
45	045	10 00 34	55 24 07	215035	21518560	16.99			
46	046	10 00 34	55 24 33	150355	20999885	16.99			
47	047	10 00 34	55 25 07	215035	21518560	16.99			
48	048	10 00 34	55 25 33	150355	20999885	16.99			
49	049	10 00 34	55 26 07	215035	21518560	16.99			
50	050	10 00 34	55 26 33	150355	20999885	16.99			
51	051	10 00 34	55 27 07	215035	21518560	16.99			
52	052	10 00 34	55 27 33	150355	20999885	16.99			
53	053	10 00 34	55 28 07	215035	21518560	16.99			
54	054	10 00 34	55 28 33	150355	20999885	16.99			
55	055	10 00 34	55 29 07	215035	21518560	16.99			
56	056	10 00 34	55 29 33	150355	20999885	16.99			
57	057	10 00 34	55 30 07	215035	21518560	16.99			
58	058	10 00 34	55 30 33	150355	20999885	16.99			
59	059	10 00 34	55 31 07	215035	21518560	16.99			
60	060	10 00 34	55 31 33	150355	20999885	16.99			
61	061	10 00 34	55 32 07	215035	21518560	16.99			
62	062	10 00 34	55 32 33	150355	20999885	16.99			
63	063	10 00 34	55 33 07	215035	21518560	16.99			
64	064	10 00 34	55 33 33	150355	20999885	16.99			
65	065	10 00 34	55 34 07	215035	21518560	16.99			
66	066	10 00 34	55 34 33	150355	20999885	16.99			
67	067	10 00 34	55 35 07	215035	21518560	16.99			
68	068	10 00 34	55 35 33	150355	20999885	16.99			
69	069	10 00 34	55 36 07	215035	21518560	16.99			
70	070	10 00 34	55 36 33	150355	20999885	16.99			
71	071	10 00 34	55 37 07	215035	21518560	16.99			
72	072	10 00 34	55 37 33	150355	20999885	16.99			
73	073	10 00 34	55 38 07	215035	21518560	16.99			
74	074	10 00 34	55 38 33	150355	20999885	16.99			
75	075	10 00 34	55 39 07	215035	21518560	16.99			
76	076	10 00 34	55 39 33	150355	20999885	16.99			
77	077	10 00 34	55 40 07	215035	21518560	16.99			
78	078	10 00 34	55 40 33	150355	20999885	16.99			
79	079	10 00 34	55 41 07	215035	21518560	16.99			
80	080	10 00 34	55 41 33	150355	20999885	16.99			
81	081	10 00 34	55 42 07	215035	21518560	16.99			
82	082	10 00 34	55 42 33	150355	20999885	16.99			
83	083	10 00 34	55 43 07	215035	21518560	16.99			
84	084	10 00 34	55 43 33	150355	20999885	16.99			
85	085	10 00 34	55 44 07	215035	21518560	16.99			
86	086	10 00 34	55 44 33	150355	20999885	16.99			
87	087	10 00 34	55 45 07	215035	21518560	16.99			
88	088	10 00 34	55 45 33	150355	20999885	16.99			
89	089	10 00 34	55 46 07	215035	21518560	16.99			
90	090	10 00 34	55 46 33	150355	20999885	16.99			
91	091	10 00 34	55 47 07	215035	21518560	16.99			
92	092	10 00 34	55 47 33	150355	20999885	16.99			
93	093	10 00 34	55 48 07	215035	21518560	16.99			
94	094	10 00 34	55 48 33	150355	20999885	16.99			
95	095	10 00 34	55 49 07	215035	21518560	16.99			
96	096	10 00 34	55 49 33	150355	20999885	16.99			
97	097	10 00 34	55 50 07	215035	21518560	16.99			
98	098	10 00 34	55 50 33	150355	20999885	16.99			
99	099	10 00 34	55 51 07	215035	21518560	16.99			
100	100	10 00 34	55 51 33	150355	20999885	16.99			





IO NO	R.A.	(1950.0) DEC.	X	Y	V25	R25	B25	TYPE
701	10 00 00	5 45 30	206650	214610	11			
702	10 00 00	5 45 30	206650	214610	11			
703	10 00 00	5 45 30	206650	214610	11			
704	10 00 00	5 45 30	206650	214610	11			
705	10 00 00	5 45 30	206650	214610	11			
706	10 00 00	5 45 30	206650	214610	11			
707	10 00 00	5 45 30	206650	214610	11			
708	10 00 00	5 45 30	206650	214610	11			
709	10 00 00	5 45 30	206650	214610	11			
710	10 00 00	5 45 30	206650	214610	11			
711	10 00 00	5 45 30	206650	214610	11			
712	10 00 00	5 45 30	206650	214610	11			
713	10 00 00	5 45 30	206650	214610	11			
714	10 00 00	5 45 30	206650	214610	11			
715	10 00 00	5 45 30	206650	214610	11			
716	10 00 00	5 45 30	206650	214610	11			
717	10 00 00	5 45 30	206650	214610	11			
718	10 00 00	5 45 30	206650	214610	11			
719	10 00 00	5 45 30	206650	214610	11			
720	10 00 00	5 45 30	206650	214610	11			
721	10 00 00	5 45 30	206650	214610	11			
722	10 00 00	5 45 30	206650	214610	11			
723	10 00 00	5 45 30	206650	214610	11			
724	10 00 00	5 45 30	206650	214610	11			
725	10 00 00	5 45 30	206650	214610	11			
726	10 00 00	5 45 30	206650	214610	11			
727	10 00 00	5 45 30	206650	214610	11			
728	10 00 00	5 45 30	206650	214610	11			
729	10 00 00	5 45 30	206650	214610	11			
730	10 00 00	5 45 30	206650	214610	11			
731	10 00 00	5 45 30	206650	214610	11			
732	10 00 00	5 45 30	206650	214610	11			
733	10 00 00	5 45 30	206650	214610	11			
734	10 00 00	5 45 30	206650	214610	11			
735	10 00 00	5 45 30	206650	214610	11			
736	10 00 00	5 45 30	206650	214610	11			
737	10 00 00	5 45 30	206650	214610	11			
738	10 00 00	5 45 30	206650	214610	11			
739	10 00 00	5 45 30	206650	214610	11			
740	10 00 00	5 45 30	206650	214610	11			
741	10 00 00	5 45 30	206650	214610	11			
742	10 00 00	5 45 30	206650	214610	11			
743	10 00 00	5 45 30	206650	214610	11			
744	10 00 00	5 45 30	206650	214610	11			
745	10 00 00	5 45 30	206650	214610	11			
746	10 00 00	5 45 30	206650	214610	11			
747	10 00 00	5 45 30	206650	214610	11			
748	10 00 00	5 45 30	206650	214610	11			
749	10 00 00	5 45 30	206650	214610	11			
750	10 00 00	5 45 30	206650	214610	11			

ID	NO	R.A.	H	M	S	DEC.	X	Y	V25	R25	B25	TYPE
751	2	34	47	10	33	42	33	360	17.18	8.21	17.95	1G1S
752	3	33	33	10	33	42	33	370	17.18			
753	4	33	33	10	33	42	33	375	17.18			
754	5	33	33	10	33	42	33	380	17.18			
755	6	33	33	10	33	42	33	385	17.18			
756	7	33	33	10	33	42	33	390	17.18			
757	8	33	33	10	33	42	33	395	17.18			
758	9	33	33	10	33	42	33	400	17.18			
759	10	33	33	10	33	42	33	405	17.18			
760	11	33	33	10	33	42	33	410	17.18			
761	12	33	33	10	33	42	33	415	17.18			
762	13	33	33	10	33	42	33	420	17.18			
763	14	33	33	10	33	42	33	425	17.18			
764	15	33	33	10	33	42	33	430	17.18			
765	16	33	33	10	33	42	33	435	17.18			
766	17	33	33	10	33	42	33	440	17.18			
767	18	33	33	10	33	42	33	445	17.18			
768	19	33	33	10	33	42	33	450	17.18			
769	20	33	33	10	33	42	33	455	17.18			
770	21	33	33	10	33	42	33	460	17.18			
771	22	33	33	10	33	42	33	465	17.18			
772	23	33	33	10	33	42	33	470	17.18			
773	24	33	33	10	33	42	33	475	17.18			
774	25	33	33	10	33	42	33	480	17.18			
775	26	33	33	10	33	42	33	485	17.18			
776	27	33	33	10	33	42	33	490	17.18			
777	28	33	33	10	33	42	33	495	17.18			
778	29	33	33	10	33	42	33	500	17.18			
779	30	33	33	10	33	42	33	505	17.18			
780	31	33	33	10	33	42	33	510	17.18			
781	32	33	33	10	33	42	33	515	17.18			
782	33	33	33	10	33	42	33	520	17.18			
783	34	33	33	10	33	42	33	525	17.18			
784	35	33	33	10	33	42	33	530	17.18			
785	36	33	33	10	33	42	33	535	17.18			
786	37	33	33	10	33	42	33	540	17.18			
787	38	33	33	10	33	42	33	545	17.18			
788	39	33	33	10	33	42	33	550	17.18			
789	40	33	33	10	33	42	33	555	17.18			
790	41	33	33	10	33	42	33	560	17.18			
791	42	33	33	10	33	42	33	565	17.18			
792	43	33	33	10	33	42	33	570	17.18			
793	44	33	33	10	33	42	33	575	17.18			
794	45	33	33	10	33	42	33	580	17.18			
795	46	33	33	10	33	42	33	585	17.18			
796	47	33	33	10	33	42	33	590	17.18			
797	48	33	33	10	33	42	33	595	17.18			
798	49	33	33	10	33	42	33	600	17.18			
799	50	33	33	10	33	42	33	605	17.18			
800	51	33	33	10	33	42	33	610	17.18			
801	52	33	33	10	33	42	33	615	17.18			
802	53	33	33	10	33	42	33	620	17.18			
803	54	33	33	10	33	42	33	625	17.18			
804	55	33	33	10	33	42	33	630	17.18			
805	56	33	33	10	33	42	33	635	17.18			
806	57	33	33	10	33	42	33	640	17.18			
807	58	33	33	10	33	42	33	645	17.18			
808	59	33	33	10	33	42	33	650	17.18			
809	60	33	33	10	33	42	33	655	17.18			
810	61	33	33	10	33	42	33	660	17.18			
811	62	33	33	10	33	42	33	665	17.18			
812	63	33	33	10	33	42	33	670	17.18			
813	64	33	33	10	33	42	33	675	17.18			
814	65	33	33	10	33	42	33	680	17.18			
815	66	33	33	10	33	42	33	685	17.18			
816	67	33	33	10	33	42	33	690	17.18			
817	68	33	33	10	33	42	33	695	17.18			
818	69	33	33	10	33	42	33	700	17.18			
819	70	33	33	10	33	42	33	705	17.18			
820	71	33	33	10	33	42	33	710	17.18			
821	72	33	33	10	33	42	33	715	17.18			
822	73	33	33	10	33	42	33	720	17.18			
823	74	33	33	10	33	42	33	725	17.18			
824	75	33	33	10	33	42	33	730	17.18			
825	76	33	33	10	33	42	33	735	17.18			
826	77	33	33	10	33	42	33	740	17.18			
827	78	33	33	10	33	42	33	745	17.18			
828	79	33	33	10	33	42	33	750	17.18			
829	80	33	33	10	33	42	33	755	17.18			
830	81	33	33	10	33	42	33	760	17.18			
831	82	33	33	10	33	42	33	765	17.18			
832	83	33	33	10	33	42	33	770	17.18			
833	84	33	33	10	33	42	33	775	17.18			
834	85	33	33	10	33	42	33	780	17.18			
835	86	33	33	10	33	42	33	785	17.18			
836	87	33	33	10	33	42	33	790	17.18			
837	88	33	33	10	33	42	33	795	17.18			
838	89	33	33	10	33	42	33	800	17.18			
839	90	33	33	10	33	42	33	805	17.18			
840	91	33	33	10	33	42	33	810	17.18			
841	92	33	33	10	33	42	33	815	17.18			
842	93	33	33	10	33	42	33	820	17.18			
843	94	33	33	10	33	42	33	825	17.18			
844	95	33	33	10	33	42	33	830	17.18			
845	96	33	33	10	33	42	33	835	17.18			
846	97	33	33	10	33	42	33	840	17.18			
847	98	33	33	10	33	42	33	845	17.18			
848	99	33	33	10	33	42	33	850	17.18			
849	100	33	33	10	33	42	33	855	17.18			
850	101	33	33	10	33	42	33	860	17.18			
851	102	33	33	10	33	42	33	865	17.18			
852	103	33	33	10	33	42	33	870	17.18			
853	104	33	33	10	33	42	33	875	17.18			
854	105	33	33	10	33	42	33	880	17.18			
855	106	33	33	10	33	42	33	885	17.18			
856	107	33	33	10	33	42	33	890	17.18			
857	108	33	33	10	33	42	33	895	17.18			
858	109	33	33	10	33	42	33	900	17.18			
859	110	33	33	10	33	42	33	905	17.18			
860	111	33	33	10	33	42	33	910	17.18			
861	112	33	33	10	33	42	33	915	17.18			
862	113	33	33	10	33	42	33	920	17.18			
863	114	33	33	10	33	42	33	925	17.18			
864	115	33	33	10	33	42	33	930	17.18			
865	116	33	33	10	33	42	33	935	17.18			
866	117	33	33	10	33	42	33	940	17.18			
867	118	33	33	10	33	42	33	945	17.18			
868	119	33	33	10	33	42	33	950	17.18			
869	120	33	33	10	33	42	33	955	17.18			
870	121	33	33	10	33	42	33	960	17.18			
871	122	33	33	10	33	42	33	965	17.18			
872	123	33	33	10	33	42	33	970	17.18			
873	124	33	33	10	33	42	33	975	17.18			
874	125	33	33	10	33	42	33	980	17.18			
875	126	33	33	10	33	42	33	985	17.18			
876	127	33	33	10	33	42	33	990	17.18			
877	128	33	33	10	33	42	33	995	17.18			
878	129	33	33	10	33	42	33	1000	17.18			
879	130	33	33	10	33	42	33	1005	17.18			
880	131	33	33	10	33	42	33	1010	17.18			
881	132	33	33	10	33	42	33	1015	17.18			
882												

IO NO	R-A. H	(1950.0) DEC.	X	Y	V25	R25	B25	TYPE
301	480	12-15-11	2337	7990	17.27			
302	233	12-15-11	1937	2255	17.27			
303	332	12-15-11	1752	1435	17.27			
304	332	12-15-11	1528	1540	17.27			
305	332	12-15-11	1528	1540	17.27			
306	332	12-15-11	1528	1540	17.27			
307	332	12-15-11	1528	1540	17.27			
308	332	12-15-11	1528	1540	17.27			
309	332	12-15-11	1528	1540	17.27			
310	332	12-15-11	1528	1540	17.27			
311	332	12-15-11	1528	1540	17.27			
312	332	12-15-11	1528	1540	17.27			
313	332	12-15-11	1528	1540	17.27			
314	332	12-15-11	1528	1540	17.27			
315	332	12-15-11	1528	1540	17.27			
316	332	12-15-11	1528	1540	17.27			
317	332	12-15-11	1528	1540	17.27			
318	332	12-15-11	1528	1540	17.27			
319	332	12-15-11	1528	1540	17.27			
320	332	12-15-11	1528	1540	17.27			
321	332	12-15-11	1528	1540	17.27			
322	332	12-15-11	1528	1540	17.27			
323	332	12-15-11	1528	1540	17.27			
324	332	12-15-11	1528	1540	17.27			
325	332	12-15-11	1528	1540	17.27			
326	332	12-15-11	1528	1540	17.27			
327	332	12-15-11	1528	1540	17.27			
328	332	12-15-11	1528	1540	17.27			
329	332	12-15-11	1528	1540	17.27			
330	332	12-15-11	1528	1540	17.27			
331	332	12-15-11	1528	1540	17.27			
332	332	12-15-11	1528	1540	17.27			
333	332	12-15-11	1528	1540	17.27			
334	332	12-15-11	1528	1540	17.27			
335	332	12-15-11	1528	1540	17.27			
336	332	12-15-11	1528	1540	17.27			
337	332	12-15-11	1528	1540	17.27			
338	332	12-15-11	1528	1540	17.27			
339	332	12-15-11	1528	1540	17.27			
340	332	12-15-11	1528	1540	17.27			
341	332	12-15-11	1528	1540	17.27			
342	332	12-15-11	1528	1540	17.27			
343	332	12-15-11	1528	1540	17.27			
344	332	12-15-11	1528	1540	17.27			
345	332	12-15-11	1528	1540	17.27			
346	332	12-15-11	1528	1540	17.27			
347	332	12-15-11	1528	1540	17.27			
348	332	12-15-11	1528	1540	17.27			
349	332	12-15-11	1528	1540	17.27			
350	332	12-15-11	1528	1540	17.27			

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SR

17.90

10.03

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7.58

18.13

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8.10

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IO	NO	R-A	DEC	X	Y	V25	R25	B25	TYPE
301	301	10	10	10	10	17.40	9.60		
302	302	10	10	10	10	17.40			
303	303	10	10	10	10	17.40			
304	304	10	10	10	10	17.40			
305	305	10	10	10	10	17.40			
306	306	10	10	10	10	17.40			
307	307	10	10	10	10	17.40			
308	308	10	10	10	10	17.40			
309	309	10	10	10	10	17.40			
310	310	10	10	10	10	17.40			
311	311	10	10	10	10	17.40			
312	312	10	10	10	10	17.40			
313	313	10	10	10	10	17.40			
314	314	10	10	10	10	17.40			
315	315	10	10	10	10	17.40			
316	316	10	10	10	10	17.40			
317	317	10	10	10	10	17.40			
318	318	10	10	10	10	17.40			
319	319	10	10	10	10	17.40			
320	320	10	10	10	10	17.40			
321	321	10	10	10	10	17.40			
322	322	10	10	10	10	17.40			
323	323	10	10	10	10	17.40			
324	324	10	10	10	10	17.40			
325	325	10	10	10	10	17.40			
326	326	10	10	10	10	17.40			
327	327	10	10	10	10	17.40			
328	328	10	10	10	10	17.40			
329	329	10	10	10	10	17.40			
330	330	10	10	10	10	17.40			
331	331	10	10	10	10	17.40			
332	332	10	10	10	10	17.40			
333	333	10	10	10	10	17.40			
334	334	10	10	10	10	17.40			
335	335	10	10	10	10	17.40			
336	336	10	10	10	10	17.40			
337	337	10	10	10	10	17.40			
338	338	10	10	10	10	17.40			
339	339	10	10	10	10	17.40			
340	340	10	10	10	10	17.40			
341	341	10	10	10	10	17.40			
342	342	10	10	10	10	17.40			
343	343	10	10	10	10	17.40			
344	344	10	10	10	10	17.40			
345	345	10	10	10	10	17.40			
346	346	10	10	10	10	17.40			
347	347	10	10	10	10	17.40			
348	348	10	10	10	10	17.40			
349	349	10	10	10	10	17.40			
350	350	10	10	10	10	17.40			
351	351	10	10	10	10	17.40			
352	352	10	10	10	10	17.40			
353	353	10	10	10	10	17.40			
354	354	10	10	10	10	17.40			
355	355	10	10	10	10	17.40			
356	356	10	10	10	10	17.40			
357	357	10	10	10	10	17.40			
358	358	10	10	10	10	17.40			
359	359	10	10	10	10	17.40			
360	360	10	10	10	10	17.40			
361	361	10	10	10	10	17.40			
362	362	10	10	10	10	17.40			
363	363	10	10	10	10	17.40			
364	364	10	10	10	10	17.40			
365	365	10	10	10	10	17.40			
366	366	10	10	10	10	17.40			
367	367	10	10	10	10	17.40			
368	368	10	10	10	10	17.40			
369	369	10	10	10	10	17.40			
370	370	10	10	10	10	17.40			
371	371	10	10	10	10	17.40			
372	372	10	10	10	10	17.40			
373	373	10	10	10	10	17.40			
374	374	10	10	10	10	17.40			
375	375	10	10	10	10	17.40			
376	376	10	10	10	10	17.40			
377	377	10	10	10	10	17.40			
378	378	10	10	10	10	17.40			
379	379	10	10	10	10	17.40			
380	380	10	10	10	10	17.40			
381	381	10	10	10	10	17.40			
382	382	10	10	10	10	17.40			
383	383	10	10	10	10	17.40			
384	384	10	10	10	10	17.40			
385	385	10	10	10	10	17.40			
386	386	10	10	10	10	17.40			
387	387	10	10	10	10	17.40			
388	388	10	10	10	10	17.40			
389	389	10	10	10	10	17.40			
390	390	10	10	10	10	17.40			
391	391	10	10	10	10	17.40			
392	392	10	10	10	10	17.40			
393	393	10	10	10	10	17.40			
394	394	10	10	10	10	17.40			
395	395	10	10	10	10	17.40			
396	396	10	10	10	10	17.40			
397	397	10	10	10	10	17.40			
398	398	10	10	10	10	17.40			
399	399	10	10	10	10	17.40			
400	400	10	10	10	10	17.40			



IO NO	R.A. H	(1950.0) S	DEC. 0	X	Y	V25	R25	B25	TYPE
351	10	0	33	5293350	89665	17.44			
352	10	0	35	377620	604860	17.44			
353	10	0	35	348740	126155	17.44			
354	10	0	35	163640	155169	17.44			
355	10	0	35	193925	180840	17.45			SP
356	10	0	35	184670	230671	17.45			
357	10	0	35	115626	149045	17.45	9.72		
358	10	0	35	839575	209127	17.45			
359	10	0	35	839663	344550	17.45			SB
360	10	0	35	715970	641650	17.45			
361	10	0	35	121160	252824	17.45			
362	10	0	35	106440	446703	17.45			
363	10	0	35	230913	679071	17.46			
364	10	0	35	192984	133075	17.46			SP
365	10	0	35	170489	133756	17.46			
366	10	0	35	161515	169222	17.46			1G1S
367	10	0	35	431123	163845	17.46			1G1S
368	10	0	35	306390	204358	17.47			
369	10	0	35	218390	106000	17.47			
370	10	0	35	217011	636335	17.47			
371	10	0	35	193456	209587	17.47			
372	10	0	35	124776	580560	17.47			
373	10	0	35	129972	187071	17.47			SP
374	10	0	35	717015	122582	17.47			
375	10	0	35	369102	222051	17.47			
376	10	0	35	209438	169474	17.48			
377	10	0	35	158380	176075	17.48			
378	10	0	35	149380	187951	17.48			
379	10	0	35	112993	203178	17.48			
380	10	0	35	114058	174651	17.48	7.02	18.69	
381	10	0	35	552745	122562	17.48			
382	10	0	35	310145	172561	17.48			
383	10	0	35	166201	158669	17.48			
384	10	0	35	229442	233479	17.49			
385	10	0	35	215370	159022	17.49			
386	10	0	35	153525	153944	17.49			
387	10	0	35	209241	156013	17.49			SP (S)
388	10	0	35	189577	237285	17.49	18.28		
389	10	0	35	155740	137030	17.49			
390	10	0	35	127732	235355	17.49			
391	10	0	35	110453	165777	17.49			
392	10	0	35	104533	146500	17.49			
393	10	0	35	373305	194693	17.49			
394	10	0	35	370650	140693	17.49	8.31	19.00	
395	10	0	35	630640	113466	17.49			
396	10	0	35	154210	183440	17.49			

## SUPPLEMENT

	R.A. h m s	(1950.0) Dec. ° ' "	$V_{25}$	$R_{25}$	$B_{25}$	Type	Comments
A	10 35 02	-25 03.5	1177000	2586000		SB(r)	420" x 370:"
B	10 34 52	-27 25.4	1200985	1319185	12.72	S	= Two spirals superposed
C	10 32 30	-28 19.5	1477000	834600	$\sim 13.2$	SB(s)	160" x 107", largest in Klemola 13
D	10 35 50	-28 31.5	1084500	727900	$\sim 13.9$	SB0:	87" x 74"
E	10 37 01	-29 19.5	948000	298000		SB(rs)	94:" x 50"
F	10 25 20	-29 32.5	2311000	172500		SB(rs)	70:" x 47:"
G	10 34 51	-27 08.6	1199800	1469000	$\sim 15.5$	E:	$\sim 0'.4$ np ✕ (BS 4162)
H	10 35 17	-27 20.4	1149410	1363090	16.0		24" sf and sup. N3316

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REFERENCES

Non-standard abbreviations are listed below (for other commonly used abbreviations and prefixes, consult Table 1.1):

A & Ap	Astronomy & Astrophysics
AJ	Astronomical Journal
Ap J	Astrophysical Journal
Mem. RAS	Memoirs of the Royal Astronomical Society
MNASSA	Monthly Notices of the Astronomical Society of Southern Africa
MNRAS	Monthly Notices of the Royal Astronomical Society
PASP	Publications of the Astronomical Society of the Pacific

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